Fluctuations of the initial condition from the Glauber models*

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Wounded nucleons

- Wounded nucleons [Białas, Błeszyński, Czyż, 1976] are a basic concept in heavy-ion analyses
- RHIC: 86% wounded + 14% binary describes the multiplicities
- Determination of centrality, impact parameter, analyses of fluctuations, ...
- Basic steps in simulations: 1) generate nucleon positions in nuclei from a WS distribution in an independent way (no correlations), 2) count collisions



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The NN force

Correlations in nuclei are important for fluctuations [Baym, Blattel, Frankfurt, Heiselberg, Strikman, 1995]



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NN correlations

Recent progress: A Monte Carlo generator of nucleon configurations in complex nuclei including Nucleon-Nucleon correlations [Alvioli, Drescher, Strikman, arXiv:0905.2670]



generate configurations with the inclusion of two-body correlations according to the wave function

 $\Psi(r_1,\ldots,r_N) = \prod_{i < j} f(r_{ij}) \Phi_{\text{indep.}}(r_1,\ldots,r_N)$

Metropolis algorithm



GLauber Initial State Simulations AND mOre [WB, Rybczyński, Bożek, Oct. 2007, Comput. Phys. Commun. 180(2009)69] read the provided distributions with correlations

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Multiplicity, NA49 setup

 N_p^{PROJ} - number of wounded nucleons in the projectile $\langle N \rangle = kqN_w/2$, where k = 3 - number of charged particles produced per wounded pair, q = 0.2 - detector acceptance (SPS setup)



std - independent with expulsion distance $a = 0.4^{\circ}$ std, d = 0 - independent with no expulsion

mod - with correlations, configurations from Alvioli et al.

fluctuations reduced

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Eccentricity, fixed axes

$$\epsilon = \frac{\langle y \rangle^2 - \langle x \rangle^2}{\langle y \rangle^2 + \langle x \rangle^2}$$



fluctuations slightly decreased

Eccentricity, variable axes

 ϵ^* computed in an event-by-event rotated frame (principal-axes frame), where it is maximal ("participant" eccentricity)



for central collisions the limit $\Delta \epsilon^*/\epsilon^* = \sqrt{4/\pi - 1} \simeq 0.52$ follows (for uncorrelated systems) from the central limit theorem [WB, Bożek, Rybczyński, Phys. Rev. C76(2007)054905]

Summary of correlation effects

- One-body observables are left intact by two-body correlations
- Event-by-event fluctuations are reduced, as the distribution with fluctuations is more regular
- Effects for the multiplicity fluctuations are substantial
- Effects for the eccentricity studies are very small

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Size fluctuations

(so far the study without correlations discussed above)



compute event-by-event the average size, here defined as

$$\langle r \rangle = \frac{1}{N_w} \sum_{i=1}^{N_w} \sqrt{x_i^2 + y_i^2}$$

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 $\left<\left<.\right>\right>$ - averaging also over events

Size fluctuation



independent of energy from SPS ($\sigma_{NN} = 32 \text{ mb}$) to LHC ($\sigma_{NN} = 63 \text{ mb}$)

Hydrodynamics

Hydro carries over the initial size fluctuation to $\langle p_T \rangle$ fluctuations

- $\bullet~$ initial state $\rightarrow~$ hydrodynamics $\rightarrow~$ freezeout $\rightarrow~$ hadrons
- more compressed initial condition leads to a faster build-up of flow, and then higher transverse velocity at freezeout, which in turn leads to higher $\langle p_T \rangle$
- $\Delta p_T / \langle p_T \rangle \simeq A \Delta \langle r \rangle / \langle \langle r \rangle \rangle$
- we estimate the proportionality constant via simulations with hydro (Lhyquid - M. Chojnacki, W. Florkowski) and THERMINATOR

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Best cases - solution of the HBT puzzle

Consider the solutions of hydro with Gaussian initial conditions







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Fluctuations of the initial condition

Fluctuations of the FO surface

Fluctuations of the size of the initial condition \to hydro \to fluctuations of the freezeout surface and velocity



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p_T fluctuations

fluctuations of the freezeout surface and velocity \to THERMINATOR \to fluctuations e-by-e of $\langle p_T \rangle$

- $c = 0 5 \quad \rightarrow N_w \simeq 350 \qquad \sigma_{\rm dyn}(p_T)$ $c = 20 30 \rightarrow N_w \simeq 160 \qquad \sigma_{\rm dyn}(p_T)$ $c = 60 70 \rightarrow N_w \simeq 30 \qquad \sigma_{\rm dyn}(p_T)$
- $\begin{aligned} \sigma_{\rm dyn}(p_T)/\langle p_T \rangle &\simeq 1.4\% \\ \sigma_{\rm dyn}(p_T)/\langle p_T \rangle &\simeq 2.3\% \\ \sigma_{\rm dyn}(p_T)/\langle p_T \rangle &\simeq 4.5\% \end{aligned}$

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p_T fluctuations

fluctuations of the freezeout surface and velocity \to THERMINATOR \to fluctuations e-by-e of $\langle p_T \rangle$

 $\begin{array}{ll} c = 0 - 5 & \rightarrow N_w \simeq 350 & \sigma_{\rm dyn}(p_T)/\langle p_T \rangle \simeq 1.4\% \\ c = 20 - 30 \rightarrow N_w \simeq 160 & \sigma_{\rm dyn}(p_T)/\langle p_T \rangle \simeq 2.3\% \\ c = 60 - 70 \rightarrow N_w \simeq 30 & \sigma_{\rm dyn}(p_T)/\langle p_T \rangle \simeq 4.5\% \end{array}$



p_T fluctuations, various data



[data from PRC 72 (2005) 044902]

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Summary of size fluctuations

- a few percent effect, explains the experimental p_T fluctuations
- proper scaling with N_w: σ²_{dyn}(p_T) ∼ 1/N_w − proper dependence on centrality
- many attempts made previously: (mini) jets, clusters, temperature fluctuations
- "geometric/statistical" origin
- weak dependence on energy

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