

Hydrodynamics in small and medium systems

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1. **p-Pb** (research with P. Bożek)
2. **α -clustered nucleus-A** (research with E. Ruiz Arriola and M. Rybczyński)

Main questions:

What is the nature of the initial state and correlations therein?

What are the limits/conditions on applicability of hydrodynamics?

What is the ground state of light nuclei?

Other analyses of collectivity in small systems:

Romatschke, Luzum, arXiv:0901.4588, Prasad et al., arXiv:0910.4844,

Bozek, arXiv:0911.2393, Werner et al., arXiv:1010.0400,

Deng, Xu, Greiner, arXiv:1112.0470, Yan et al., arXiv: 0912.3342,

Bozek, arXiv:1112.0912 Shuryak, Zahed, arXiv:1301.4470,

Bzdak et al., arXiv:1304.3403, Qin, Müller, arXiv:1306.3439,

Werner et al., arXiv:1307.4379

Method

Flow develops



3-stage approach

Our three-phase approach (“Standard Model of heavy-ion collisions”):

initial \rightarrow hydro \rightarrow statistical hadronization

(successful in description of A+A collisions)

- **Initial phase** - Glauber model GLISSANDO
- **Hydrodynamics** - 3+1 D viscous event-by-event
- **Statistical hadronization** - THERMINATOR

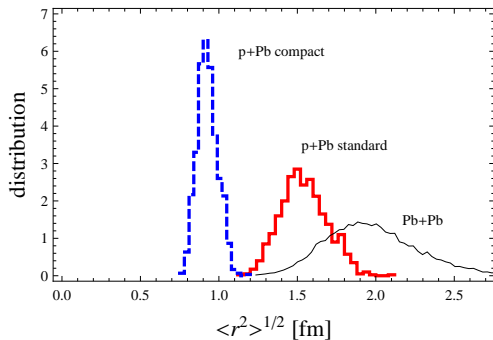
3+1 D viscous Israel-Stewart event-by-event hydrodynamics (viscous corrections essential due to large gradients)

- $\tau_{\text{init}} = 0.6 \text{ fm}/c$, $\eta/s = 0.08$ (**shear**), $\zeta/s = 0.04$ (**bulk**)
- Gaussian smearing of the sources, $r = 0.4 \text{ fm}$ – physical effect
- average initial temperature in the center of the fireball adjusted to fit the multiplicity
- realistic equation of state (lattice + hadron gas [based on Chojnacki, Florkowski 2007])
- freezeout at $T_f = 150 \text{ MeV}$
- lattice spacing of 0.15 fm (**thousands of CPU hours for one reaction**)

Highlights of p-Pb

Size in p-Pb vs Pb-Pb

fixed $N_{\text{part}} = 19$



smaller size in p-Pb \rightarrow larger entropy density \rightarrow more rapid expansion

standard - source positioned at the center of the nucleons

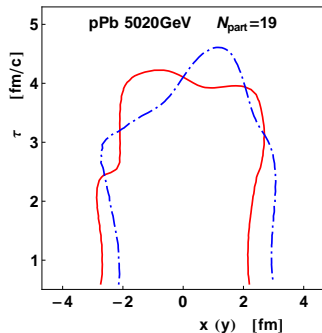
compact - source at the CM of the colliding pair

[see also Bzdak, Schenke, Tribedy, Venugopalan, arXiv:1304.3403]

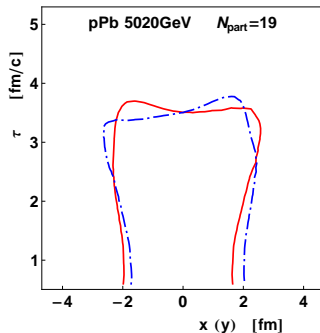
All in all, initial conditions in most central p-Pb not very far from those in peripheral Pb-Pb

