

Forward-backward multiplicity correlations in a superposition approach

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Unreasonable effectiveness of statistical approaches to high-energy collisions,
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[based on AO & W. Broniowski, PRC 88 (2013) 044913, arXiv:1303.5280v2]

F-B correlations with wide rapidity separation provide information on the earliest stages of the collision

- [1] B. Back et al. (PHOBOS Collaboration), Phys. Rev. C 74,011901 (2006)
- [2] T. J. Tarnowsky, Ph.D thesis, Purdue University, 2008
- [3] B. Abelev et al. (STAR Collaboration), Phys. Rev. Lett. 103, 172301 (2009)
- [4] P. Brogueira and J. Dias de Deus, Phys. Lett. B 653, 202 (2007)
- [5] T. Lappi and L. McLerran, Nucl. Phys. A 832, 330 (2010)
- [6] A. Bzdak, Phys. Rev. C 85, 051901 (2012)
- [7] A. Bialas, J. Phys. G 35, 044053 (2008)
- [8] A. Bialas and K. Zalewski, Phys. Rev. C 82, 034911 (2010)
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- F-B multiplicity correlation measured by STAR Cu+Cu collaboration for Au+Au at $\sqrt{s_{NN}} = 200$ GeV, at RHIC [1,2,3]

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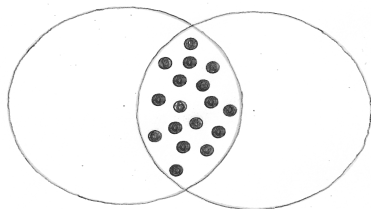
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- Three stage superposition model

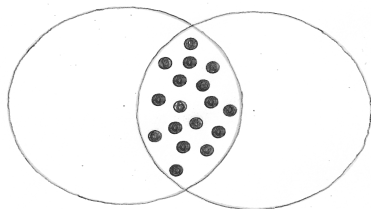
- Three stage superposition model
- Comparison to LHC
 - 1 Glauber model analysis
 - 2 Model independent analysis

