

CONDENSATION

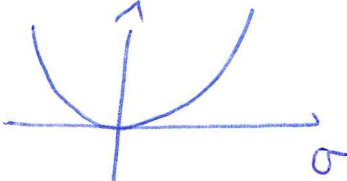
TWO IMPORTANT POINTS:
ABOUT CONDENSATION

CONDENSATION

FLUCTUATION AROUND
THE CONDENSATION

• The "condensation" itself means nothing..

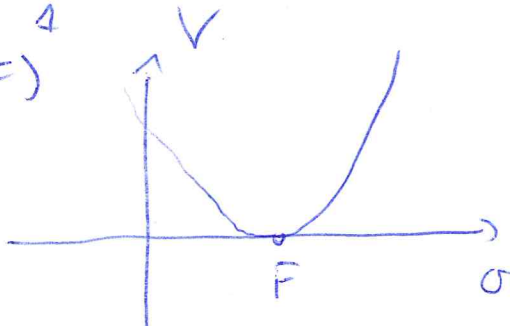
① $V(\sigma) = \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{\lambda}{4} \sigma^4$



Minimum for $\sigma=0$. Mass m_σ . Scattering: λ .

• Now, suppose I gave you

② $V(\sigma) = \frac{1}{2} m_\sigma^2 (\sigma - F)^2 + \frac{\lambda}{4} (\sigma - F)^4$



Now, I have a condensate...

$$\sigma = F!!!$$

But, how can I see it physically? NO WAY!!!

m_σ is the mass, λ the scattering!!!

Theories ① and ② are perfectly identical. The emergence of the condensate in the second case is an artifact of the choice of the γ -axis... (ref. frame choice)!

it's, things change drastically, if we add another field (which we can add)
(mixed) scalar field

$$\textcircled{1} \quad V_1 = \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{\lambda}{4} \sigma^4 + c \sigma^2 \varphi^2$$

$$m_\sigma \quad \times$$

$$m_\varphi = 0 \quad \times$$

$$\textcircled{2} \quad V_2 = \frac{1}{2} m_\sigma^2 (\sigma - F)^2 + \frac{\lambda}{4} (\sigma - F)^4 + c \sigma^2 \varphi^2$$

$$\sigma \mapsto F + \sigma$$

$$V_2 = \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{\lambda}{4} \sigma^4 + c F^2 \varphi^2 + 2c F \sigma \varphi^2 + c \sigma^2 \varphi^2$$

$$m_\sigma \quad \times$$

$$m_\varphi = 2F$$

$$\times$$

$$\textcircled{2} \neq \textcircled{1}$$

$$\textcircled{3} \quad V = \frac{1}{2} m_\sigma^2 (\sigma - F)^2 + \frac{\lambda}{4} (\sigma - F)^4 + c (\sigma - F)^2 (\varphi - \bar{F})^2$$

then the minimum is for $\sigma = F, \varphi = \bar{F} \mapsto$ shift of both $\sigma \mapsto F + \sigma$
 $\varphi \mapsto \bar{F} + \varphi$

$\textcircled{3} = 1$ (Both fields condense, but you don't see that...)

How, things change drastically if we add another field (not a fermion field)

$$\textcircled{1} \quad V_1(\sigma) = \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{\lambda}{4} \sigma^4 + c \sigma \bar{\psi} \psi$$

$$m_\sigma, \lambda, \quad \times$$

$$m_\psi = 0, \quad \frac{\cancel{\psi}}{c} \frac{\cancel{\sigma}}{\psi}$$

c dimensionless...

$$\textcircled{2} \quad V_2(\sigma) = \frac{1}{2} m_\sigma (\sigma - F)^2 + \frac{\lambda}{4} (\sigma - F)^4 + c \sigma \bar{\psi} \psi$$

$$\sigma \mapsto F + \sigma$$

$$\begin{aligned} V(\sigma) &= \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{\lambda}{4} \sigma^4 + c(\sigma + F) \bar{\psi} \psi \\ &= \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{\lambda}{4} \sigma^4 + \underbrace{c F \bar{\psi} \psi} + c \sigma \bar{\psi} \psi \end{aligned}$$

$$m_\sigma, \lambda$$

$$m_\psi = c F \neq 0!$$

$\textcircled{2}$ is $\neq$ from $\textcircled{1}$.

$\textcircled{3}$ The interactions that arise upon condensation are the physical realisation of it...

B.C.W.: This is exactly the way how the Higgs gives mass to the fermions of the SM. $\rightarrow$ STATCHER ANALOGY!

Back to our σ -model.

