## Alpha clustering from relativistic collisions

#### Wojciech Broniowski

UJK Kielce & IFJ PAN Cracow

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[based on WB& E. Ruiz Arriola, arXiv:1312.0289]

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#### Instead of outline

(WPCF 2013 
$$\rightarrow \alpha$$
)

#### Two phenomena are related:

 $\alpha$  clustering in light nuclei



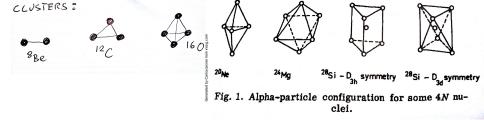
harmonic flow in ultra-relativistic A+B collisions

low-energy structure ←→ highest energy mini bangs (!)

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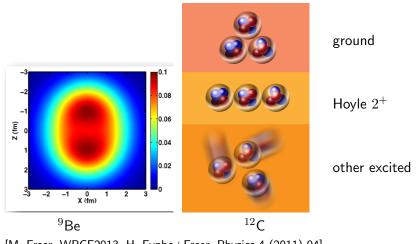
## History

David Brink: After Gamow's theory of  $\alpha$ -decay it was natural to investigate a model in which nuclei are composed of  $\alpha$ -particles. Gamow developed a rather detailed theory of properties in his book "Constitution of Nuclei" published in 1931 before the discovery of the neutron in 1932. He supposed that 4n-nuclei like  $^8\text{Be},\ ^{12}\text{C},\ ^{16}\text{O}$  ... were composed of  $\alpha$ -particles



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### Present status

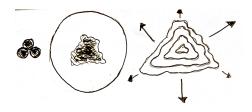


[M. Freer, WPCF2013, H. Fynbo+Freer, Physics 4 (2011) 94]

ab initio calculations up to  $^{16}O\longleftrightarrow$  strong  $\alpha$  clusterization

#### From $\alpha$ clusters to flow in relativistic collisions

 $\begin{array}{c} \alpha \text{ clusters} \to \text{asymmetry of shape} \to \text{asymmetry of initial fireball} \to \\ & \to \text{hydro or transport} \to \text{collective harmonic flow} \end{array}$ 

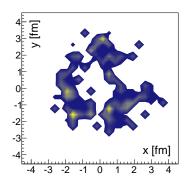


What are the chances of detection?

Related idea:  $triton/^3He$ -Au at RHIC in 2015 [Sickles (PHENIX) 2013] The case of light nuclei is more promising, as it leads to abundant fireballs

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# <sup>12</sup>C-<sup>208</sup>Pb – single event



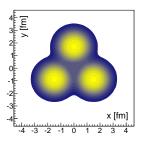
Imprints of the  $\alpha$  clusters clearly visible

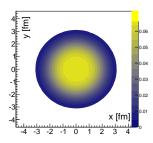
4 11 2 4 4 12 2 4 12 2 2 4 12 2 2 4 12 2

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## $^{12}\text{C}$ - $^{208}\text{Pb}$ – intrinsic average over events

#### Intrinsic distributions: 3 $\alpha$ 's in a triangular arrangement



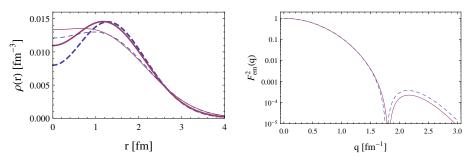


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## Constraints from EM form factor



Electric charge density (thin lines) and the corresponding distribution of the centers of nucleons (thick lines) in  $^{12}\mathrm{C}$  for the data and BEC calculations (dashed lines), and for the FMD calculations (solid lines), plotted against the radius.

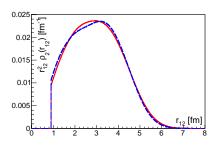
Central depletion

4 D > 4 A > 4 E > 4 E > E 9 Q P

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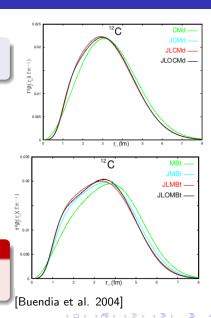
## Distribution of pairs

# Radial density in the relative NN distance $r_{12}$



#### Our Monte Carlo

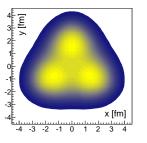
The  $\alpha$  cluster structure is modeled sufficiently accurately