- Two colliding nuclei in the transverse plane

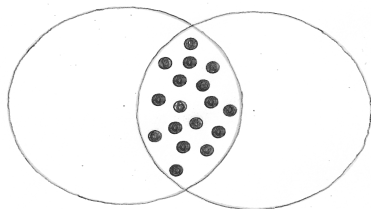


sources in
the transverse plane

- Two colliding nuclei in the transverse plane
- Nucleons create **sources**

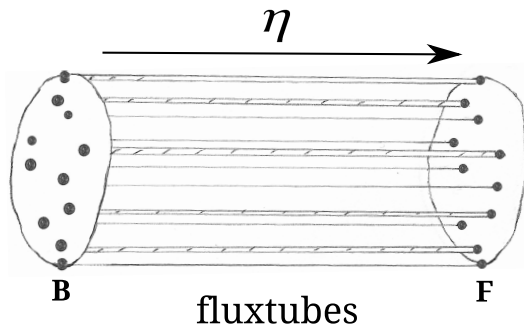


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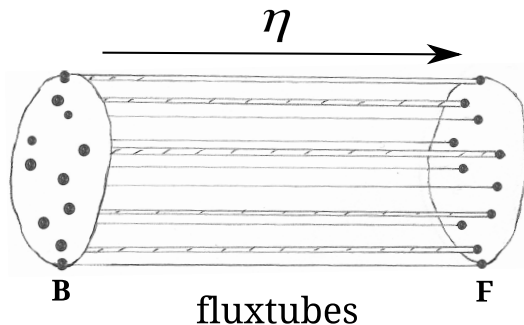
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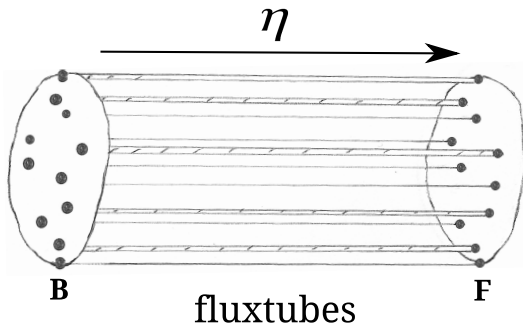
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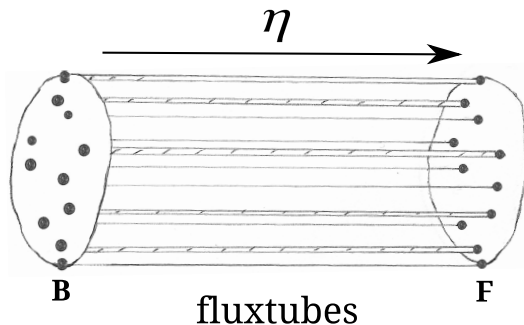
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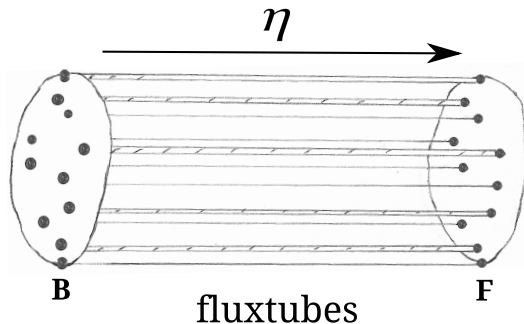


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A possible simple 3D picture:

- The concept of the longitudinal strings (fluxtubes)
- Number of fluxtubes increases with centrality
- We assume well separated F and B rapidity bins
- η corresponds to spatial rapidity

$$\eta = \frac{1}{2} \ln \left(\frac{|\mathbf{p}| + p_L}{|\mathbf{p}| - p_L} \right)$$



Hypothesis of maximum F-B correlations

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Superposition model

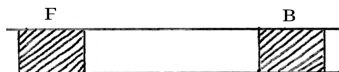
"Statistical evolution" of the fireball

Many successful models are based on the three stage approach

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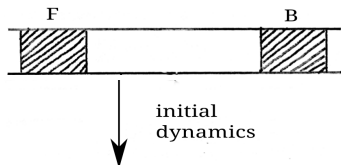
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s - number of initial sources

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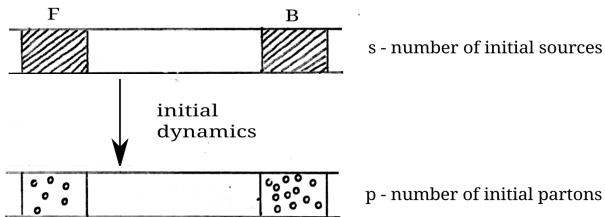


s - number of initial sources

- initial tubes extend along the f-b range
- partons formed from braking fluxtubes

