

Ultrarelativistic proton-beryllium collisions

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Message:

Nuclear structure effects can be seen in ultrarelativistic collisions

Light nucleus “hitting a wall”

[WB + Enrique Ruiz Arriola, PRL 112 (2014) 112501,
PB+WB+ERA+MR, PRC 90 (2014) 064902]

deformation in light nuclei



harmonic flow in collisions with a heavy nucleus

[see talks by P. Bożek and M. Rybczyński]

This talk: p – polarized A collisions

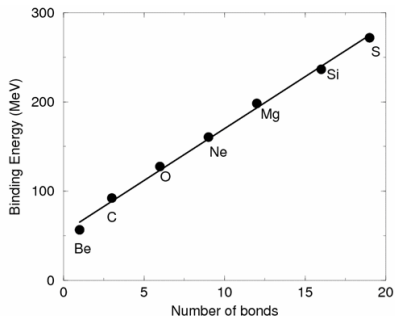
- Effects of clustering in light nuclei can be seen in a very robust way in distributions of number of participants
- Polarization needed (spin in the ground state + magnetic field)

Isotopes of beryllium

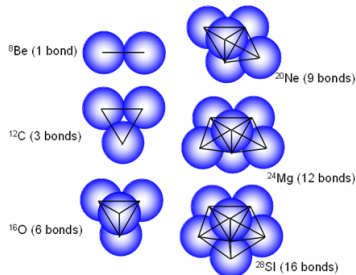
symbol	Z	N	lifetime	decay	J^P
${}^7\text{Be}$	4	3	53 days	e capture	$3/2^-$
${}^8\text{Be}$	4	4	7×10^{-17} s	α	0^+
${}^9\text{Be}$	4	5	stable		$3/2^-$

Clusters in light nuclei

Some features of clustering



- strong binding
- small overlap
- → can treat approximately as elementary
- specific excitation spectra
- fragmentation experiments
- **Generalization:** ${}^7\text{Be} = {}^4\text{He} + {}^3\text{He}$,
 ${}^9\text{Be} = 2{}^4\text{He} + n$



[see C. Beck ed., *Clusters in Nuclei*,
Lecture Notes in Physics 818, 848, 875,
Springer (2010, 2012, 2014)]

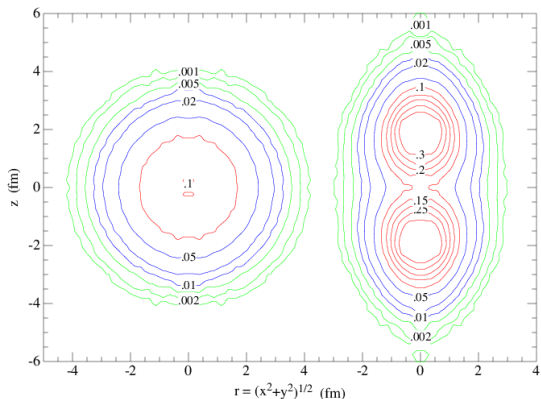


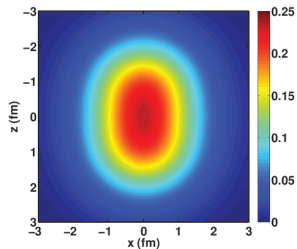
Fig. 15 (Wiringa, et al.)

laboratory frame

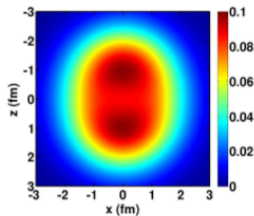
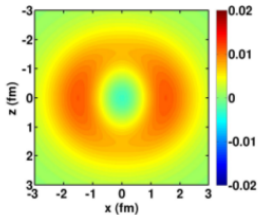
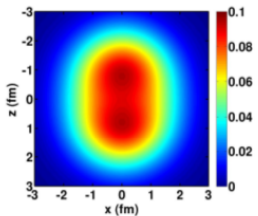
intrinsic frame

Contours of constant density for ^8Be , plotted in cylindrical coordinates.
 [R. Wiringa et al., PRC 62 (2000) 014001]

No-core shell model ${}^7\text{Be}$ and ${}^9\text{Be}$



${}^7\text{Be} (p + n)$



$${}^9\text{Be} \begin{pmatrix} p & n \\ n - p & n \end{pmatrix}$$

[Robert Chase Cockrell, PhD Thesis, Iowa State U.]

Making states with good quantum numbers

Simplest case: ${}^8\text{Be}$. GS is a 0^+ state (rotationally symmetric w.f.).

