

# STRANGE PARTICLE PRODUCTION IN A SINGLE-FREEZE-OUT MODEL

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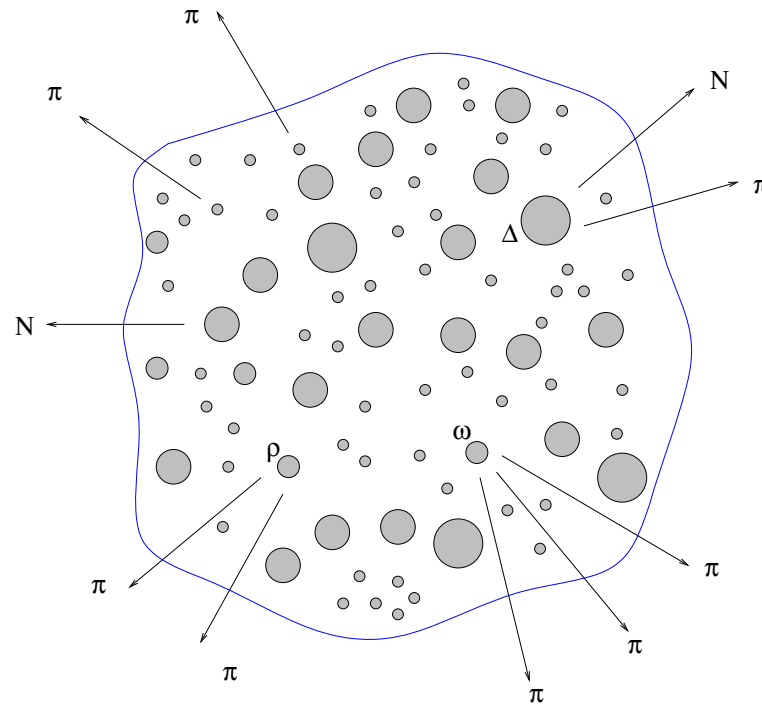
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Cape Town, SOUTH AFRICA, 2004

main topic of this talk: transverse-momentum spectra and elliptic flow of  
multistrange particles produced at RHIC

# Thermal model

Koppe (1948), Fermi (1950), Landau, Hagedorn, Rafelski, Letessier, Torrieri, Bjorken, Gorenstein, Gaździcki, Sinyukov, Kostyuk, Heinz, Sollfrank, Braun-Munzinger, Stachel, Redlich, Csörgő, Lörstad, Becattini, Cleymans, Wheaton



$$\sim e^{-(E-\mu)/T}$$

## our variant

WB + WF, **PRL 87 (2001) 272302** ( $p_{\perp}$  spectra of pions, kaons, and protons)

WB + WF, **PRC 65 (2002) 064905** ( $p_{\perp}$  spectra of strange particles)

WB + WF + Anna Baran, **AIP 660** (HBT radii and  $v_2$ )

WB + WF + Brigitte Hiller, **PRC 68 (2003) 034911** (pion invariant-mass distributions)

Piotr Bożek + WB + WF, **Heavy-Ion Physics (2004)** (pion balance functions)

## single freeze-out model

1.  $T_{\text{chem}} = T_{\text{kin}} \equiv T$
2. Complete treatment of resonances
3. Special choice of the freeze-out hypersurface,  $\tau = \sqrt{t^2 - x^2 - y^2 - z^2} = \text{const}$
4. Only 4 parameters:  $T, \mu_B$  (fixed by the ratios of the particle abundances), invariant time at freeze-out  $\tau$  (controls the overall normalization), transverse size  $\rho_{\text{max}}$  ( $\rho_{\text{max}}/\tau$  controls the slopes of the  $p_{\perp}$  spectra)
5. Hubble-like flow,  $u^{\mu} = \frac{x^{\mu}}{\tau} = \frac{t}{\tau}(1, \frac{x}{t}, \frac{y}{t}, \frac{z}{t})$

definition of the hypersurface not unique! – see papers by Rafelski and Torrieri

the choice  $t = \text{const}$ , á la Blast-Wave, also possible

## Ratios → Transverse-momentum spectra

initial step: standard analysis of the particle ratios gives  $T$  and  $\mu_B$  (relative normalization), later  $\tau$  and  $\rho_{\max}$  fitted from the spectra (absolute normalization and shape)

recent developments: **SHARE** - Statistical **H**adronization with **R**esonances, Cracow - Arizona Collaboration supported by the NATO grant, nucl-th/0404083, submitted to Communications in Physics Computing

G. Torrieri, W. Broniowski, J. Letessier, J. Rafelski, and WF

a set of programs in Fortran and Mathematica devoted to the statistical analysis of the ratios of hadron resonances, offered for the community via our web page:

[www.physics.arizona.edu/~torrieri/share/share.html](http://www.physics.arizona.edu/~torrieri/share/share.html)

alternative approach: **THERMUS** - a thermal model package for ROOT

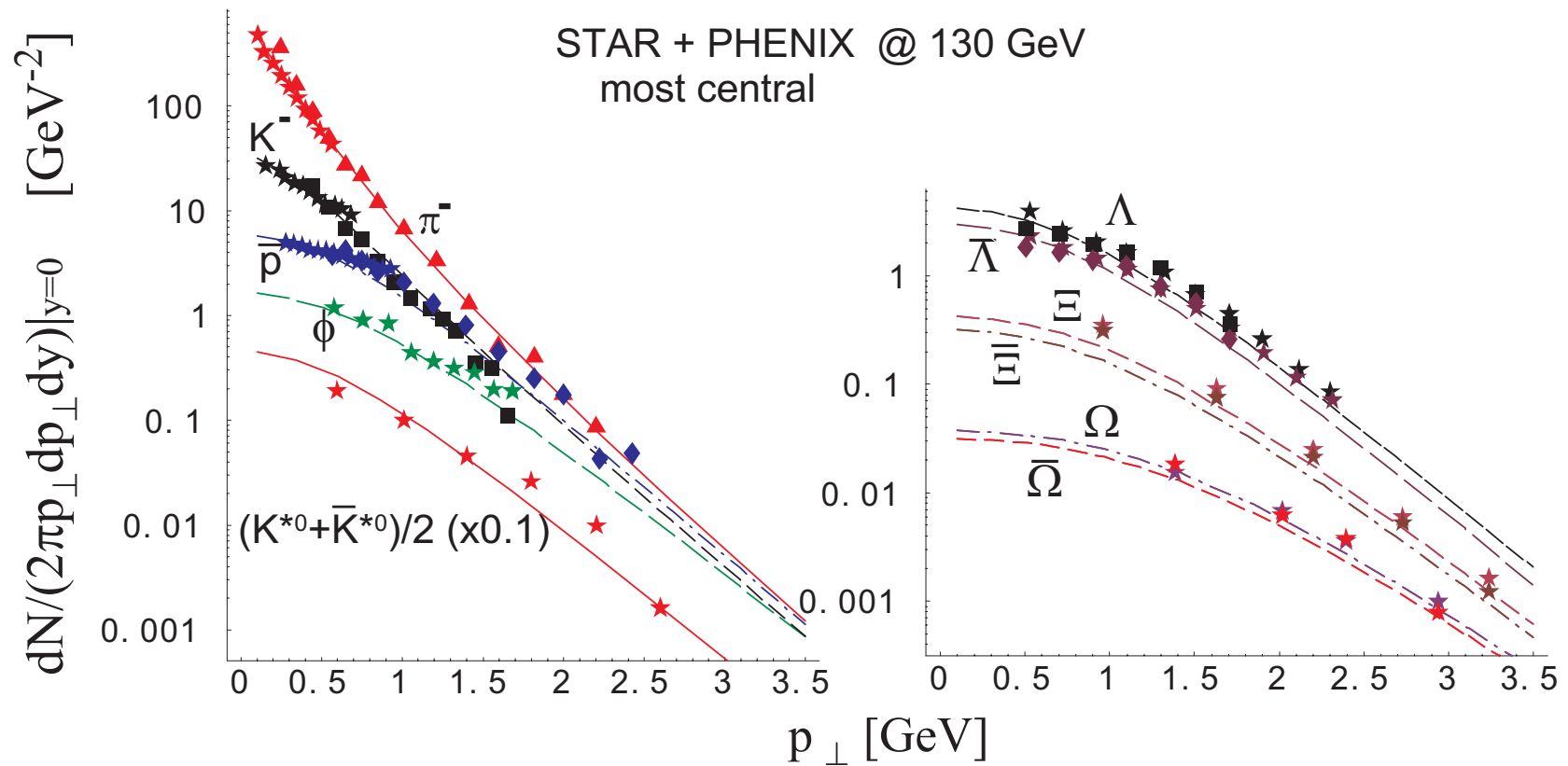
S. Wheaton and J. Cleymans, hep-ph/0407174

## 2 thermal parameters fitted from particle ratios

$$T \text{ [MeV]} = 165 \pm 7, \mu_B \text{ [MeV]} = 41 \pm 5, \quad @ 130 \text{ GeV}$$