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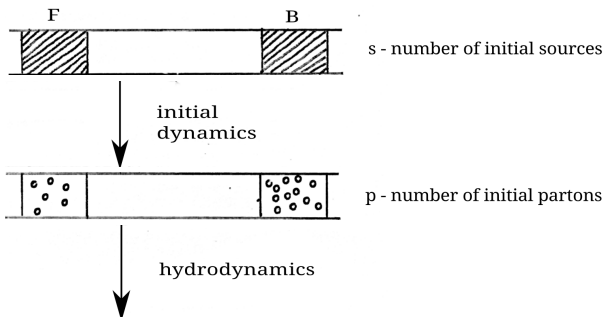
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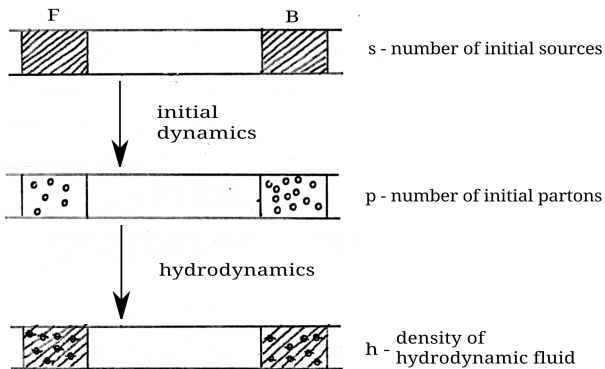
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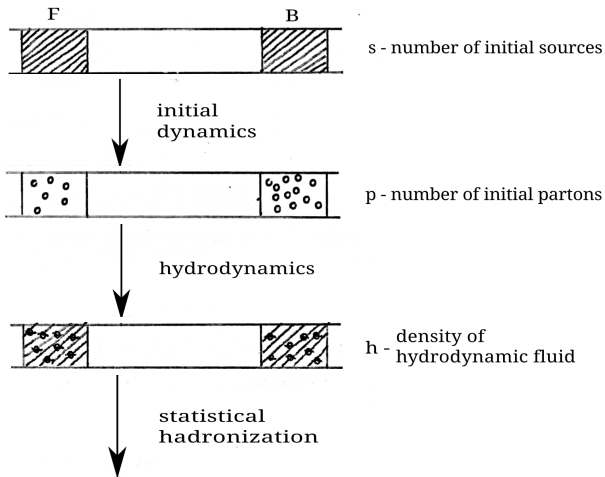
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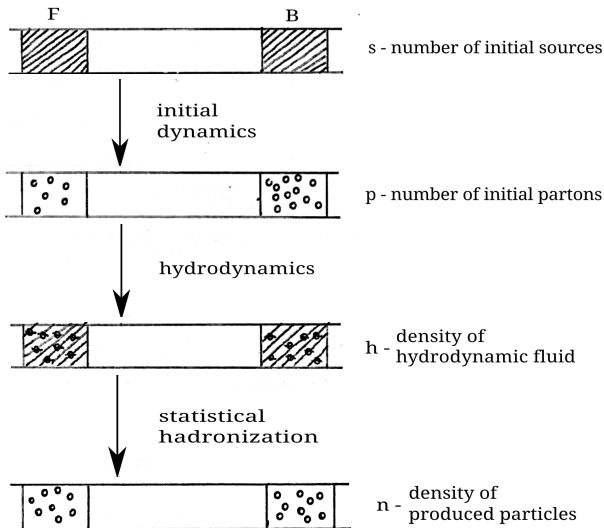
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- partons formed from braking fluxtubes
- partons with the same distribution
- deterministic evolution
- creation of the fluid sources
- hadronization
- overlapped random distribution - freezeout of the fluid

- The production occurs from each source in the same universal manner

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$$p_A = \sum_{i=1}^{s_A} \mu_i, \quad A = F, B$$

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- Distribution of μ is universal \rightarrow independence from cell location
- Using superposition model \rightarrow well known formulas

$$s \xrightarrow{\text{init. production}} p$$

$$\begin{aligned}\langle p_A \rangle &= \langle \mu \rangle \langle s_A \rangle \\ \text{var}(p_A) &= \text{var}(\mu) \langle s_A \rangle + \langle \mu \rangle^2 \text{var}(s_A) \\ \text{cov}(p_F, p_B) &= \langle \mu \rangle^2 \text{cov}(s_F, s_B)\end{aligned}$$

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- When fluctuations are not too large \rightarrow

$$h = t_0 \langle p \rangle + t_1 (p - \langle p \rangle) + \mathcal{O} \left((p - \langle p \rangle)^2 \right)$$

t_i 's depend on properties of hydrodynamics

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- **The hydrodynamics is complicated but deterministic**

$$p \xrightarrow{\text{hydro}} h$$

$$\begin{aligned}\langle h_A \rangle &= t_0 \langle p_A \rangle \\ \text{var}(h_A) &= t_1^2 \text{var}(p_A) \\ \text{cov}(h_F, h_B) &= t_1^2 \text{cov}(p_F, p_B)\end{aligned}$$

\leftarrow Formulas link statistical properties of initial partons and hydrodynamics sources

- Cell emits n hadrons into a region of phase space with some statistical distribution superimposed over h .

$$n_A = \sum_{i=1}^{h_A} m_i, \quad A = F, B$$

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$$h \xrightarrow{\text{hadronization}} n$$

