Exact solutions of kinetic equation for massive particles

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Outline

- Motivation
    - Successes of viscous hydrodynamics in description of relativistic heavy-ion collisions — intensive studies of transport coefficients
    - Our idea is to perform comparisons of exact solutions of simple kinetic equations with hydro approaches, which allows us to select correct forms of these coefficients
- Kinetic equation
    - Boltzmann equation
    - Boost-invariant variables
    - Moments of equation
    - Landau matching
    - Numerical method
- Results
    - Time dependence of thermodynamics-like variables
    - Bulk viscousity
    - Shear viscosity
- Conclusions

Motivation

- Experimental and theoretical studies of heavy-ion collisions showed that the behavior of matter produced in such collisions is very well described by viscous hydrodynamics, with a very small viscosity to entropy density ratio
- These results brought a lot of attention to the studies of kinetic coefficients whose values determine the magnitude of important observables such as the elliptic flow
- Interestingly, different methods lead to different values of the kinetic coefficients
- Moreover, the form of the second order hydrodynamic equations depends on the specific values of the kinetic coefficients

Motivation

- Our idea is to perform comparisons of exact solutions of simple kinetic equations with hydrodynamic approaches — this allows for numerical determination of the kinetic coefficients
- Instead of performing complicated simulations based on the Boltzmann equation we analyze its simple form which can be solved exactly (Baym, Heiselberg, Wang, Wong)
- We extend here some of the recent results obtained for massles particles:

  W. Florkowski, R. Ryblewski, M. Strickland, Phys. Rev. C88 (2013) 024903

  W. Florkowski, R. Ryblewski, M. Strickland, Nucl. Phys. A916 (2013) 249

Motivation

LIMITATIONS OF OUR MODEL:

- Collision term treated in the relaxation time approximation (RTA) with a constant equilibration time
- Only longitudinal expansion included (along the z-axis) — justified for early stages of the evolution (1–2 fm/c)
- Boost invariance — justified in the central region ($z \approx 0$)
- All particles have the same mass m

Motivation

ADVANTAGES OF OUR MODEL:

- We find exact solutions of the kinetic equation numerically
- We find the proper forms of shear and bulk viscosities by studying the system's approach towards equilibrium

Kinetic equation
General setup

- Boltzmann equation (BE) in the relaxation-time approximation (RTA)

$$p^\mu \partial_\mu G(x,p) = C[G(x,p)] \qquad C[G] = p\cdot u\, \frac{G^{\rm eq} - G}{\tau_{\rm eq}}$$

background thermal distribution

$$G^{\rm eq} = \frac{2}{(2\pi)^3} \exp(-p\cdot u/T)$$

- boost-invariant variables (Bialas, Czyz)

$$w = tp_\parallel - zE \qquad v = tE - zp_\parallel = \sqrt{w^2 + \left(m^2 + \vec{p}_\perp^{\,2}\right)\tau^2}$$

$$E = \frac{vt + wz}{\tau^2} \qquad p_\parallel = \frac{wt + vz}{\tau^2}$$

- boost-invariant form of the kinetic equation

$$\frac{\partial G}{\partial \tau} = \frac{G^{\rm eq} - G}{\tau_{\rm eq}}$$

$$G^{\rm eq}(\tau, w, p_\perp) = \frac{2}{(2\pi)^3} \exp\left[-\frac{\sqrt{w^2 + (m^2 + p_\perp^2)\tau^2}}{T(\tau)\tau}\right]$$

Kinetic equation
Moments

- zeroth moment (describes particle production)

$$\partial_\mu \int dP\, p^\mu G = \int dP\, C \qquad \frac{dn}{d\tau} + \frac{n}{\tau} = \frac{n^{\mathrm{eq}} - n}{\tau_{\mathrm{eq}}}$$

- first moment (describes energy-momentum conservation)

$$\partial_\mu \underbrace{\int dP\, p^\nu p^\mu G}_{} = \int dP\, p^\nu C = 0 \qquad \frac{d\mathcal{E}}{d\tau} = -\frac{\mathcal{E} + \mathcal{P}_\parallel}{\tau}$$

$$T^{\mu\nu} = (\mathcal{E} + \mathcal{P}_\perp) u^\mu u^\nu - \mathcal{P}_\perp g^{\mu\nu} + (\mathcal{P}_\parallel - \mathcal{P}_\perp) V^\mu V^\nu$$

$$u^\mu = \left(\frac{t}{\tau}, 0, 0, \frac{z}{\tau}\right) \qquad V^\mu = \left(\frac{z}{\tau}, 0, 0, \frac{t}{\tau}\right)$$

