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# Elliptic emission of $K^+$ and $\pi^+$ in 158 A · GeV Pb + Pb collisions

WA98 Collaboration

M.M. Aggarwal<sup>a</sup>, A. Agnihotri<sup>b</sup>, Z. Ahammed<sup>c</sup>, A.L.S. Angelis<sup>d</sup>, V. Antonenko<sup>e</sup>,  
V. Arefiev<sup>f</sup>, V. Astakhov<sup>f</sup>, V. Avdeitchikov<sup>f</sup>, T.C. Awes<sup>g</sup>, P.V.K.S. Baba<sup>h</sup>,  
S.K. Badyal<sup>h</sup>, A. Baldine<sup>f</sup>, L. Barabach<sup>f</sup>, C. Barlag<sup>i</sup>, S. Bathe<sup>i</sup>, B. Batiounia<sup>f</sup>,  
T. Bernier<sup>j</sup>, K.B. Bhalla<sup>b</sup>, V.S. Bhatia<sup>a</sup>, C. Blume<sup>i</sup>, R. Bock<sup>k</sup>, E.-M. Bohne<sup>i</sup>,  
Z. Böröcz<sup>i</sup>, D. Bucher<sup>i</sup>, A. Buijs<sup>l</sup>, H. Büsching<sup>i</sup>, L. Carlen<sup>m</sup>, V. Chalyshev<sup>f</sup>,  
S. Chattopadhyay<sup>c</sup>, R. Cherbatchev<sup>e</sup>, T. Chujo<sup>n</sup>, A. Claussen<sup>i</sup>, A.C. Das<sup>c</sup>,  
M.P. Decowski<sup>r</sup>, H. Delagrange<sup>j</sup>, V. Djordjadze<sup>f</sup>, P. Donni<sup>d</sup>, I. Doubovik<sup>e</sup>,  
S. Dutt<sup>h</sup>, M.R. Dutta Majumdar<sup>c</sup>, K. El Chenawi<sup>m</sup>, S. Eliseev<sup>o</sup>, K. Enosawa<sup>n</sup>,  
P. Foka<sup>d</sup>, S. Fokin<sup>e</sup>, V. Frolov<sup>f</sup>, M.S. Ganti<sup>c</sup>, S. Garpman<sup>m</sup>, O. Gavrishchuk<sup>f</sup>,  
F.J.M. Geurts<sup>l</sup>, T.K. Ghosh<sup>p</sup>, R. Glasow<sup>i</sup>, S.K. Gupta<sup>b</sup>, B. Guskov<sup>f</sup>,  
H.Å. Gustafsson<sup>m</sup>, H.H. Gutbrod<sup>j</sup>, R. Higuchi<sup>n</sup>, I. Hrivnacova<sup>o</sup>, M. Ippolitov<sup>e</sup>,  
H. Kalechofsky<sup>d</sup>, R. Kamermans<sup>l</sup>, K.-H. Kampert<sup>i</sup>, K. Karadjev<sup>e</sup>, K. Karpio<sup>q</sup>,  
S. Kato<sup>n</sup>, S. Kees<sup>i</sup>, B.W. Kolb<sup>k</sup>, I. Kosarev<sup>f</sup>, I. Koutcheryaev<sup>e</sup>, T. Krümpel<sup>i</sup>,  
A. Kugler<sup>o</sup>, P. Kulinich<sup>r</sup>, M. Kurata<sup>n</sup>, K. Kurita<sup>n</sup>, N. Kuzmin<sup>f</sup>, I. Langbein<sup>k</sup>,  
A. Lebedev<sup>e</sup>, Y.Y. Lee<sup>k</sup>, H. Löhner<sup>p</sup>, L. Luquin<sup>j</sup>, D.P. Mahapatra<sup>s</sup>, V. Manko<sup>e</sup>,  
M. Martin<sup>d</sup>, G. Martínez<sup>j</sup>, A. Maximov<sup>f</sup>, R. Mehdiyev<sup>f</sup>, G. Mgebrichvili<sup>e</sup>,  
Y. Miake<sup>n</sup>, D. Mikhalev<sup>f</sup>, Md.F. Mir<sup>h</sup>, G.C. Mishra<sup>s</sup>, Y. Miyamoto<sup>n</sup>,  
D. Morrison<sup>t</sup>, B. Mohanty<sup>s</sup>, D.S. Mukhopadhyay<sup>c</sup>, V. Myalkovski<sup>f</sup>, H. Naef<sup>d</sup>,  
B.K. Nandi<sup>s</sup>, S.K. Nayak<sup>j</sup>, T.K. Nayak<sup>c</sup>, S. Neumaier<sup>k</sup>, A. Nianine<sup>e</sup>,  
V. Nikitine<sup>f</sup>, S. Nikolaev<sup>e</sup>, P. Nilsson<sup>m</sup>, S. Nishimura<sup>n</sup>, P. Nomokonov<sup>f</sup>,  
J. Nystrand<sup>m</sup>, F.E. Obenshain<sup>t</sup>, A. Oskarsson<sup>m</sup>, I. Otterlund<sup>m</sup>, M. Pachr<sup>o</sup>,  
A. Parfenov<sup>f</sup>, S. Pavliouk<sup>f</sup>, T. Peitzmann<sup>i</sup>, V. Petracek<sup>o</sup>, F. Plasil<sup>g</sup>,  
W. Pinganaud<sup>j</sup>, M.L. Porschke<sup>k</sup>, B. Raeven<sup>l</sup>, J. Rak<sup>o</sup>, R. Raniwala<sup>b</sup>,  
S. Raniwala<sup>b</sup>, V.S. Ramamurthy<sup>s</sup>, N.K. Rao<sup>h</sup>, F. Retiere<sup>j</sup>, K. Reygers<sup>i</sup>,  
G. Roland<sup>r</sup>, L. Rosselet<sup>d</sup>, I. Roufanov<sup>f</sup>, C. Roy<sup>j</sup>, J.M. Rubio<sup>d</sup>, H. Sako<sup>n</sup>,  
S.S. Sambyal<sup>h</sup>, R. Santo<sup>i</sup>, S. Sato<sup>n</sup>, H. Schlagheck<sup>i</sup>, H.-R. Schmidt<sup>k</sup>, Y. Schutz<sup>j</sup>,  
G. Shabrato<sup>f</sup>, T.H. Shah<sup>h</sup>, I. Sibiriak<sup>e</sup>, T. Siemiarczuk<sup>q</sup>, D. Silvermyr<sup>m</sup>,

B.C. Sinha <sup>c</sup>, N. Slavine <sup>f</sup>, K. Söderström <sup>m</sup>, N. Solomey <sup>d</sup>, S.P. Sørensen <sup>t</sup>,  
 P. Stankus <sup>g</sup>, G. Stefanek <sup>q</sup>, P. Steinberg <sup>r</sup>, E. Stenlund <sup>m</sup>, D. Stüken <sup>i</sup>, M. Sumbera <sup>o</sup>,  
 T. Svensson <sup>m</sup>, M.D. Trivedi <sup>c</sup>, A. Tsvetkov <sup>e</sup>, L. Tykarski <sup>q</sup>, J. Urbahn <sup>k</sup>,  
 E.C. v.d. Pijll <sup>l</sup>, N. v. Eijndhoven <sup>l</sup>, G.J. v. Nieuwenhuizen <sup>r</sup>, A. Vinogradov <sup>e</sup>,  
 Y.P. Vijoyi <sup>c</sup>, A. Vodopianov <sup>f</sup>, S. Vörös <sup>d</sup>, B. Wysłouch <sup>r</sup>, K. Yagi <sup>n</sup>, Y. Yokota <sup>n</sup>,  
 G.R. Young <sup>g</sup>

<sup>a</sup> University of Panjab, Chandigarh 160014, India

<sup>b</sup> University of Rajasthan, Jaipur 302004, Rajasthan, India

<sup>c</sup> Variable Energy Cyclotron Centre, Calcutta 700 064, India

<sup>d</sup> University of Geneva, CH-1211 Geneva 4, Switzerland

<sup>e</sup> RRC Kurchatov Institute, RU-123182 Moscow, Russia

<sup>f</sup> Joint Institute for Nuclear Research, RU-141980 Dubna, Russia

<sup>g</sup> Oak Ridge National Laboratory, Oak Ridge, TN 37831-6372, USA

<sup>h</sup> University of Jammu, Jammu 180001, India

<sup>i</sup> University of Münster, D-48149 Münster, Germany

<sup>j</sup> SUBATECH, Ecole des Mines, Nantes, France

<sup>k</sup> Gesellschaft für Schwerionenforschung (GSI), D-64220 Darmstadt, Germany

<sup>l</sup> Universiteit Utrecht / NIKHEF, NL-3508 TA Utrecht, The Netherlands

<sup>m</sup> Lund University, SE-221 00 Lund, Sweden

<sup>n</sup> University of Tsukuba, Ibaraki 305, Japan

<sup>o</sup> Nuclear Physics Institute, CZ-250 68 Rez, Czech Republic

<sup>p</sup> KVI, University of Groningen, NL-9747 AA Groningen, The Netherlands

<sup>q</sup> Institute for Nuclear Studies, 00-681 Warsaw, Poland

<sup>r</sup> MIT, Cambridge, MA 02139, USA

<sup>s</sup> Institute of Physics, 751-005 Bhubaneswar, India

<sup>t</sup> University of Tennessee, Knoxville, TN 37966, USA

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## Abstract

An event-by-event analysis of the azimuthal angular correlation with respect to the reaction plane has been carried out for  $K^+$  and  $\pi^+$  emission near mid-rapidity in 158 A · GeV Pb + Pb collisions. In semi-central collisions,  $K^+$  mesons are found to be preferentially emitted out of the reaction plane, while  $\pi^+$  mesons are emitted in the reaction plane. The results suggest that the kaon emission is influenced by in-medium potential effects in addition to collective flow effects. © 1999 Published by Elsevier Science B.V. All rights reserved.

## 1. Introduction

Relativistic heavy ion collisions provide a unique tool for the study of nuclear properties at high temperature and density. The study of collective motion in the final state of the produced hadrons is expected to provide information on both the dynamics of heavy-ion collisions and the equation of state [1]. The observation of collective behavior in the system provides a validation of a hydrodynamical description of the dynamics. Once a hydrodynamical

interpretation is established the collective motion will result from pressure gradients in the matter, which will reflect the compressibility of the underlying equation of state. In the case of a phase transition from ordinary matter to quark-gluon plasma, it is expected that the compressibility should exhibit a softening due to the increased number of degrees of freedom [2].

A standard way to search for collective motion in the final state is to analyze the azimuthal distribution of particle emission with respect to the reaction

plane in terms of a Fourier expansion. The first two coefficients in the Fourier decomposition are termed the directed and elliptic flow components. The importance of collective motion measurements has been emphasized by several authors [2–5] and measurements have been reported over a range of energies; directed and elliptic flow at  $0.1 \sim 1 A \cdot \text{GeV}$  [6–9], at  $10 A \cdot \text{GeV}$  [10] and also at  $158 A \cdot \text{GeV}$  [11,12]. One of the striking features of these results is the change of the elliptic flow direction from out-of-plane (squeeze-out) to in-plane in the region between  $1 A \cdot \text{GeV}$  and  $100 A \cdot \text{GeV}$  incident energy.

The azimuthal distribution of a particular particle species in the final state is also expected to be sensitive to the in-medium potential of the particle. Namely, it is interesting to study the distributions of  $K^+$  and  $K^-$  mesons because the kaon-nucleon potential in nuclear matter is expected to be repulsive for  $K^+$  but attractive for  $K^-$  mesons [13–15].

Recently, the elliptic emission of nucleons, pions, and  $K^+$  mesons has been observed in the mid-rapidity region in  $0.8 \sim 1 A \cdot \text{GeV}$   $^{197}\text{Au} + ^{197}\text{Au}$  and  $^{209}\text{Bi} + ^{209}\text{Bi}$  collisions [7–9]. The observed out-of-plane emission of  $K^+$  mesons is claimed as a consequence of the repulsive  $K^+$ -nucleon potential. On the other hand, pions and nucleons also show out-of-plane emission at this energy, which is interpreted as dominantly due to the shadowing effect of the spectator nucleons [7,16,17].

At higher incident energies, particles produced near mid-rapidity are not expected to interact with the spectator matter due to the short crossing time of the collisions. Therefore, the in-medium potential effects might be more clearly visible at SPS energies [18]. In this letter we present the azimuthal distributions, relative to the reaction plane, of  $K^+$  and  $\pi^+$  mesons near mid-rapidity in  $158 A \cdot \text{GeV}$   $\text{Pb} + \text{Pb}$  collisions. This is the first observation of the elliptic emission of  $K^+$  mesons at SPS energies.

## 2. Experimental setup

The data presented here were taken with a subset of the WA98 experiment detector system using the  $158 A \cdot \text{GeV}$   $^{208}\text{Pb}$  beams of the CERN-SPS on a  $\text{Pb}$  target of  $213 \mu\text{m}$  thickness. In the WA98 setup the incident  $\text{Pb}$  beam provides a trigger by a valid signal in a gas Cherenkov start counter with a timing

resolution of 30 ps and no signal in a downstream veto counter with a 3 mm diameter hole. The minimum bias trigger and centrality of the collision is determined by the total transverse energy,  $E_T$ , measured with the mid-rapidity calorimeter (MIRAC) which covers the pseudo-rapidity range of  $3.5 < \eta < 5.5$ . The Plastic Ball detector has full azimuthal coverage in the pseudo-rapidity range of  $-1.7 < \eta < 0.5$ . It identifies pions, protons, deuterons, and tritons with kinetic energies of 50 to 250 MeV by simultaneous  $\Delta E$  and  $E$  measurement using the mass and energy dependence of  $dE/dx$  [11]. The measurement of identified particles near mid-rapidity is performed using two tracking spectrometer arms with a large (1.6 m) aperture dipole magnet (GOLIATH) which provides 1.6 Tm bending power. The particle identification is based on a measurement of momentum and time-of-flight.

The present analysis has been performed on the full sample of WA98 positive charged particle tracking arm data. This data sample is limited due to the fact that the positive charged particle tracking arm was installed only for the final run period of WA98. The tracking arm has a momentum resolution of  $\Delta p/p \approx 0.97\% + 0.16\% p + 0.023\% p^2$  ( $p$  in  $\text{GeV}/c$ ) and a time-of-flight resolution of  $< 90$  ps. The rapidity and transverse momentum ( $y, p_T$ ) coverages of the Plastic Ball and the tracking arm are shown in Fig. 1. The ( $y, p_T$ ) coverages for the

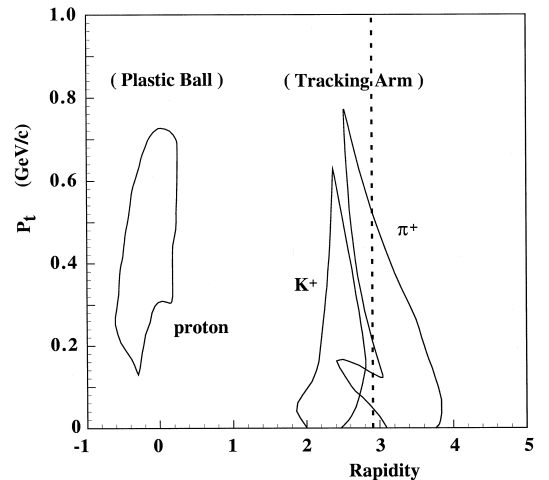


Fig. 1. The rapidity and transverse momentum ( $y, p_T$ ) coverages of the Plastic Ball detector and the tracking arm. The midrapidity of the collisions is shown as dashed line.

Table 1

Results of the Fourier analysis of  $K^+$  and  $\pi^+$  mesons for semi-central 158 A·GeV Pb + Pb collisions. The values are integrated over the indicated  $y$  and  $p_T$  ranges. Positive  $v_1$  corresponds to an emission in the same direction as nucleons. Observed negative value of  $v_1$  for  $\pi^+$  indicates so-called anti-flow of pions. Positive and negative  $v_2$  corresponds to an in-plane and out-of-plane elliptic emission, respectively. The values are obtained numerically by evaluating  $\langle \cos(\phi - \Phi_0) \rangle$  and  $\langle \cos(2(\phi - \Phi_0)) \rangle$ , with corrections of the reaction plane resolution at each  $E_T$  bin. The errors include the statistical errors and the errors of the experimental resolution of the reaction plane determination

Particle	$y$	$p_T$ (GeV/c)	$v_1$	$v_2$
$K^+$	2.2–2.8	0.05–0.60	$-0.004 \pm 0.031$	$-0.24 \pm 0.14$
$\pi^+$	2.4–3.4	0.05–0.80	$-0.010 \pm 0.006$	$0.047 \pm 0.024$

present analysis are listed in Table 1. Detailed information about the positive tracking arm [20] and the experiment can be found elsewhere [11,19].