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$$\text{var}(n_A) = \text{var}(m) \langle h_A \rangle + \langle m \rangle^2 \text{var}(h_A)$$

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Joining all stages

$$s \xrightarrow{\text{3 stage}} n$$

$$\langle n_A \rangle = \alpha \langle s_A \rangle$$

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- The importance of $\gamma \rightarrow$ occurs with variance and covariance

Comparison to LHC

- Wounded nucleon [1] with binary collision [2] → Mixed model [3]

[1] A. Bialas, M. Bleszynski and W. Czyz, Nucl. Phys. B111 (1976) 461

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- We obtain number of sources

$$s_A = \frac{1}{2} (1 - a) s_w + a s_{\text{bin}}$$

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- Mixing parameter $\rightarrow a = 11\%$
- Inelastic cross section $\rightarrow \sigma_{NN}^{\text{inel}} = 65\text{mb}$

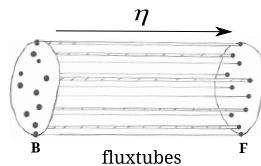
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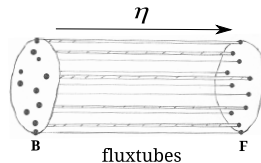
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Basic methodology:



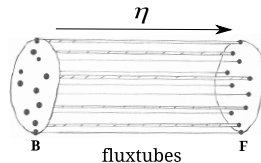
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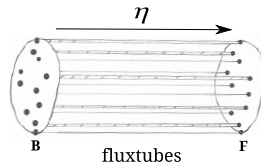
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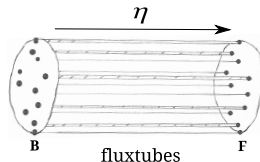
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- calculate theoretical values of statistical properties

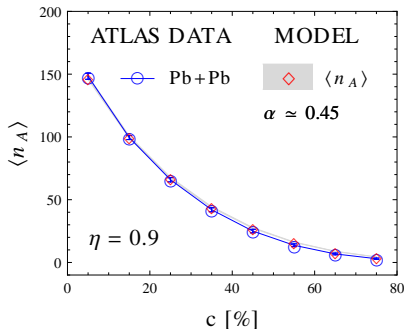
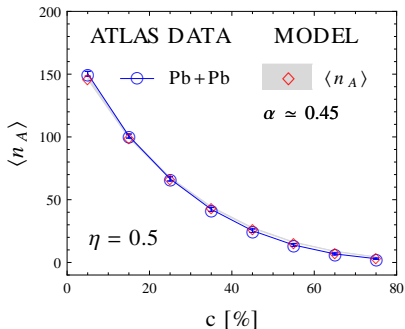


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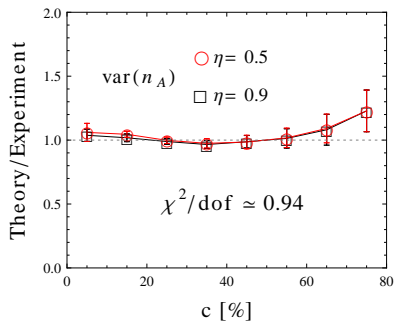
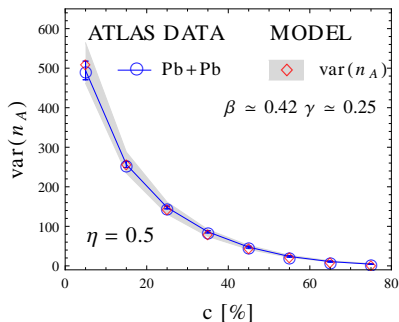