(with σ and π ($N_f=1$))

$$V = \frac{\Lambda}{4} (\sigma^2 + \pi^2 - F^2)^2 - \epsilon \sigma$$

$$\sigma \mapsto F + \sigma$$

at all that we have seen that a plenty of things happen:

- $m_\sigma \neq 0$
- $m_\pi = 0$ (for $\epsilon=0$, otherwise $m_\pi^2 = \epsilon/F$)

•  Decay curve



is very small in the chiral limit

(and for $K_u \rightarrow 0$... soft limit!!!)

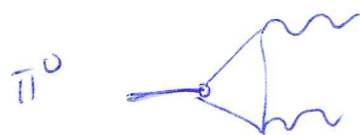
The "condensation" is visible in the interactions which generate, not in the value of the field itself... This is also valid for the σ and for the Higgs!!!

Second point :

$\pi \mapsto$ "very small known hadron" ... the lightest hadron.

$$\pi^\pm, \pi^0$$

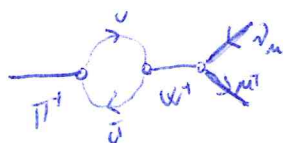
$$\pi^+ \mapsto \mu^+ \nu_\mu \rightarrow e^+ \nu_e$$



$\gamma \gamma$

$$\pi^0 F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$



Why longer than $e^+ \nu_e$?

(Mixing) R-L ... the mass term is responsible for that...

$$\Gamma \propto m_e^2$$

... but : where is the σ ? which is the particle corresponding to the σ ?

Search for the state with the corresponding quantum number...

the lightest one is the state $f_0(500)$! Is this the χ -partner of the pion? The

This state has a very long and difficult history.

fluctuation of the chiral condensate?

1S_0 $f_0(600)$ $M_0 \approx 400 - 1200$ MeV.

$^1S_0 \rightarrow$ removed

$^1S_0 \rightarrow$ returned!

$^1S_0 \rightarrow$ removed $f_0(500)$ and
man!

"450 - 1250" has quoted

(some

$\Rightarrow$ why the pole?

$\Rightarrow$ very broad state ...

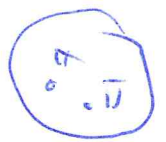
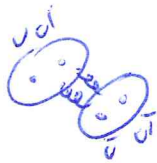
"An unstable state does not have a properly defined mass, but has an energy distribution".

Now, although the existence is "certain", the question is if this is the state we need...

Indeed, many works are in agreement that $f_2(500)$ is not $\frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}) \dots$

The possibilities are:

- $[u, d][\bar{u}, \bar{d}]$
- $\pi\pi$ molecule (or dyn. generated enhancement)



= The next state to look for is

$f_2(980)$

but this state couples very strongly to $K \dots$ it must contain a strange quark!

Indeed, this state can be also well interpreted as a tq :

$$\frac{1}{\sqrt{2}}([u, s][\bar{u}, \bar{s}] + [d, s][\bar{d}, \bar{s}])$$

$\Rightarrow \phi_2(500), f_2(980), \omega_2(980), K_0^*(800) \equiv$ NUMBER
OF TQ
STATES...

... then, the next available state is the resonance $f_0(1370)$

It couples strongly to $\pi\pi$ (and less strongly to $K\bar{K}$).

it is broad...

it is "heavy"... $M_{f_0} \sim M_{f_1} \sim M_{\omega_1} \sim M_{f_2} \sim M_{\omega_2}$

Namely, recall that a "scalar" $J^{PC} = 0^{++}$ arises from

$L = 1 = S \rightarrow$ similar mass than the other p-wave states.

$$M_{f_0} > M_{\omega}$$

$$L=1 \\ S=1$$

$$L=0 \\ S=1$$

What does the σ model tell us? Interesting point.

LSM

$$f_0(500) \rightarrow \pi\pi$$

$$f_0 \approx 400 \text{ MeV}$$

correct...

this is indeed one of the reasons why $f_0(500)$ was long-time considered as such...

eLSM
(with VM)

$$f_0(500) \rightarrow \pi\pi \approx 150 \text{ MeV} \Rightarrow \text{too narrow!!!}$$

This has to do with technical details...

but the result is obviously

VMore here, cannot be omitted.

→

... but then:

$$f_0(1370) \begin{cases} \xrightarrow{\text{LSM (no VM)}} 2-3 \text{ GeV} \\ \xrightarrow{\text{eLSM}} \end{cases}$$

$\approx 500 \text{ MeV}$ OK!!!

Indeed, we also find that $f_0(1370)$ is $\frac{1}{\sqrt{2}}(\bar{u}u + \bar{d}d)$ and is the π partner of the pion. Then, $f_0(1500) \approx \pi_1$, $f_0(1710) \approx \eta_8$.

PDG comes to the same conclusion.

Here, care is needed:

- lattice $\rightarrow$ huge problems in the scalar sector.
- various models and calculations... very plausible solution...
but not yet confirmed!
- Mixing:

