

Forward-backward flow correlations in relativistic heavy-ion collisions¹

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¹Based on P. Bozek, WB, J. Moreira, PRC 83, 034911 (2011)

Outline

- 1 Introduction
- 2 Collision geometry
 - Emission profiles from d -Au
 - Emission profiles in AA
- 3 Fluctuations
 - Initial-state fluctuations in Glauber
- 4 Torque
 - Torque from fluctuations and asymmetric emission
 - Measures of the torque
 - Cumulants

Basic idea

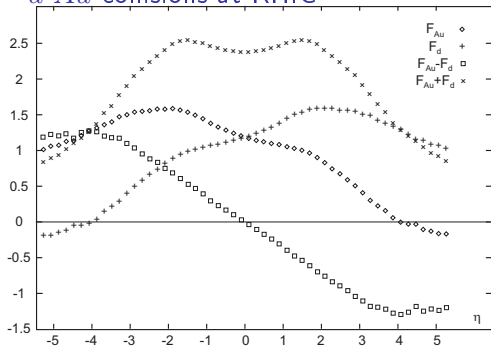
- determination of the initial state is crucial [review by S. Bass]
- event-by-event fluctuations

These elements jointly lead to interesting new effects, in particular to **forward-backward flow correlations**

Collision geometry

[Bialas and Czyz, Acta Phys. Polon. B36, 905 (2005)]

d-Au collisions at RHIC

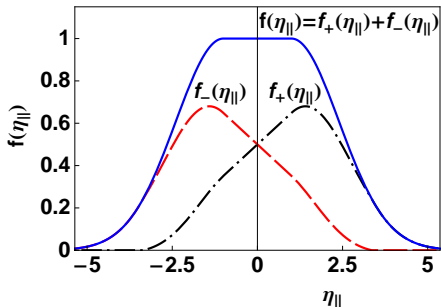


Wounded nucleons emit
 predominantly in their
 forward hemisphere

[similar ideas: Brodsky, Gunion, Kuhn (1977); Adil, Gyulassy (2005); Adil, Gyulassy, Hirano (2006)]

Extension to AA collisions

[Gazdzicki, Gorenstein (2006); Bzdak (2009); Bzdak, Wozniak (2010); Bozek, Wyskiel, PRC 81 (2010) 054902]



mixed model: three sources

$$N_{bin}, N_w^+, N_w^-$$

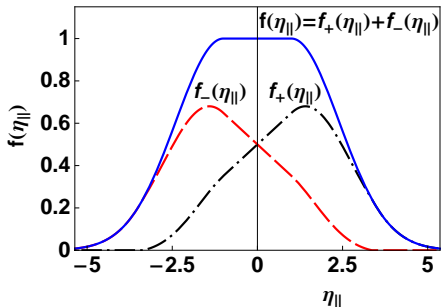
$$N_{prod} \sim \frac{1-\alpha}{2} N_w + \alpha N_{bin}$$

$\alpha \sim 14\%$ at RHIC

explains the F-B multiplicity
 fluctuations (Bzdak+Wozniak)

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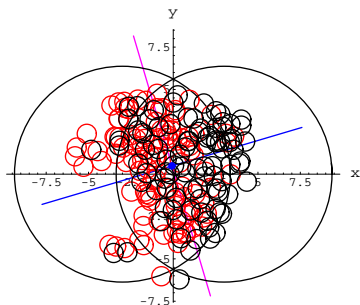
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Key formula: space-time rapidity ($\eta_{||}$) and transverse (x, y) distribution:

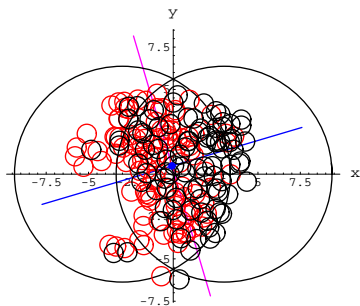
$$F(\eta_{||}, x, y) = (1-\alpha)[\rho_+(x, y)f_+(\eta_{||}) + \rho_-(x, y)f_-(\eta_{||})] + \alpha\rho_{bin}(x, y)f(\eta_{||})$$

Fluctuations in the Glauber approach



- shape and size fluctuations originate from the statistical nature of the distribution of sources in the transverse plane
- increased eccentricity \rightarrow hydro \rightarrow increased v_2
- e-by-e fluctuations of v_2
- triangular flow, higher Fourier components, the ridge

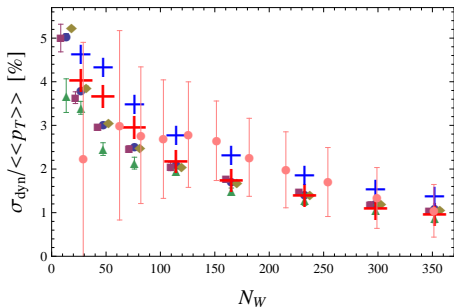
Fluctuations in the Glauber approach



clusters with $N_W^A \neq N_W^B$!

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Digression: p_T fluctuations



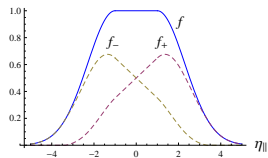
radial size fluctuations \rightarrow
 hydro \rightarrow radial flow
 fluctuations \rightarrow explanation
 of p_T fluctuations at RHIC

GLISSANDO + 3+1 perfect hydro
 + THERMINATOR

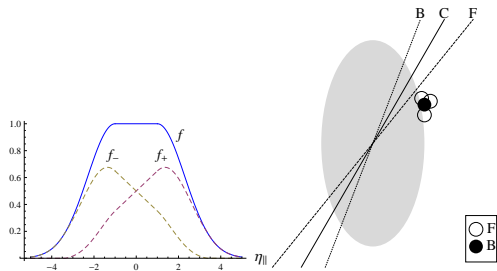
mixed model, wounded-nucleon model, STAR and PHENIX data

[WB, Chojnacki, Obara, PRC 80 (2009) 051902(R)]

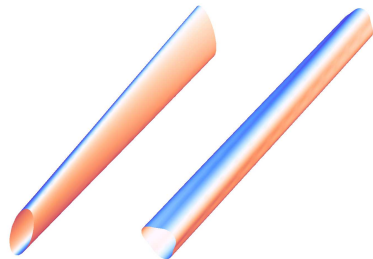
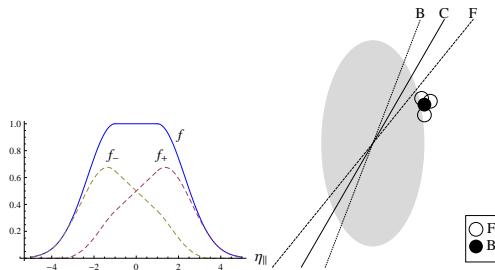
The torque



The torque



The torque

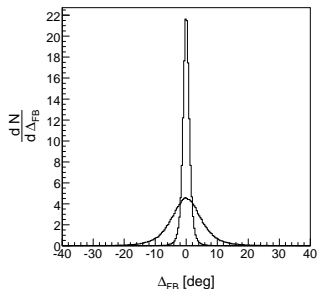


e-by-e fluctuation of the relative angle of the F and B principal axes

Distribution of the torque angle

GLISSANDO simulations in the mixed model

Δ_{FB} - torque angle between the forward and backward rapidity regions

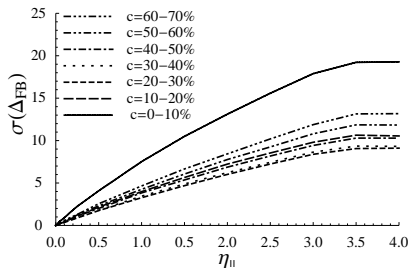
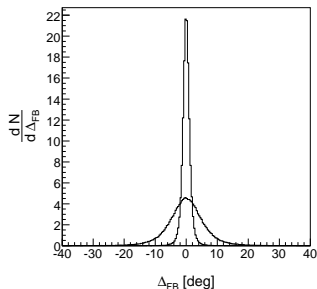


narrow: $\eta_{||} = 0.5$, broad: $\eta_{||} = 2.5$
 ($c = 20 - 30\%$)

Distribution of the torque angle

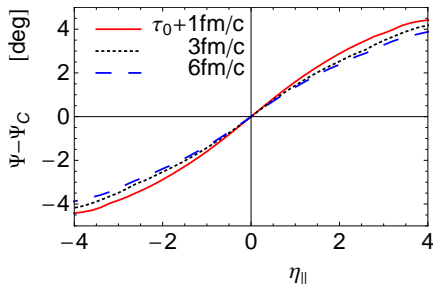
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Hydrodynamic evolution



- hydro (here perfect 3+1 with realistic EoS) evolution \rightarrow the torque survives
- statistical hadronization leads to additional fluctuations - can we sort out the effect?