- Landau matching

$$\int dP\, p^\nu C = 0$$

- 0th and 1st moments are fulfilled automatically for the exact solution of BE

Kinetic equation

Landau matching

- Landau matching allows us to find effective temperature T

$$\begin{aligned}\mathcal{E}(\tau) &= \mathcal{E}^{\mathrm{eq}}(\tau)\\ \mathcal{E}(\tau) &= \frac{g_0}{\tau^2}\int dP\, v^2\, G(\tau, w, p_\perp)\\ &= \frac{g_0}{\tau^2}\int dP\, v^2\, G^{\mathrm{eq}}(\tau, w, p_\perp)\\ &= \frac{g_0 T m^2}{\pi^2}\left[3TK_2\left(\frac{m}{T}\right) + mK_1\left(\frac{m}{T}\right)\right]\end{aligned}$$

- In the limit of vanishing particle masses:

$$\frac{g_0 T m^2}{\pi^2}\left[3TK_2\left(\frac{m}{T}\right) + mK_1\left(\frac{m}{T}\right)\right] \xrightarrow[\mathrm{m=0}]{} \frac{6 g_0 T^4}{\pi^2}$$

Kinetic equation

Formal solution

- formal structure of the solutions (Baym, Heiselberg, Wang, Wong)

$$G(\tau, w, p_\perp) = D(\tau,\tau_0) G_0(\tau, w, p_\perp) + \int\limits_{\tau_0}^{\tau} \frac{d\tau'}{\tau_{\rm eq}(\tau')} D(\tau,\tau')\, G^{\rm eq}(\tau', w, p_\perp)$$

$$D(\tau_2,\tau_1) = \exp\left[-\int\limits_{\tau_1}^{\tau_2} \frac{d\tau''}{\tau_{\rm eq}(\tau'')} \right]$$

- equilibration time in our calculations is constant

$$\tau_{eq} = 0.25 \ {\rm fm/c}$$

- Romatschke-Strickland (RS) form of the initial condition

$$G_0(w, p_\perp) = \frac{1}{4\pi^3} \exp\left[-\frac{\sqrt{(1+\xi_0) w^2 + (m^2 + p_\perp^2)\tau_0^2}}{\Lambda_0 \, \tau_0} \right]$$

- $1 + \xi_0 = x_0$ - initial value of the anisotropy parameter, Λ_0 defines initial transverse-momentum scale (transverse temperature)

Kinetic equation

Numerical method

$$\begin{aligned}\frac{g_0 T m^2}{\pi^2}\left[3TK_2\left(\frac{m}{T}\right)+mK_1\left(\frac{m}{T}\right)\right] &= \frac{g_0}{2\pi^2}\left[D(\tau,\tau_0)\Lambda_0^4\widetilde{\mathcal{H}}_2\left(\frac{\tau_0}{\tau\sqrt{1+\xi_0}},\frac{m}{\Lambda_0}\right)\right.\\ &+ \left.\int_{\tau_0}^{\tau}\frac{d\tau'}{\tau_{\rm eq}(\tau')}D(\tau,\tau')T'^4\widetilde{\mathcal{H}}_2\left(\frac{\tau'}{\tau},\frac{m}{T'}\right)\right].\end{aligned}$$

- iterative method (Banerjee, Bhalerao, Ravishankar):
  1) use a trial function $T' = T(\tau')$ on the RHS of the dynamic equation
  2) the LHS of the dynamic equation determines the new $T = T(\tau)$
  3) use the new $T(\tau)$ as the trial one
  4) repeat steps 1-3 until the stable $T(\tau)$ is found
- particle density, transverse and longitudinal pressure

$$\begin{aligned}n(\tau) &= \frac{g_0}{\tau}\int dP\, v\, G(\tau,w,p_\perp)\\ \mathcal{P}_\parallel(\tau) &= \frac{g_0}{\tau^2}\int dP\, w^2\, G(\tau,w,p_\perp)\\ \mathcal{P}_\perp(\tau) &= \frac{g_0}{2}\int dP\, p_T^2\, G(\tau,w,p_\perp)\end{aligned}$$