### 3. Experimental results

The reaction plane is determined from the azimuthal direction of the total transverse momentum vector of fragments ( $p$ ,  $d$ , and  $t$ ) observed by the Plastic Ball detector [11]. The azimuthal angle of the reaction plane,  $\Phi_0$  in the laboratory is determined as

$$\Phi_0 = \tan^{-1} \left( \frac{\sum_{i=1}^N p_{Ti} \sin(\phi_i)}{\sum_{i=1}^N p_{Ti} \cos(\phi_i)} \right), \quad (1)$$

where the sum runs over all fragments. Here  $\phi_i$  and  $p_{Ti}$  are the azimuthal angle in the laboratory and the transverse momentum of the  $i$ -th fragment, respectively. The multiplicity of protons in the Plastic Ball detector is around 8 in semi-central collisions and a minimum of three protons is required for this analysis. The  $\Phi_0$ 's accumulated for many events should be uniformly distributed if there is no detector bias. The observed  $\Phi_0$  distribution has a variation of less than 2% due to the detector biases such as dead channels and inefficiency. In the following analysis, we have corrected for this effect by weighting with the inverse of the yield although there is no significant difference in the results obtained with or without this correction.

The angular correlation between the azimuthal angle  $\phi$  of  $K^+$  and  $\pi^+$  near mid-rapidity and the  $\Phi_0$  determined in the target rapidity region has been studied. Since the acceptance of  $\pi^+$  crosses mid-rapidity,  $\Phi_0$  was rotated by  $180^\circ$  for particles in the region forward of mid-rapidity. Otherwise, the extracted value of the directed flow would be underestimated since it changes sign at mid-rapidity. Fig. 2 shows the azimuthal distributions of  $K^+$  and  $\pi^+$  mesons with respect to the reaction plane for semi-central ( $50 < E_T < 250$  GeV in MIRAC) and central ( $320 < E_T < 500$  GeV) Pb + Pb collisions. For  $\pi^+$  mesons in semi-central collisions, the azimuthal distribution indicates weak maxima at  $\phi - \Phi_0 = 0^\circ$  and  $\pm 180^\circ$ , which indicates an enhanced emission in the reaction plane. On the other hand, the  $K^+$  azimuthal distribution in semi-central collisions exhibits maxima at  $\phi - \Phi_0 = \pm 90^\circ$  which demonstrates an enhanced emission out of the reaction plane. Although the statistics is limited, the results provide evidence that the  $K^+$  emission axis is orthogonal to the in-plane emission axis of the  $\pi^+$ . In central collisions, as shown in Fig. 2, the azimuthal distributions for both particle types are nearly flat, as expected from the near azimuthal symmetry of the collision system at small impact parameter.

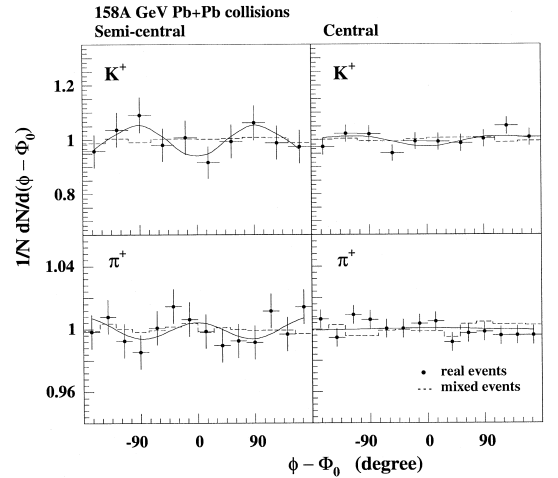


Fig. 2. The azimuthal distributions of  $K^+$  and  $\pi^+$  mesons with respect to the reaction plane for semi-central and central 158 A·GeV Pb + Pb collisions. Solid circles with error bars show the real events and dashed histograms show the results of the mixed events. The solid curves show the fits using Eq. (2). The  $(y, p_T)$  covers for this analysis are listed in Table 1.

To test for detector effects, such as geometrical acceptance and detector efficiency, fake events were created using a mixed-event method which were then analyzed in the same manner as the real events. The mixed-event results are shown as histograms in Fig. 2. The azimuthal distributions for the mixed-events are flat which indicates that the observed anisotropies are not due to detector effects.