Typical evolution in p-Pb

standard



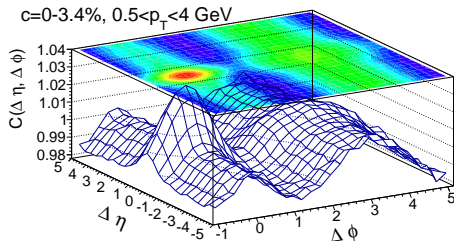
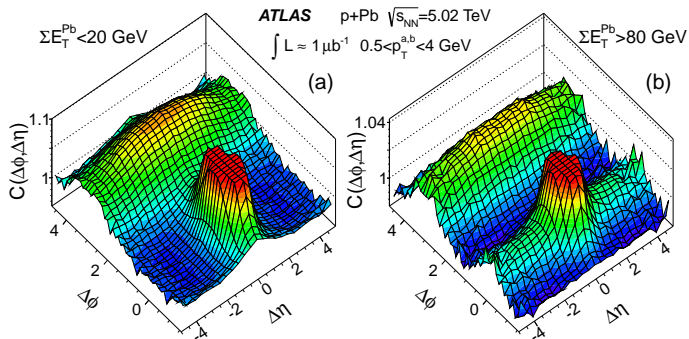
compact



isotherms at freeze-out $T_f = 150$ MeV for two sections in the transverse plane

evolution lasts about 4 fm/c - shorter but more rapid than in A+A

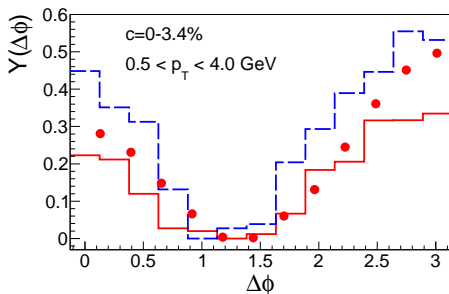
Ridge in p-Pb, ATLAS



Projection on $2 \leq |\Delta\eta| \leq 5$, ATLAS

$$Y(\Delta\phi) = \frac{\int B(\Delta\phi)d(\Delta\phi)}{N}C(\Delta\phi) - b_{ZYAM}$$

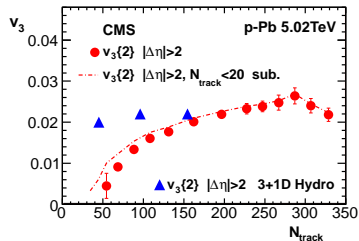
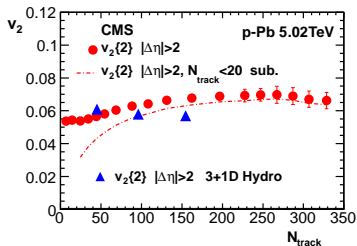
The near-side ridge from our model:



red - standard, blue - compact

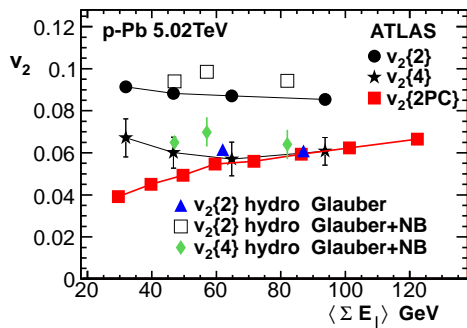
[CGC: Dusling, Venugopalan, arXiv:1210.3890, 1211.3701, 1302.7018]