Measures of the torque based on cumulants

$$\langle e^{i2(\phi_F - \phi_B)} \rangle = \frac{1}{N_{\text{events}}} \sum_{\text{events}} \frac{1}{n_F n_B} \sum_{i=1}^{n_F} \sum_{j=1}^{n_B} e^{i2(\phi_i - \phi_j)}$$

When no correlations, then

$$f(\phi) = v_0 + 2 \sum_{k=1} v_k \cos[k(\phi - \Psi^{(k)})]$$

and

$$\langle e^{i2(\phi_F - \phi_B)} \rangle = \langle v_{2,F} v_{2,B} \cos(2\Delta_{FB}) \rangle_{\text{events}}$$

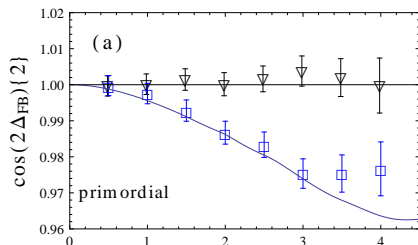
Since we are interested in measuring the average $\cos[2(\Psi_F - \Psi_B)]$, we need to divide by $v_{2,F} v_{k,B}$. This can be achieved by considering the ratio

$$\cos(2\Delta_{FB}) \{2\} \equiv \frac{\langle e^{i2(\phi_F - \phi_B)} \rangle}{\sqrt{\langle e^{i2(\phi_{F,1} - \phi_{F,2})} \rangle \langle e^{i2(\phi_{B,1} - \phi_{B,2})} \rangle}} = \langle \cos(2\Delta_{FB}) \rangle_{\text{events}} + \text{nonflow}$$

(similar for higher Fourier components, also more-particle cumulants may be used)

Simulations with THERMINATOR

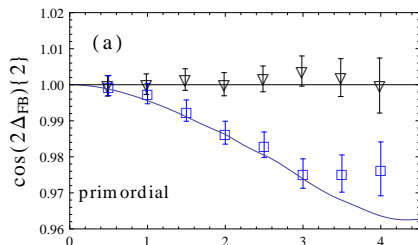
triangles - no torque, squares - with torque (100 000 events)



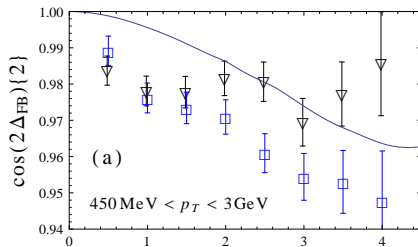
only the primordial particles created at freeze-out included
 (test)

Simulations with THERMINATOR

triangles - no torque, squares - with torque (100 000 events)



only the primordial particles created at freeze-out included
 (test)



all final particles (with resonance decays)
 effect is observable !

Summary

- Rapidity dependent reaction plane from simple assumptions: asymmetric emission profiles in rapidity + fluctuations \rightarrow torque
- $\sim 20^\circ$ torque angle between F and B regions in most central and $\sim 10^\circ$ in mid-peripheral Au+Au collisions at the highest RHIC energy
- Similar for higher-order Fourier components
- Torqued fireball \rightarrow torqued collective flow \rightarrow torqued principal axes of the p_T distributions at different rapidities
- Signal detectable through the use of cumulants involving F and B particles
- With the fantastically high statistics available in experiments the effect can be examined in experiments, but **large rapidity coverage needed**
- The effect influences the elliptic and directed flow studies, which use the reaction planes determined in different pseudorapidity intervals
- Because of the increase in the statistical noise, one should look for the torque effects in mid-central/mid-peripheral collisions ($v_2\sqrt{n}$ should be maximal)

- The effects momentum, angular-momentum, and charge conservation should be examined
- The torque fluctuations in other multi-source models should be investigated
- The experimental detection of the torque would provide independent information on the initial state, in particular would confirm the assumptions on the initial emission profile from the sources

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Please, analyze the data!