The strength of the azimuthal anisotropy is evaluated by means of a Fourier expansion [21,22]. The Fourier coefficients  $v_n$  ( $n=1,2$ ) can be extracted from the azimuthal distribution of identified particles with respect to  $\Phi_0$ ;

$$\frac{1}{N} \frac{dN}{d(\phi - \Phi_0)} = 1 + 2v'_1 \cos(\phi - \Phi_0) + 2v'_2 \cos(2(\phi - \Phi_0)), \quad (2)$$

where  $\phi$  is the azimuthal angle of the particle in the laboratory frame. The Fourier coefficient  $v'_1$  quantifies the directed flow, whereas  $v'_2$  quantifies the elliptic flow. Positive  $v'_1$  corresponds to an emission in the same direction as nucleons. (Please note that this definition is opposite to our previous reports [11].) The coefficient  $v'_2$  is negative for an out-of-plane emission and positive for an in-plane emission.

The observed distributions are affected by the experimental resolution of the reaction plane determination. Since the measured Fourier coefficients are reduced by this effect, they should be corrected as [21–23]

$$v_n = \frac{v'_n}{\langle \cos(n\Delta\Phi) \rangle}, \quad (n=1,2), \quad (3)$$

where  $\Delta\Phi$  is the deviation from the true reaction plane and  $\langle \rangle$  indicates the mean value summed over all events. The reaction plane resolution is determined by randomly subdividing the full event into two equal size subevents and extracting  $\langle \cos(\Delta\Phi) \rangle \approx \sqrt{2} \cdot \sqrt{\langle \cos(\Phi_a - \Phi_b) \rangle}$ , where  $\Phi_a$  and  $\Phi_b$  are the azimuthal angles of the reaction plane of two equal size subevents. Using the more accurate interpolation formula of Ref. [22] one obtains  $\langle \cos(\Delta\Phi) \rangle = 0.360 \pm 0.018$ ,  $\langle \cos(2\Delta\Phi) \rangle = 0.084 \pm 0.005$  for semi-central collisions.

The values of the Fourier coefficients  $v_n$  ( $n=1,2$ ) are obtained numerically by evaluating  $\langle \cos(\phi - \Phi_0) \rangle$  and  $\langle \cos(2(\phi - \Phi_0)) \rangle$ , respectively. The re-

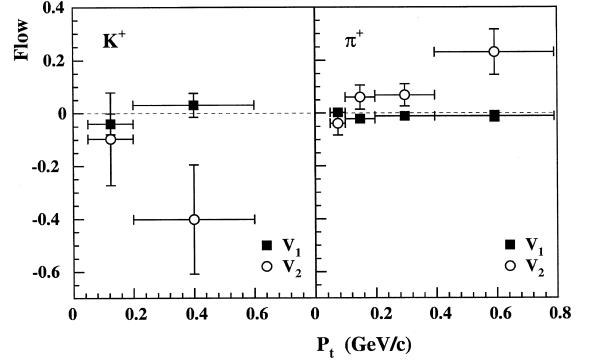


Fig. 3. The  $p_T$  dependence of the  $v_1$  and  $v_2$  values of  $K^+$  and  $\pi^+$  mesons for semi-central 158 A-GeV Pb + Pb collisions. The rapidity coverages are listed in Table 1. The results are corrected for the experimental resolution of the reaction plane determination. Solid squares are used for  $v_1$  and open circles for  $v_2$ .

sults for semi-central collisions are listed in Table 1<sup>1</sup>. The values have been corrected for the reaction plane resolution according to Eq. (3) at each  $E_T$  bin. The value of  $v_1$ , indicating the strength of the directed flow, is negative for  $\pi^+$ , which corresponds to the direction opposite to the nucleons. This so called anti-flow of pions has been previously reported [11,12]. The value of  $v_1$  for the  $K^+$  meson is compatible with zero. The value of  $v_2$  for the  $K^+$  meson is negative with a  $1.7\sigma$  separation from zero, while the value of  $v_2$  for the  $\pi^+$  mesons is positive with a  $1.9\sigma$  separation from zero. Under the assumption of gaussian errors these results correspond to confidence levels of 95.7% and 97.5% for the assignment of the sign of  $v_2$  for  $K^+$  and  $\pi^+$ , respectively. Although the significance of the result is limited by statistics, one may conclude at the 93% confidence level that the elliptic emission planes for  $K^+$  and  $\pi^+$  are oriented orthogonally to one another. The systematic errors on the  $v_2$  values due to the contamination from other particle species are estimated to be less than 8% for  $K^+$  and 2% for  $\pi^+$ .