v_2, v_3 vs CMS

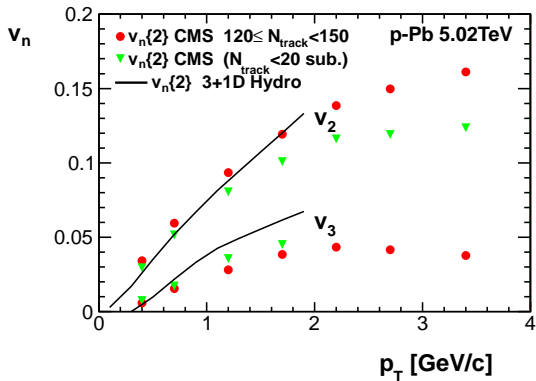


v_3 somewhat too large for lower-multiplicity collisions

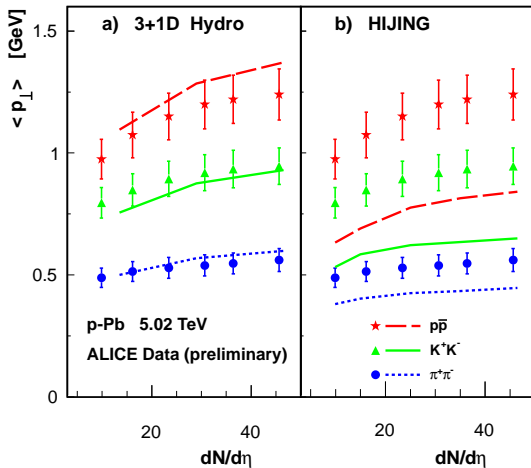
→ limit of validity of the model



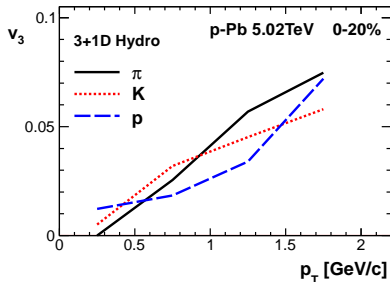
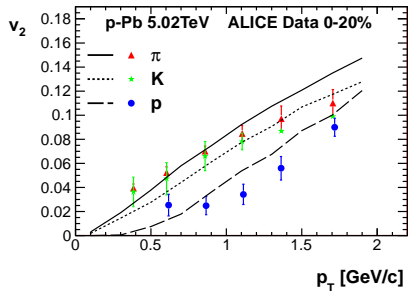
(NB – additional fluctuations of the strength of the sources, adjusted in such a way that the experimental multiplicity distribution is reproduced)



[Bożek, WB, Torrieri, PRL 111 (2013) 172303]



Identified v_2 and v_3



Resonance decays affect the mass ordering in v_3

α clusters

Some history

David Brink: After Gamow's theory of α -decay it was natural to investigate a model in which nuclei are composed of α -particles. Gamow developed a rather detailed theory of properties in his book "Constitution of Nuclei" published in 1931 before the discovery of the neutron in 1932. He supposed that $4n$ -nuclei like ${}^8\text{Be}$, ${}^{12}\text{C}$, ${}^{16}\text{O}$... were composed of α -particles

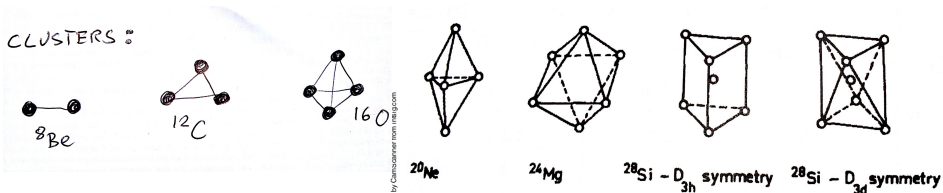


Fig. 1. Alpha-particle configuration for some $4N$ nuclei.

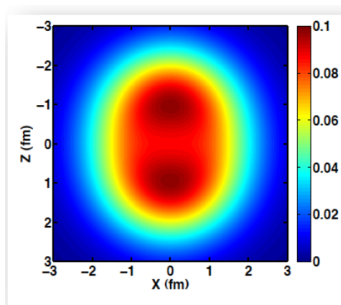
Generated by CamScanner from intsig.com

Shell model (and its problems)

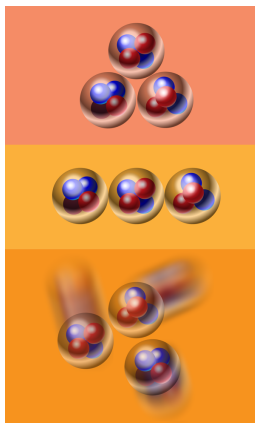
Eugene Wigner, Maria Goeppert-Mayer, Hans Jensen, Nobel in 1963

Michael P. Carpenter: *However, in the 1960s, excited states in nuclei that comprise equal numbers of protons and neutrons, (e.g., ^{12}C and ^{16}O) were identified that could not be described by the shell model, and it was suggested by Ikeda and others that these states could be associated with configurations composed of α particles*

Present theory status



${}^9\text{Be}$



${}^{12}\text{C}$

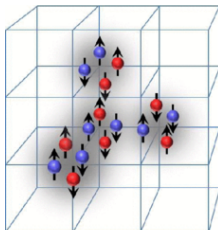
ground

Hoyle 0^+

other excited, 2^+ ...