Deformation concerns **multiparticle correlations** between the nucleons

Superposition over orientations:

$$|\Psi_{0^+}(x_1, \dots, x_8)\rangle = \frac{1}{4\pi} \int d\Omega \Psi_{\text{intr}}(x_1, \dots, x_8; \Omega)$$

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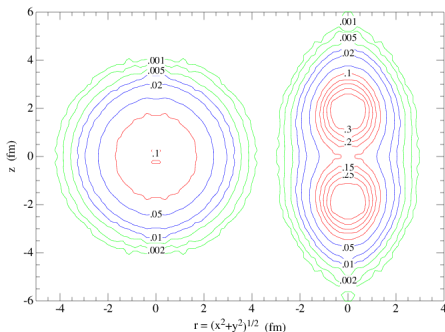


Fig. 15 (Wiringa, et al.)

Making ${}^7\text{Be}$ with good quantum numbers

$${}^7\text{Be} = {}^4\text{He} + {}^3\text{He} \quad (\text{treated as elementary})$$

$$\frac{3}{2}^- = 0^+ + \frac{1}{2}^+ + 1^- \quad (\text{orbital motion of } {}^4\text{He} \text{ and } {}^3\text{He})$$

$$|\frac{3}{2}, m = \frac{3}{2}\rangle = |\frac{1}{2}, \frac{1}{2}\rangle \otimes |1, 1\rangle$$

$$|\frac{3}{2}, m = \frac{1}{2}\rangle = \sqrt{\frac{2}{3}}|\frac{1}{2}, \frac{1}{2}\rangle \otimes |1, 0\rangle + \sqrt{\frac{1}{3}}|\frac{1}{2}, -\frac{1}{2}\rangle \otimes |1, 1\rangle$$

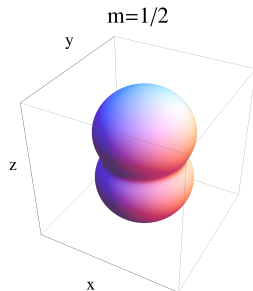
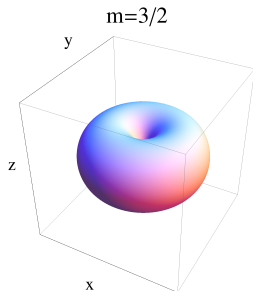
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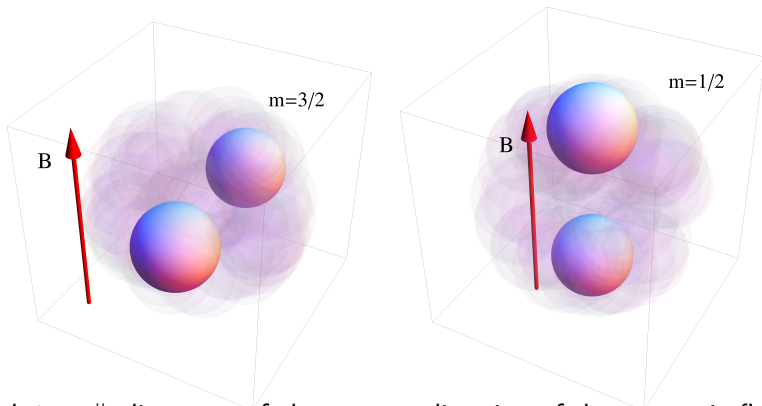


spatial density of orbitals $|\langle \frac{3}{2}, m | \frac{3}{2}, m \rangle|^2$

Overlaying the intrinsic distributions

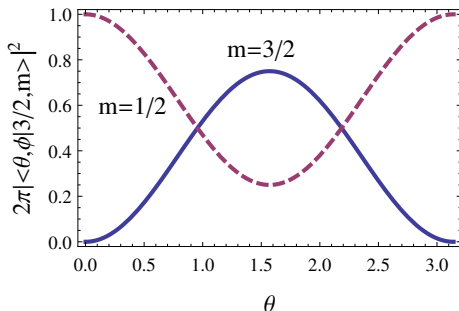
For nuclear physicists: equivalent to the Peierls-Yoccoz projection

$$\Psi_{3/2,m}(\vec{r}) = \sum_{k=\pm\frac{1}{2}} \int d\Omega D_{m,k}^{3/2}(\Omega) \Psi_k^{\text{intr}}(\vec{r}; \Omega)$$



(partial \perp or \parallel alignment of clusters wrt direction of the magnetic field)

Distribution in the axial angle



$m = 3/2$ peaks at the equator (\perp), and $m = 1/2$ at the poles (\parallel)