Fig. 3 shows the  $v_1$  and  $v_2$  values as a function of the transverse momentum  $p_T$  for  $K^+$  and  $\pi^+$  mesons in semi-central collisions. For the  $K^+$  data

<sup>1</sup> One can obtain the values of  $v_1$  and  $v_2$  from fitting the azimuthal distribution with Eq. (2). Within the errors, the fit values agree with those listed in Table 1.

only two bins are shown due to the limited statistics. The magnitude of  $v_2$  for  $K^+$  and  $\pi^+$  mesons tends to increase with  $p_T$ , while the  $v_1$  values remain close to zero for both particles.

#### 4. Discussion

To compare the present results with other measurements the  $v_2$  values for protons, pions, and kaons near mid-rapidity are plotted as a function of the beam energy in Fig. 4. For both protons and pions a transition from out-of-plane to in-plane emission occurs around 5–10  $A \cdot \text{GeV}$ . At SPS energy, results from NA49 [12] have shown that protons and pions exhibit in-plane emission near mid-rapidity. Our  $\pi^+$  data agree with the NA49 results within errors. Unlike the case for protons and pions, the present results indicates that  $K^+$  mesons exhibit out-of-plane emission, which is similar to observations for 1  $A \cdot \text{GeV}$  Au + Au collisions [8].

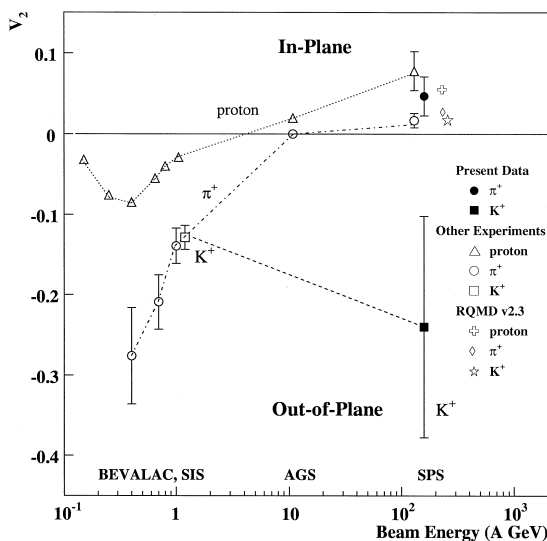


Fig. 4. Beam energy dependence of the  $v_2$  value near mid-rapidity [6–8,10,12]. Solid symbols indicate the present data. Note that the region in  $(y, p_T)$  is different for each experiment. The RQMD (v2.3 cascade mode) calculations for proton,  $\pi^+$  and  $K^+$  in 158  $A \cdot \text{GeV}$  Pb + Pb collisions are also shown for the impact parameter range  $b = 6.5\text{--}12$  fm with the filter of experimental acceptance. The data points are shifted horizontally where they overlap.

Results of RQMD model calculations (v2.3 cascade mode) [24] filtered with the experimental acceptance are also shown in Fig. 4. The RQMD results show in-plane emission for  $\pi^+$ , proton, and also for  $K^+$ . The RQMD calculation agrees with the measured results for  $\pi^+$  and proton, but it fails to reproduce the out-of-plane elliptic emission of  $K^+$ . These results might suggest that a new ingredient to the calculation such as the in-medium potentials of the  $K^+$  and  $K^-$  mesons is required [13–15]. A simple model calculation [25] in which  $K^+$  mesons propagate through a static anisotropic distribution of nucleons with a repulsive  $K^+N$  potential demonstrates that a final out-of-plane elliptic emission pattern can emerge from an initially isotropic azimuthal distribution of  $K^+$ .

It is clear that if the azimuthal distribution of the pions reflects their preferential in-plane collective motion, then the out-of-plane  $K^+$  azimuthal distribution would indicate that the kaons carry little direct information about the collective motion of the pion matter. The results therefore appear to question the assumption of a common collective flow velocity for pions and kaons.

#### 5. Summary

In summary, we have measured the elliptic emission patterns of  $K^+$  and  $\pi^+$  mesons near mid-rapidity in 158  $A \cdot \text{GeV}$  Pb + Pb collisions. In semi-central collisions, the  $K^+$  mesons are found to be emitted preferentially perpendicular to the reaction plane, while  $\pi^+$  mesons tend to be emitted in the reaction plane. The RQMD cascade calculation reproduces the  $v_2$  values for  $\pi^+$  and protons, but it fails to explain the out-of-plane emission of  $K^+$  mesons. The results suggest that the kaon emission is influenced by in-medium potential effects in addition to collective flow effects.

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