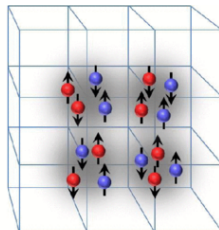
[M. Freer: WPCF13, H. Fynbo+Freer: Physics 4 (2011) 94]

Ab initio calculations of ^{16}O with chiral NN force (Juelich 2014)
→ strong α clusterization



(a) Initial state "A",
8 equivalent orientations.

ground state



(b) Initial states "B" and "C",
3 equivalent orientations.

excited

Computational techniques

(massive effort)

Funaki et al.: *certain states in self-conjugated nuclei ... can be described as product states of α particles, all in the lowest $0S$ state. We define a state of condensed α particles in nuclei as a bosonic product state in good approximation, in which all bosons occupy the lowest quantum state of the corresponding bosonic mean-field potential (α BEC)*

Another approach: Fermionic Molecular Dynamics (FMD)

Quantum Variational Monte Carlo (with 2- and 3-body forces) for $A=2-12$
[R. Wiringa et al., <http://www.phy.anl.gov/theory/research/density/>]

All approaches to light nuclei give clusters

Goal (not yet accurately reached):

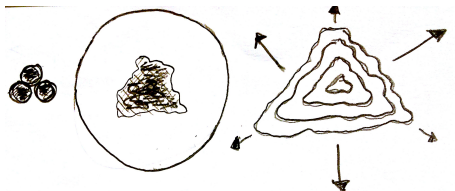
reproduce ground-state energy, excitation spectrum, EM form factor, ...

Merge the two ideas: α 's and flow

From α clusters to flow in relativistic collisions

[WB, Ruiz Arriola, PRL 112 (2014) 112501]

α clusters \rightarrow asymmetry of shape \rightarrow asymmetry of initial fireball \rightarrow
 \rightarrow hydro or transport \rightarrow collective harmonic flow



nuclear triangular geometry \rightarrow fireball triangular geometry \rightarrow triangular flow

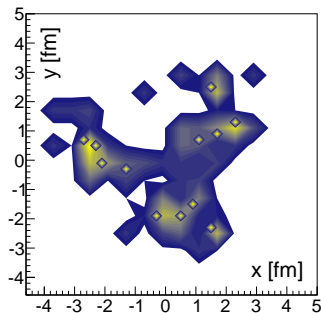
What are the signatures, chances of detection?

Related idea: triton/ ^3He -Au at RHIC in 2015 [Sickles (PHENIX) 2013]

The case of light nuclei is more promising, as it leads to abundant fireballs.

why ultrarelativistic?

reaction time is much shorter than time scales of the structure
→ a frozen “snapshot” of the nuclear configuration

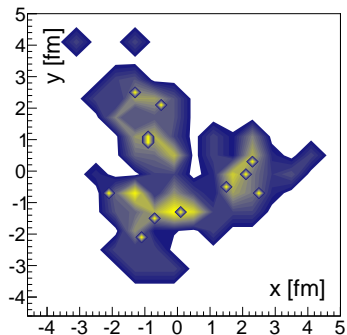


($N_w > 70$ - flat-on orientation)

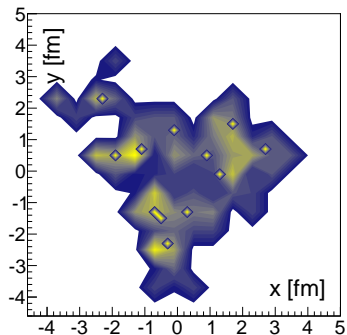
Imprints of the three α clusters clearly visible

[simulations with GLISSANDO 2]