Mean



$$\langle n_A \rangle = \alpha \langle s_A \rangle$$

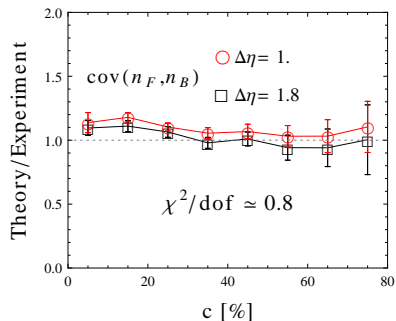
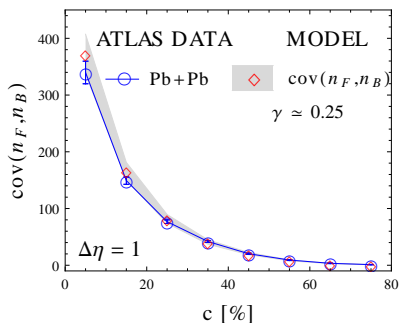
Variance



The α, β, γ practically independent of bin separation $\Delta\eta$

$$\text{var}(n_A) = \beta \langle s_A \rangle + \gamma \text{var}(s_A)$$

Covariance prediction

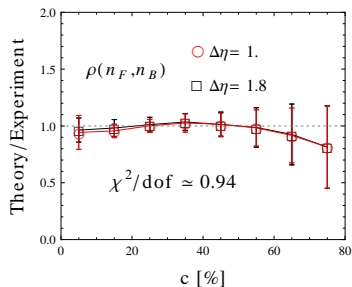
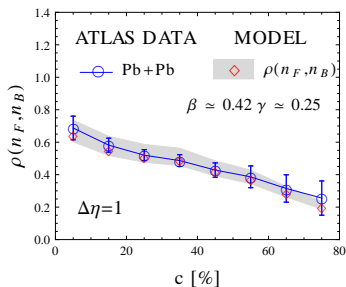


We don't fit anything, but use already fitted (test of consistency) γ

$$\text{cov}(n_F, n_B) = \gamma \text{cov}(s_F, s_B) = \rho(s_F, s_B) \gamma \text{var}(s_A)$$

$$\text{cov}(n_F, n_B) = \gamma \text{var}(s_A) \text{ (maximum correlated sources)}$$

Correlation



$$\rho(n_F, n_B) = \frac{\text{cov}(n_F, n_B)}{\text{var}(n_A)} \quad \omega(s_A) = \frac{\text{var}(s_A)}{\langle s_A \rangle}$$

$$\rho(n_F, n_B) = \frac{\rho(s_F, s_B)}{1 + \beta/\gamma\omega(s_A)}$$

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Consider parameters from model

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$$\frac{\alpha^2}{\gamma} = \frac{t_0}{t_1} \implies t_0 \simeq 0.9 t_1$$

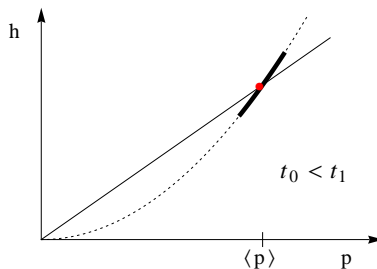
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Verification of two possibilities



- Hydrodynamic growth is faster than linear function

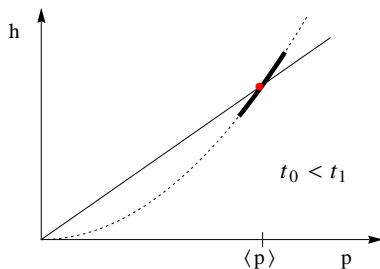
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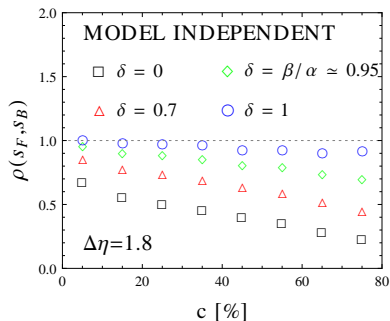
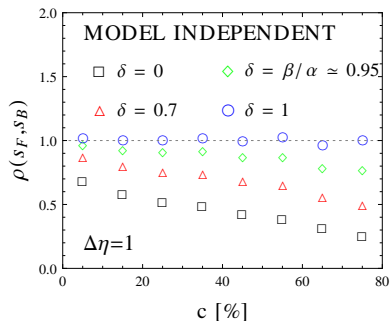
$$\frac{\alpha^2}{\gamma} = \frac{t_0}{t_1} \implies t_0 \simeq 0.9 t_1$$