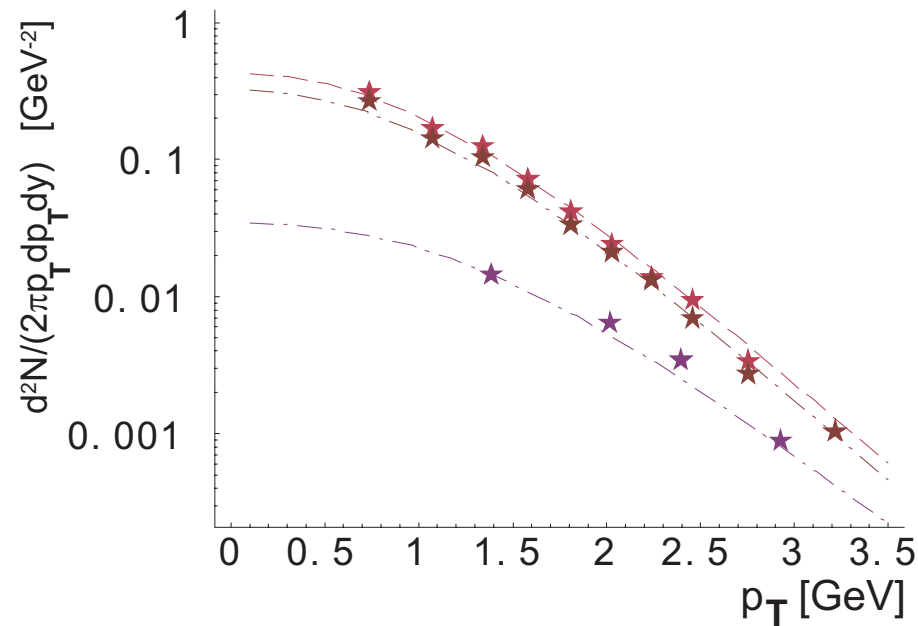
$$T \text{ [MeV]} = 166 \pm 5, \mu_B \text{ [MeV]} = 29 \pm 4, \quad @ 200 \text{ GeV}$$

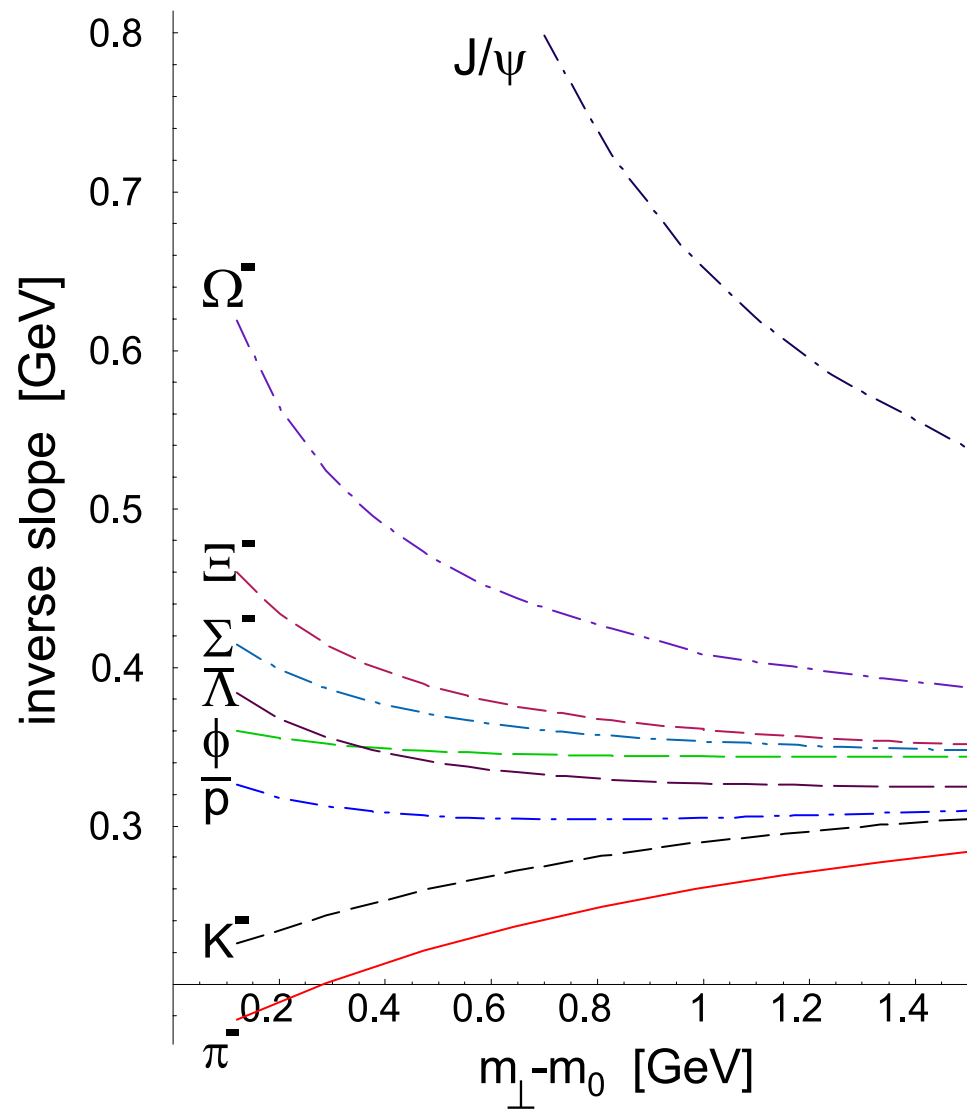
## 2 geometric parameters fitted from the spectra of $\pi^\pm, K^\pm, p,$ and $\bar{p}$