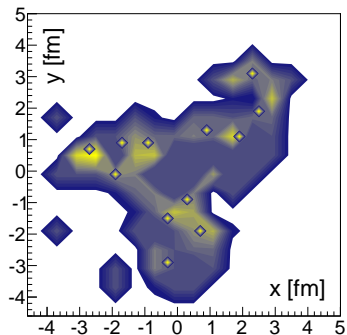
... more events



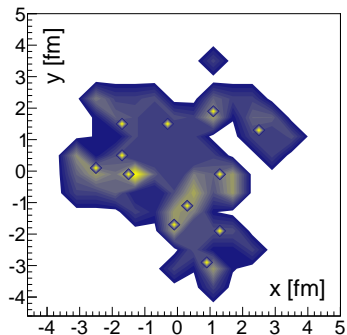
... more events



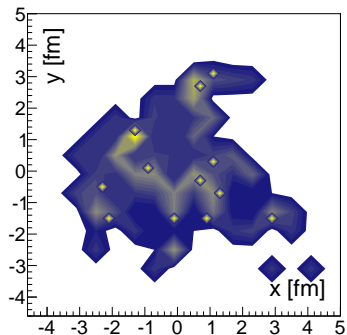
... more events



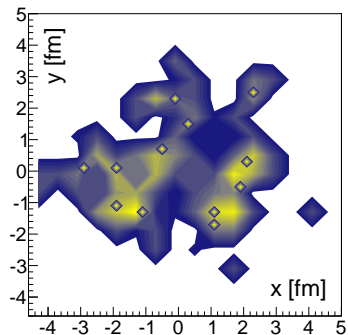
... more events



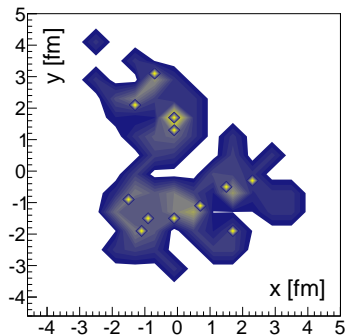
... more events



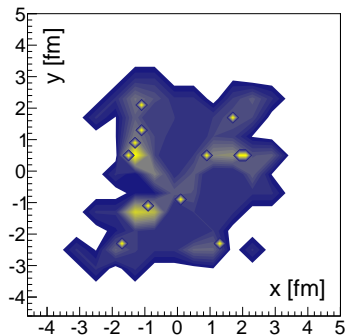
... more events



... more events



... more events



The meaning of *intrinsic*

Ground state of ^{12}C is a 0^+ state (rotationally symmetric wave function).
The meaning of *deformation* concerns **multiparticle correlations** between the nucleons

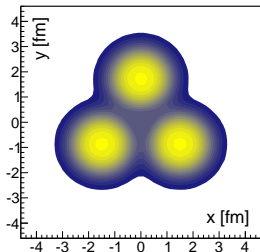
Superposition over orientations:

$$|\Psi_{0^+}(x_1, \dots, x_N)\rangle = \frac{1}{4\pi} \int d\Omega \Psi_{\text{intr}}(x_1, \dots, x_N; \Omega)$$

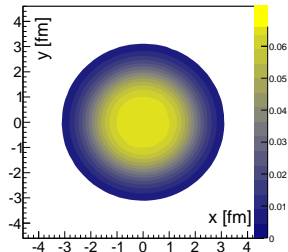
The *intrinsic* density of sources of rank n is defined as the average over events, where the distributions in each event have aligned principal axes:
 $f_n^{\text{intr}}(\vec{x}) = \langle f(R(-\Phi_n)\vec{x}) \rangle$. Brackets indicate averaging over events and $R(-\Phi_n)$ is the inverse rotation by the principal-axis angle in each event

Back to ^{12}C – intrinsic density

Intrinsic distributions in ^{12}C : three α 's in a triangular arrangement

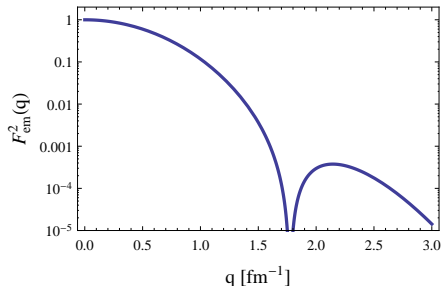
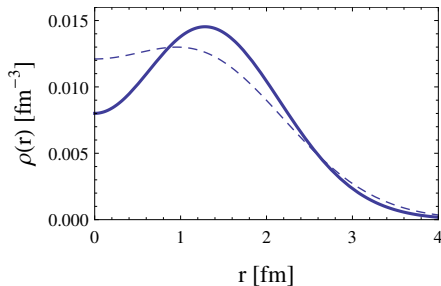