$\widetilde{\mathcal{H}}$ functions

- $\widetilde{\mathcal{H}}_2$, $\widetilde{\mathcal{H}}_{2\|}$, and $\widetilde{\mathcal{H}}_{2\perp}$ functions are defined as integrals:

$$\begin{aligned}
\widetilde{\mathcal{H}}_2(y,z) &= \int_0^\infty dr\, r^3 \mathrm{e}^{-\sqrt{r^2+z^2}} \mathcal{H}_2\left(y, \frac{z}{r}\right), \\
\widetilde{\mathcal{H}}_{2\|}(y,z) &= \int_0^\infty dr\, r^3 \mathrm{e}^{-\sqrt{r^2+z^2}} \mathcal{H}_{2\|}\left(y, \frac{z}{r}\right), \\
\widetilde{\mathcal{H}}_{2\perp}(y,z) &= \int_0^\infty dr\, r^3 \mathrm{e}^{-\sqrt{r^2+z^2}} \mathcal{H}_{2\perp}\left(y, \frac{z}{r}\right)
\end{aligned}$$

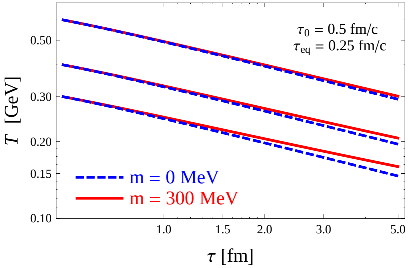
$\mathcal{H}$ functions

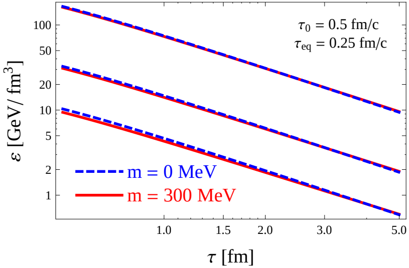
- $\mathcal{H}_2$, $\mathcal{H}_{2\|}$, and $\mathcal{H}_{2\perp}$ functions are defined similarly as:

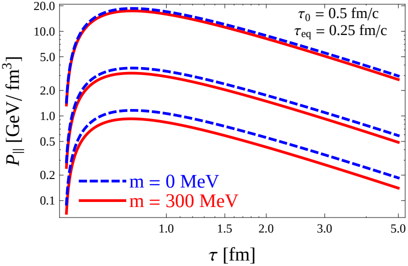
$$\begin{aligned}
\mathcal{H}_2\left(y,\frac{z}{r}\right) &= y\int_0^{\pi} d\phi\,\sin\phi\sqrt{y^2\cos^2\phi+\sin^2\phi+\left(\frac{z}{r}\right)^2},\\
\mathcal{H}_{2\|}\left(y,\frac{z}{r}\right) &= y^3\int_0^{\pi} d\phi\,\frac{\sin\phi\cos^2\phi}{\sqrt{y^2\cos^2\phi+\sin^2\phi+\left(\frac{z}{r}\right)^2}},\\
\mathcal{H}_{2\perp}\left(y,\frac{z}{r}\right) &= y\int_0^{\pi} d\phi\,\frac{\sin^3\phi}{\sqrt{y^2\cos^2\phi+\sin^2\phi+\left(\frac{z}{r}\right)^2}}
\end{aligned}$$

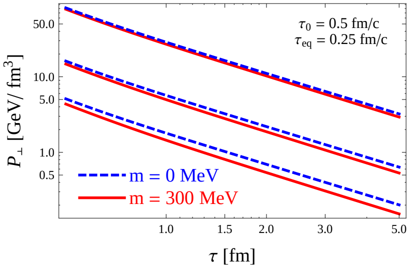
These integrals are analytic but the results are rather lengthy and not shown here.

Thermodynamics-like variables









Bulk viscous pressure

- Bulk pressure in the kinetic theory may be defined as:

$$\Pi^k_\zeta \quad = \quad \frac{1}{3}\left[\mathcal{P}_{\parallel}(\tau) + 2\mathcal{P}_{\perp}(\tau) - 3\mathcal{P}_{\rm eq}(\tau)\right].$$