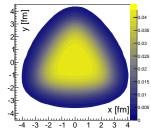


## <sup>12</sup>C-<sup>208</sup>Pb collision

Mixed Glauber model at SPS conditions:  $n \sim \frac{1-a}{2} N_{\rm w} + a N_{\rm bin}$ , a=0.12

Intrinsic distributions in the transverse plane in the fireball,  $N_w > 70$  – large multiplicity





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## **Eccentricity parameters**

Eccentricity parameters

$$\epsilon_n e^{in\Phi_n} = \frac{\sum_j \rho_j^n e^{in\phi_j}}{\sum_j \rho_j^n}$$

describe the shape (j labels the sources in the event, n=rank)

#### Two components:

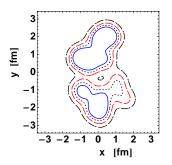
- intrinsic (from existent mean deformation of the fireball)
- from fluctuations



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## Digression: d-Pb

Initial entropy density in a d-Pb collision with  $N_{
m part}=24$  [Bozek 2012]



Fluctuations around the intrinsic ellipticity

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## Geometry vs multiplicity in <sup>12</sup>C-Pb

The triangle plane parallel or perpendicular to the transverse plane:

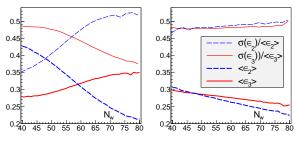


higher multiplicity higher triangularity lower ellipticity



lower multiplicity lower triangularity higher ellipticity

## Ellipticity and triangularity vs multiplicity



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#### Clusters:

When  $N_w \nearrow$  then  $\langle \epsilon_3 \rangle \nearrow$  and  $\langle \epsilon_2 \rangle \searrow$ 

and  $\langle \sigma(\epsilon_3)/\epsilon_3 \rangle$  ,  $\langle \sigma(\epsilon_2)/\epsilon_2 \rangle$  / tending to  $\sqrt{4/\pi-1} \sim 0.52$ 

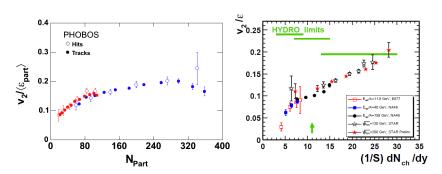
#### No clusters:

similar behavior for n=2 and n=3

## Shape-flow transmutation

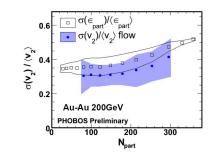
The eccentricity parameters are transformed (in all models based on collective dynamics) into asymmetry of the transverse-momentum flow. It has been found that

$$\langle v_n \rangle \simeq A \langle \epsilon_n \rangle$$



## E-by-e fluctuations

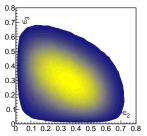
$$\frac{\sigma(v_n)}{\langle v_n \rangle} \simeq \frac{\sigma(\epsilon_n)}{\langle \epsilon_n \rangle}$$

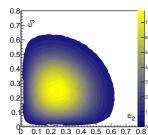


 $\longrightarrow$ 

Measured flow coefficients reflect the initial shape eccentricities

## Triangularity vs ellipticity





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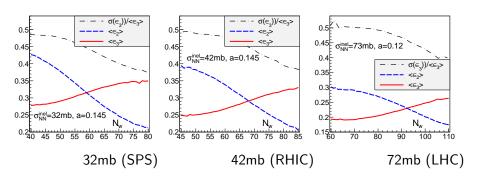
unclustered

#### Clusters:

Anticorrelation:  $\rho(\epsilon_2, \epsilon_3) \simeq -0.3$ 

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## Dependence on the collision energy



Qualitative conclusions remain from SPS to the LHC

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#### Conclusions

Signatures of clustered  $^{12}\text{C-}^{208}\text{Pb}$  collisions  $\rightarrow$ 

- $\bullet$  Increase of  $\epsilon_3$  and  $v_3$  with multiplicity for the highest multiplicity events
- ullet Decrease of scaled variance  $\epsilon_3$  and  $v_3$  with multiplicity for the highest multiplicity events
- Anticorrelation of  $\epsilon_2$  and  $\epsilon_3$ , or  $v_2$  and  $v_3$

#### Extensions:

- Other systems
- More detailed modeling

Possible future data (NA61?) in conjunction with a detailed knowledge of the dynamics of the evolution of the fireball would allow to place constrains on the  $\alpha$ -cluster structure of the colliding nuclei. Conversely, the knowledge of the clustered nuclear distributions may help

Conversely, the knowledge of the clustered nuclear distributions may help to verify the fireball evolution models