clustered



unclustered

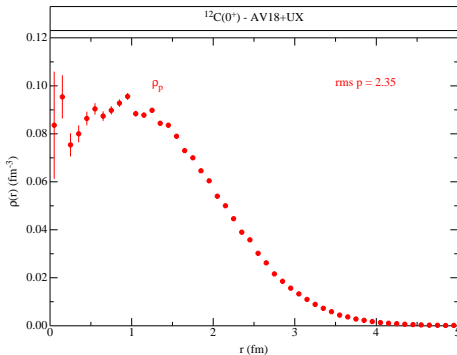
Constraints on ^{12}C from EM form factor



Electric charge density (dashed line) and the corresponding distribution of the centers of protons (solid line) in ^{12}C for the data plotted against the radius, for the **BEC** calculation – agrees with the experimental data for the charge form factor

Central depletion naturally explained with the hole between the clusters

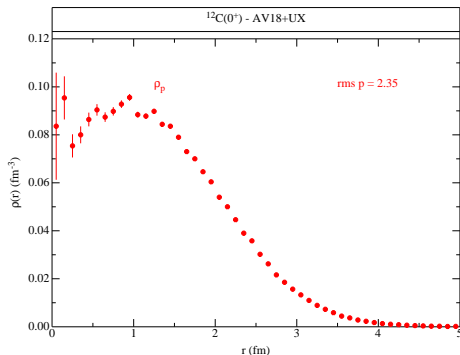
^{12}C from Wiringa's MC



Distribution of the centers of protons = neutrons in ^{12}C

smaller central depletion

^{12}C from Wiringa's MC



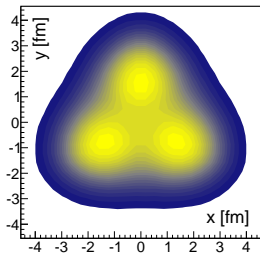
Distribution of the centers of protons = neutrons in ^{12}C

smaller central depletion

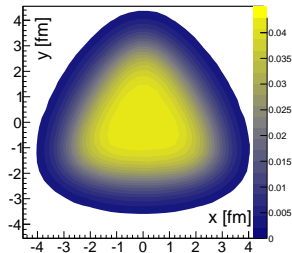
GLISSANDO implements these clustered distributions

→ carry out detailed simulations

Intrinsic distributions in the *transverse plane* of the fireball (here with $N_w > 70$ – large multiplicity enforcing the flat-on collision)

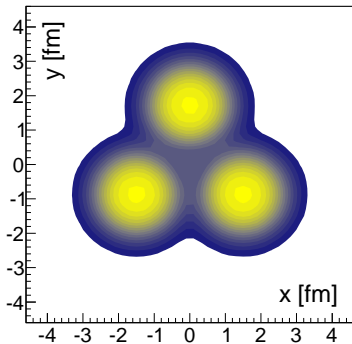


clustered



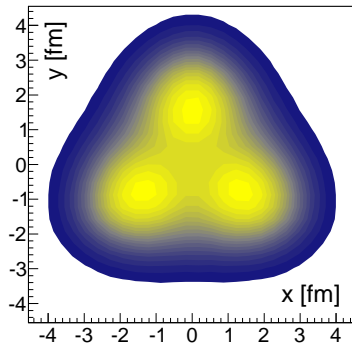
unclustered

Geometry of nucleus \rightarrow geometry of fireball



intrinsic density of ^{12}C

\rightarrow



geometry of the fireball

Eccentricity parameters ϵ_n (Fourier analysis)

$$\epsilon_n e^{in\Phi_n} = \frac{\sum_j \rho_j^n e^{in\phi_j}}{\sum_j \rho_j^n}$$

describe the shape of each event (j labels the sources in the event, n =rank, Φ_n is the principal axis angle)