${}^9\text{Be}$ with good quantum numbers

$${}^9\text{Be} = n + {}^4\text{He} + {}^4\text{He}$$

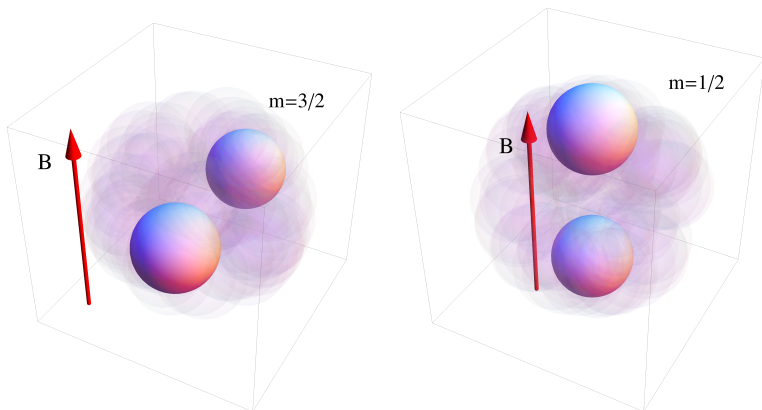
$$\frac{3}{2}^- = \frac{1}{2}^+ + 0^+ + 0^+ + 1^- \text{ (orbital motion of } {}^4\text{He and } {}^4\text{He)}$$

– then the projection is exactly as for ${}^7\text{Be}$

p – polarized Be collisions

Why ultra-relativistic?

Reaction time is much shorter than time scales of the structure
→ a frozen “snapshot” of the nuclear configuration



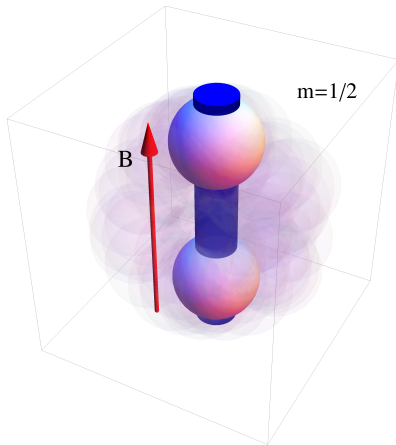
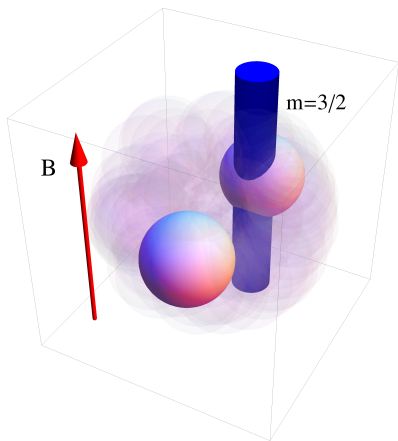
Probability of inelastic p-A interaction at impact parameter \vec{b} :

$$P_{pA}(\vec{b}) = \prod_{i=1}^A \int d^3x_i |\psi_A(\vec{x}_1, \dots, \vec{x}_A)|^2 \\ \times \frac{1}{\sigma_{pA}^{\text{inel}}(s)} \left\{ 1 - \prod_{i=1}^A \left[1 - \sigma_{NN}^{\text{inel}} P_{NN}(\vec{b} - \vec{x}_{i,T}) \right] \right\}$$

(correlated) GS wave function enters

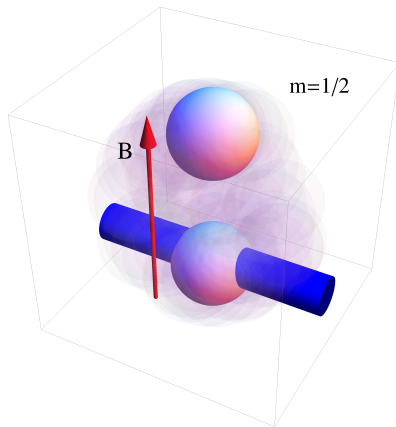
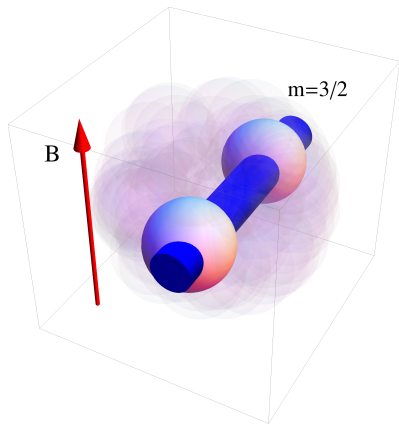
Simulations with GLISSANDO 2, $\sigma_{NN}^{\text{inel}} = 32 \text{ mb}$ (SPS)