- When the system approaches equilibrium, we expect

$$\Pi_\zeta(\tau) \quad = \quad -\frac{\zeta(T(\tau))}{\tau},$$

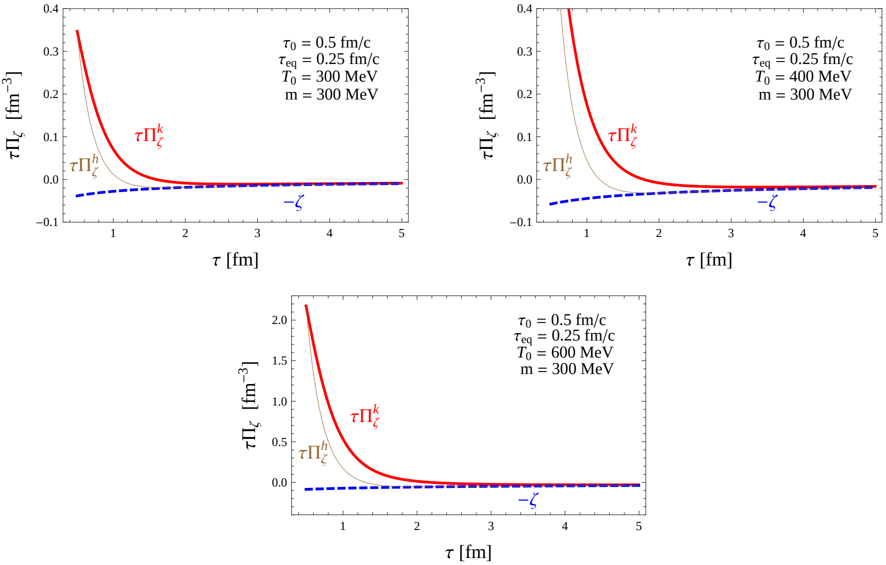
- where the bulk viscosity is given by the formula (Redlich and Sasaki, PRC **79** (2009) 055207; Bożek, PRC **81** (2010) 034909):

$$\zeta(T) \quad = \quad \frac{g_0 m^2}{3\pi^2 T}\int\limits_0^\infty p^2 e^{-\frac{\sqrt{m^2+p^2}}{T}}\left[c_s^2(T) - \frac{p^2}{3(m^2+p^2)}\right] dp.$$

- Hydrodynamic predictions for the time dependence of the bulk viscous pressure (Israel, Stewart):

$$\tau_{\rm eq}\left(\frac{d\Pi^h_\zeta}{d\tau} + \frac{4\Pi^h_\zeta}{3\tau}\right) + \Pi^h_\zeta \quad = \quad -\frac{\zeta}{\tau}.$$

Comparison with exact solutions - bulk viscosity



Bulk viscosity - Anderson and Witting

- Anderson and Witting formula, Physica **74** (1974) 466, Physics for bulk viscosity:

$$\zeta = \frac{\tau p}{3}\frac{m}{T}\left[\frac{3\left(G^2\zeta - 5G - \zeta\right)}{\zeta^2 + 5G\zeta - G^2\zeta^2 - 1} + \frac{\zeta^2}{3}\left(\frac{3G}{\zeta^2} - \frac{1}{\zeta} + \frac{K_1}{K_2} - \frac{Ki_1}{K_2}\right)\right]$$