Verification of two possibilities



- Hydrodynamic growth is faster than linear function
- Nonlinearity of hydrodynamics

Correlations prediction



- Formula use only measured quantities (no Glauber model)
- Using only one free parameter \rightarrow maximum correlated sources

$$\delta = \beta/\alpha \text{ (model parameter), } \rho(S_F, S_B) = \frac{\rho(n_F, n_B)}{1 - \delta/\omega(n_A)}$$

Conclusions

- Simple formulas linking the statistical properties of the F-B correlations in the data and in the original sources have been derived in the three-stage model. Together with the Glauber model for the sources leads to natural description of the early LHC data

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- The effect of hydrodynamics may be, under reasonable assumptions, incorporated in terms of just two parameters. Our study shows that the hydrodynamic growth faster than linear function
- The hypothesis of maximal correlation of sources (continuous fluxtubes), $\rho(s_F, s_B) = 1$, works for LHC.

Thank YOU!

Backup slides



- STAR measurement is **affected** by correlations to the *reference bin* n_R

STAR measurement

$$\langle n_A \rangle_{n_R} = c_0 + c_1 n_R$$

$$\rho^*(n_F, n_B) = \frac{\rho(n_F, n_B) - R^2}{1 - R^2}$$

$$\omega^*(n_A) = \omega(n_A) (1 - R^2)$$

$$c_1 = R \frac{\sigma(n_A)}{\sigma(n_R)}, \quad R = \rho(n_A, n_R)$$



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- STAR measurement is **affected** by correlations to the *reference bin* n_R
- Analysis sets the multiplicity \rightarrow computes the variance and correlation
- Much more different method
- Complicated formula linking F-B and P-C properties

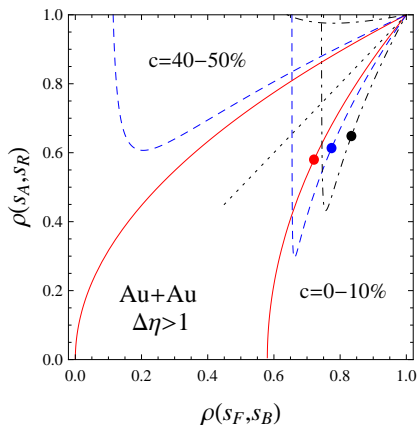
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F-b and p-c relation

$$\rho(\mathbf{s}_A, \mathbf{s}_R)^2 = \frac{\left\{ \left[1 - \frac{\delta}{\omega^*(n_A)} \right] \rho(\mathbf{s}_F, \mathbf{s}_B) - \rho^*(n_F, n_B) \right\}^2}{\left\{ 1 - \rho^*(n_F, n_B) - \frac{\delta}{\omega^*(n_A)} \right\} \left\{ \rho(\mathbf{s}_F, \mathbf{s}_B) - \rho^*(n_F, n_B) - \frac{\delta}{\omega^*(n_A)} \left[\frac{\langle n_A \rangle}{\langle n_R \rangle} (\rho(\mathbf{s}_F, \mathbf{s}_B) - 1) + \rho(\mathbf{s}_F, \mathbf{s}_B) \right] \right\}}$$

Forward-backward and peripheral-center correlation

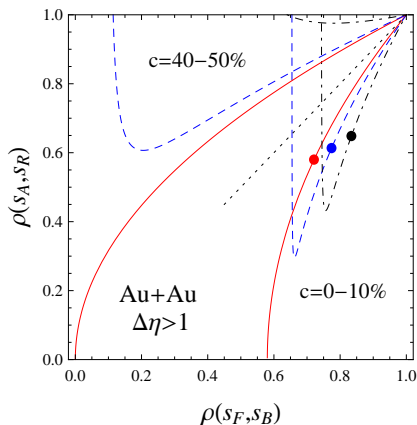


- The dots indicate the estimate for $\rho(n_F, n_B) \simeq 0.72$ [2]