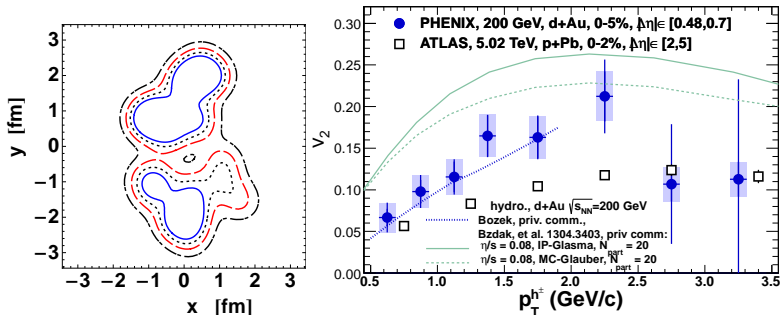
Two components:

- internal (from existent mean deformation of the fireball)
- from fluctuations

Digression: d-A by Bożek

The deuteron has an intrinsic dumbbell shape with very large deformation: rms $\simeq 2$ fm

Initial entropy density in a d-Pb collision with $N_{\text{part}} = 24$ [Bożek 2012]



Resulting large elliptic flow confirmed with the later RHIC analysis

Geometry vs multiplicity correlations in ^{12}C -Pb

Specific feature of the ^{12}C collisions:

The cluster plane parallel or perpendicular to the transverse plane:

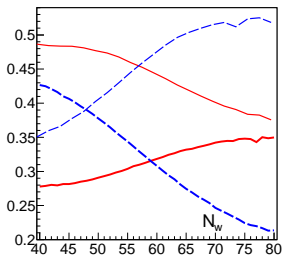


higher multiplicity
higher triangularity
lower ellipticity

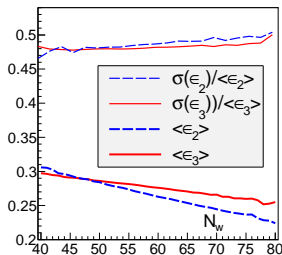


lower multiplicity
lower triangularity
higher ellipticity

Ellipticity and triangularity vs multiplicity



clustered



unclustered

Clusters:

When $N_w \nearrow$ then $\langle \epsilon_3 \rangle \nearrow$ and $\langle \epsilon_2 \rangle \searrow$

and $\langle \sigma(\epsilon_3)/\epsilon_3 \rangle \searrow$, $\langle \sigma(\epsilon_2)/\epsilon_2 \rangle \nearrow$ tending to $\sqrt{4/\pi - 1} \sim 0.52$

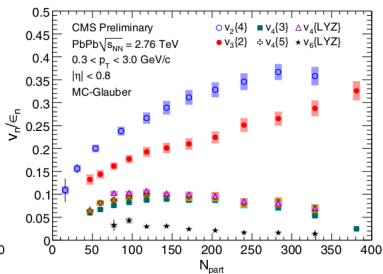
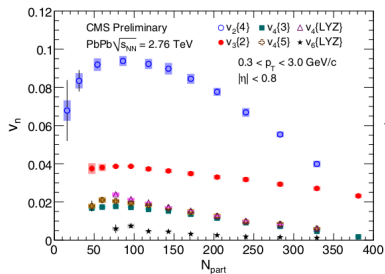
No clusters:

similar behavior for $n = 2$ and $n = 3$

Shape-flow transmutation

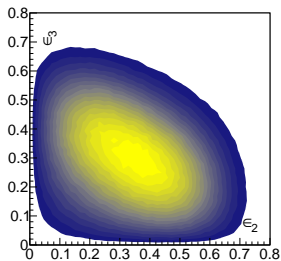
The eccentricity parameters are transformed (in all models based on collective dynamics) into asymmetry of the transverse-momentum flow
It has been found that

$\langle v_n \rangle$ grows with $\langle \epsilon_n \rangle$

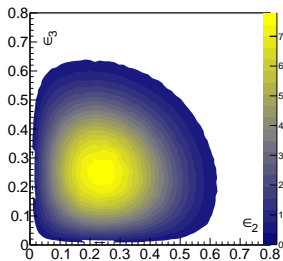


→ for ^{12}C collisions v_3 should grow with multiplicity even stronger than ϵ_3

Triangularity vs ellipticity



clustered

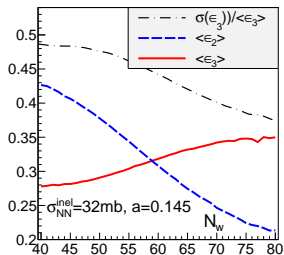


unclustered

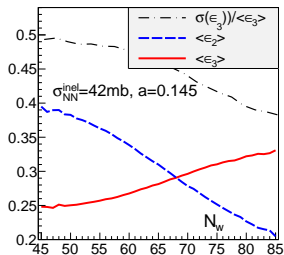
Clusters:

Anticorrelation of ϵ_2 and ϵ_3

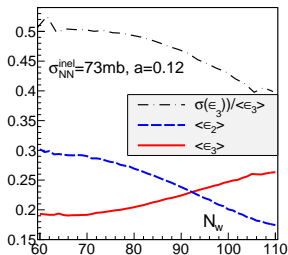
Dependence on the collision energy



32mb (SPS)



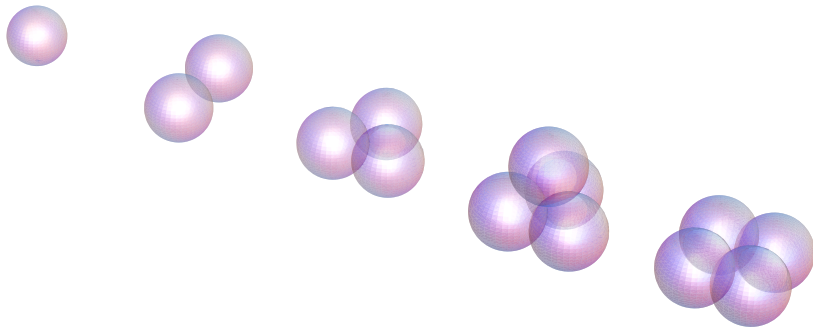
42mb (RHIC)



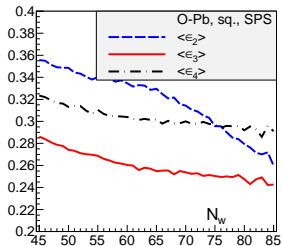
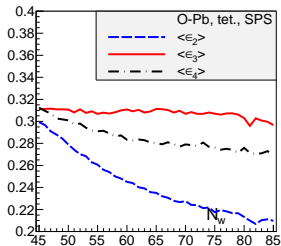
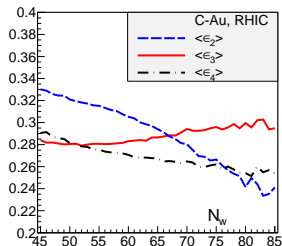
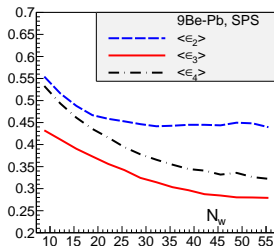
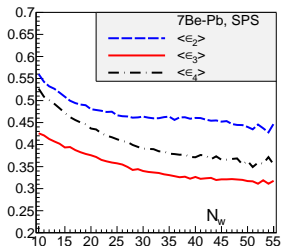
72mb (LHC)

Qualitative conclusions hold from SPS to the LHC

Other systems



Other systems (Wiringa's distributions)



[work with Maciej Rybczyński]

Conclusions

Is there collectivity in p-Pb?

Collective dynamics is compatible with the high-multiplicity soft LHC data for p-Pb

- Large v_2 and v_3 coefficients measured in p-Pb reasonably reproduced, including the p_T dependence
- Model 2D correlations exhibit the two ridges, in particular the **near-side ridge**
- Mass ordering in $\langle p_T \rangle$ and flow coefficients reproduced

- Numerous effects should still be incorporated (jets, core-corona, ...) – more important for the lower-multiplicity events

Clustered geometry of the ground state \rightarrow harmonic flow

Best signatures of $^{12}\text{C}-^{208}\text{Pb}$ collisions

- Increase of triangularity with multiplicity for the highest multiplicity events
- Decrease of scaled variance of triangularity with multiplicity for the highest multiplicity events
- Anticorrelation of ellipticity and triangularity

Extensions (in progress)

- Other systems, more detailed modeling involving e-by-e hydro

Possible future data (NA61, RHIC?) in conjunction with a detailed knowledge of the dynamics of the evolution of the fireball would allow to place constraints on the α -cluster structure of the colliding nuclei.

Conversely, the knowledge of the clustered nuclear distributions may help to verify the fireball evolution models