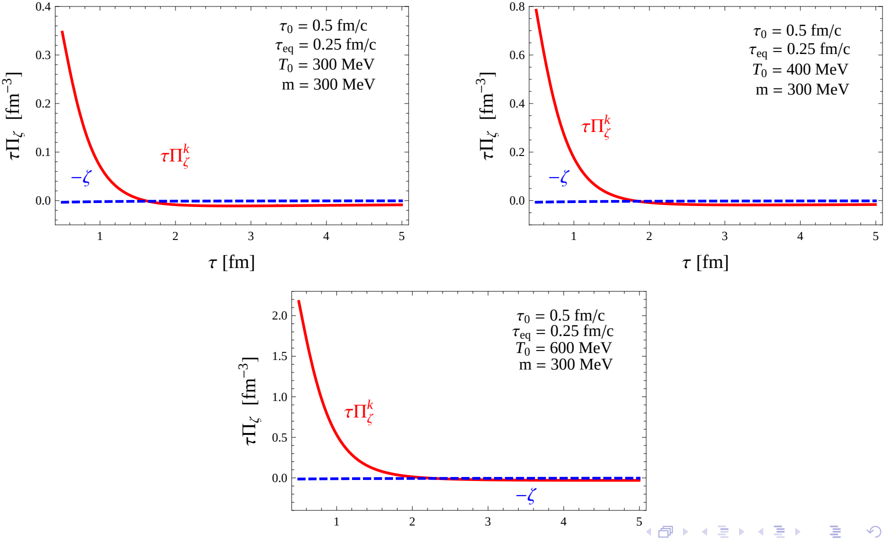
here

$$\zeta = \frac{m}{T},$$

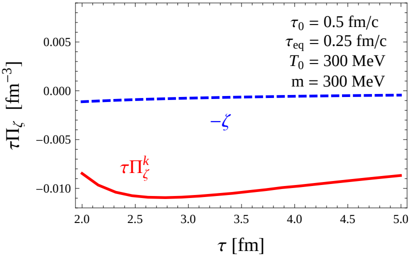
$$G = \frac{K_3}{K_2},$$

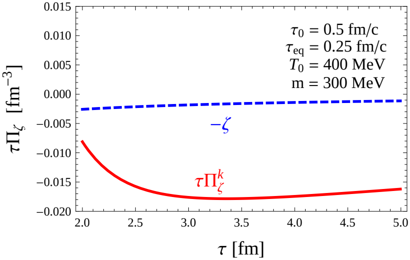
$$K_{i,n}(\zeta) = \int\limits_{\zeta}^{\infty} K_{i,n-1}(t)dt = \int\limits_{0}^{\infty} \frac{\mathrm{e}^{-\zeta \cosh t}}{\cosh^n t}dt$$

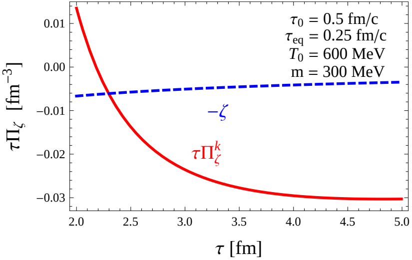
Comparison with exact solutions - Anderson-Witting formula for bulk viscosity



Comparison with exact solutions - bulk viscosity







Shear viscous pressure

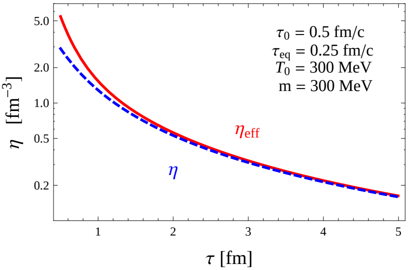
- The calculation by Anderson and Witting, Physica **74** (1974) 466, gives the shear viscosity coefficient in the form

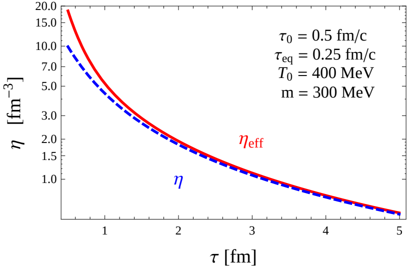
$$\eta = \frac{\tau p}{15}\left(\frac{m}{T}\right)^3\left[\frac{3T^2}{m^2}\frac{K_3}{K_2} - \frac{T}{m} + \frac{K_1}{K_2} - \frac{Ki_1}{K_2}\right]$$

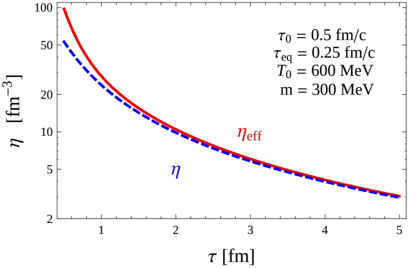
- From the kinetic equation we obtain the effective shear viscosity as

$$\eta_{\rm eff} = \frac{\left(\mathcal{P}_\perp - \mathcal{P}_\parallel\right)\tau}{2}$$

Comparison with exact solutions - shear viscosity







Conclusions

- We have constructed exact solutions of the one-dimensional boost-invariant kinetic equation treated in the relaxation time approximation.
- The previous approaches valid for massless particles have been generalized.
- We have established the correspondence between the late, near equilibrium evolution of the system described by the kinetic theory and by the viscous hydrodynamics.
- We have shown that the late time behavior of the bulk viscous pressure is determined by the bulk viscosity formula used, e.g., by Bozek and Redlich. On the other hand, a disagreement has been found with the Anderson-Witting formula.
- On the other hand, Anderson-Witting formula for the shear viscosity works well in the case of massive particles (and also for massless particles, as it was shown before).

Thank You