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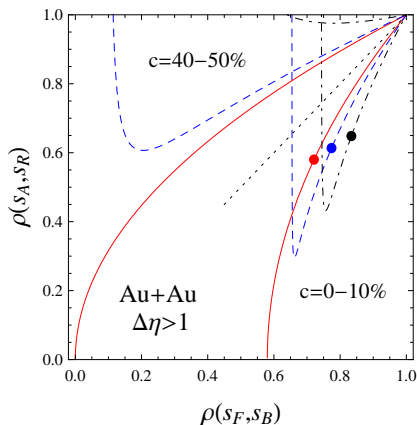


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Forward-backward and peripheral-center correlation



- The dots indicate the estimate for $\rho(n_F, n_B) \simeq 0.72$ [2]
- We consider three δ values and accept the rising parts of the curves
- The most central collisions $\rightarrow \rho(s_F, s_B) > \rho(s_A, s_R)$. The puzzle [1,2]!

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Independent emissions from s sources,

$$n = \sum_{i=1}^n m_i,$$

m_i - number of particles produced by the i th source from some distribution ->

$$\langle n \rangle = \langle s \rangle \langle m \rangle.$$

$$\begin{aligned} \text{var}(n) &= \left\langle \sum_{i=1}^s (\delta m_i + \langle m \rangle) \sum_{j=1}^s (\delta m_j + \langle m \rangle) \right\rangle - (\langle s \rangle \langle m \rangle)^2 \\ &= \left\langle \sum_{i=1}^s \delta m_i^2 \right\rangle + \left\langle \sum_{i,j=1, j \neq i}^s \delta m_i \delta m_j \right\rangle \\ &+ 2 \langle m \rangle \left\langle \sum_{i=1}^s \delta m_i \right\rangle + \langle m \rangle^2 \left\langle \sum_{i=1}^s \sum_{j=1}^s \right\rangle - \langle s \rangle^2 \langle m \rangle^2. \end{aligned}$$

where $\delta m_i = m_i - \langle m \rangle$, with $\langle \delta m_i \rangle = 0$.

$$\text{var}(n) = \langle s \rangle \text{var}(m) + \langle m \rangle^2 \text{var}(s).$$

This leads us to simple relation

Next, we look at the covariance between two well-separated bins, which means $\langle m_i m_j \rangle = m^2$, with i and j belonging to two different bins. We have

$$\langle n_1 n_2 \rangle = \left\langle \sum_{i=1}^s m_i \sum_{j=1}^s m_j \right\rangle = \langle m \rangle^2 \langle s_1 s_2 \rangle,$$

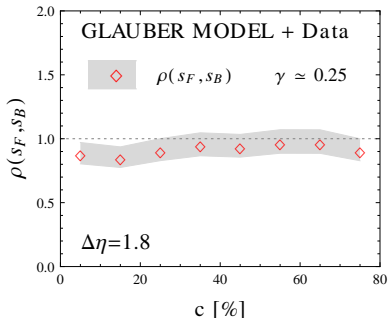
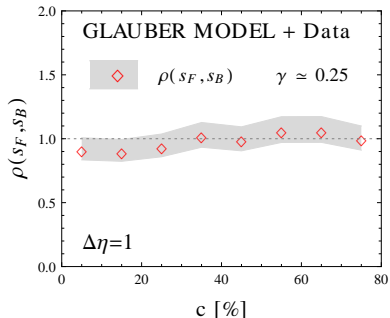
and we get

$$\text{cov}(n_1, n_2) = \langle m \rangle^2 \text{cov}(s_1, s_2).$$

For the correlation coefficient it follows that

$$\rho(n_1, n_2) = \frac{\rho(s_1, s_2)}{\sqrt{1 + \frac{\omega(m)}{\langle m \rangle \omega(s_1)}} \sqrt{1 + \frac{\omega(m)}{\langle m \rangle \omega(s_2)}}}.$$

Correlation prediction



$$\rho(S_F, S_B) = \frac{\text{cov}(n_F, n_B)}{\gamma \text{var}(S_A)}$$