$$B \parallel z$$

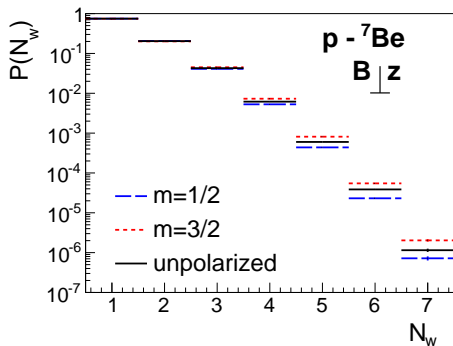
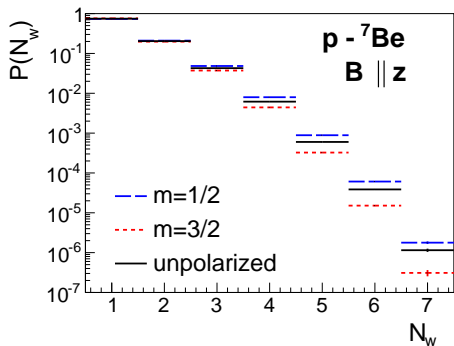


(radius of the cylinder corresponds to the NN wounding cross section)

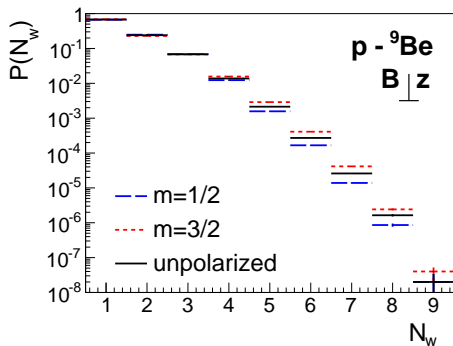
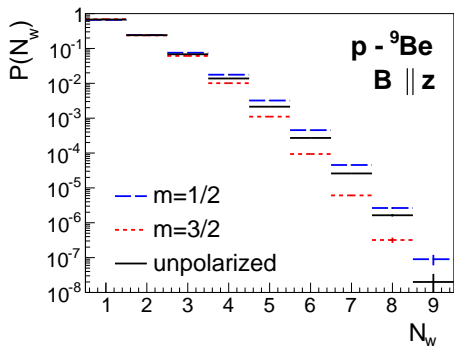
$$B \perp z$$



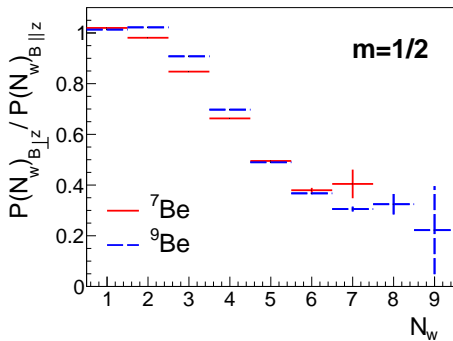
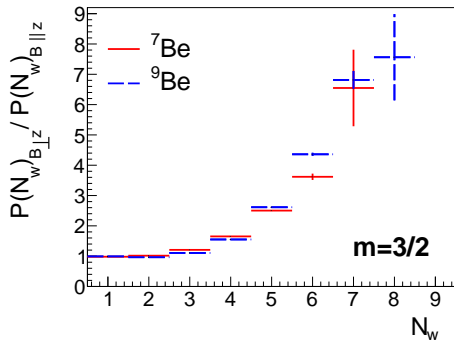
Results for $p - {}^7\text{Be}$



Results for $p - {}^9\text{Be}$



Ratios



Conclusions

Light nuclei structure from relativistic collisions with protons

- High-energy protons probe the nuclear wave function

Needed: magnetic field → sensitivity to orientation

- Factor-of-a-few effects in the distribution for large N_w
- Need to distinguish m or the orientation wrt magnetic field
- Sensitivity to the nuclear structure: clustering gives the strong effects
- Applicable to nuclei with nonzero spin in GS
- No polarization of Be → small effects (clustered vs uniform: $\sim 50\%$)
 - MSc Theses by Milena Sołtysiak, UJK (2014)
- No polarization in Be-Be collisions → small effects

Experimental prospects? There is magnetic field in beam lines ...