comparison of our predictions from (2001) with the finally published data, STAR Collaboration, Multi-Strange Baryon Production in Au-Au Collisions at  $\sqrt{s_{NN}} = 130$  GeV, J. Adams et al., Phys. Rev. Lett. 92, 182301 (2004)

### spectra of $\Xi$ and $\Omega + \bar{\Omega}$

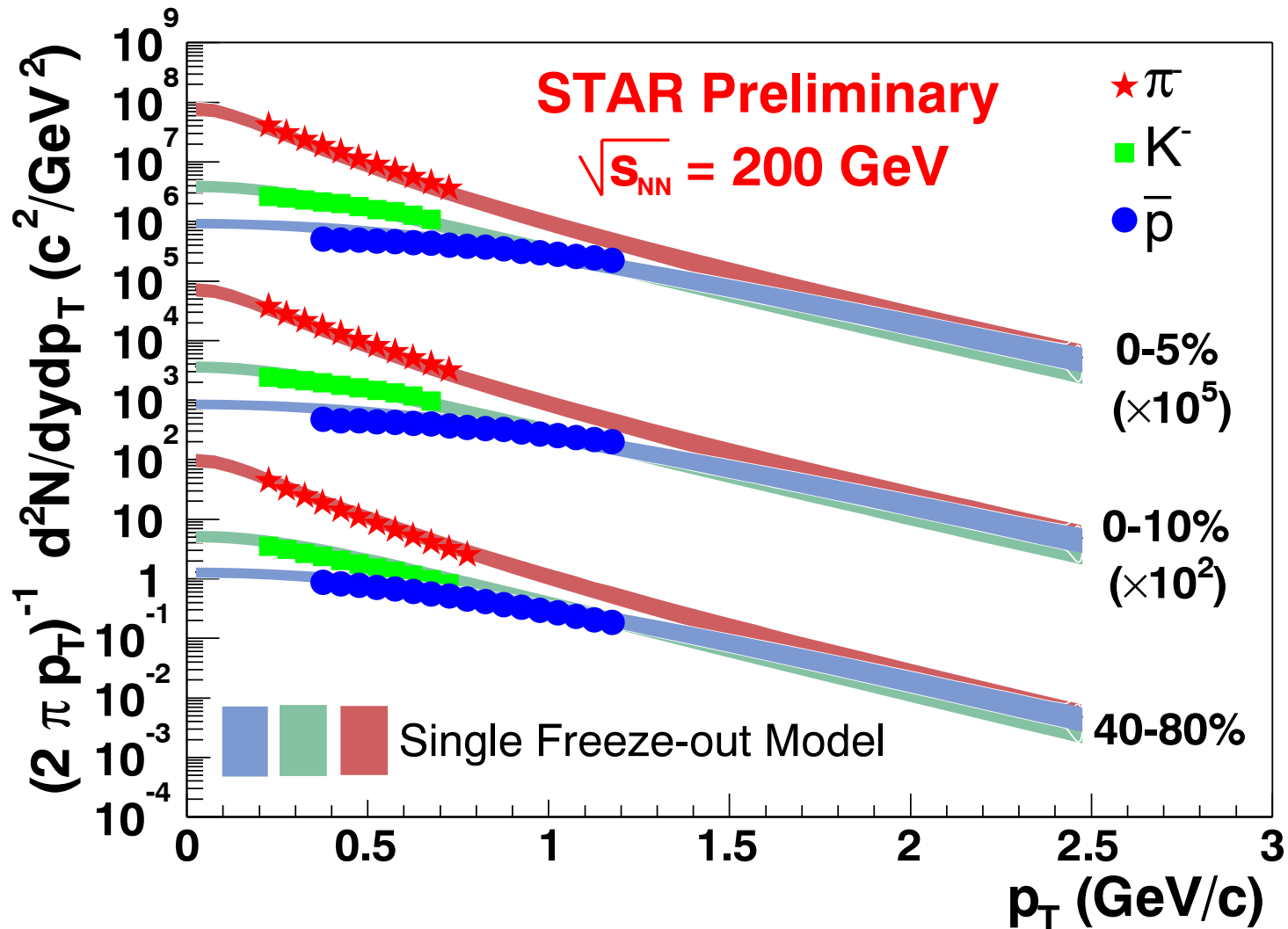




**130  $\longrightarrow$  200 GeV**

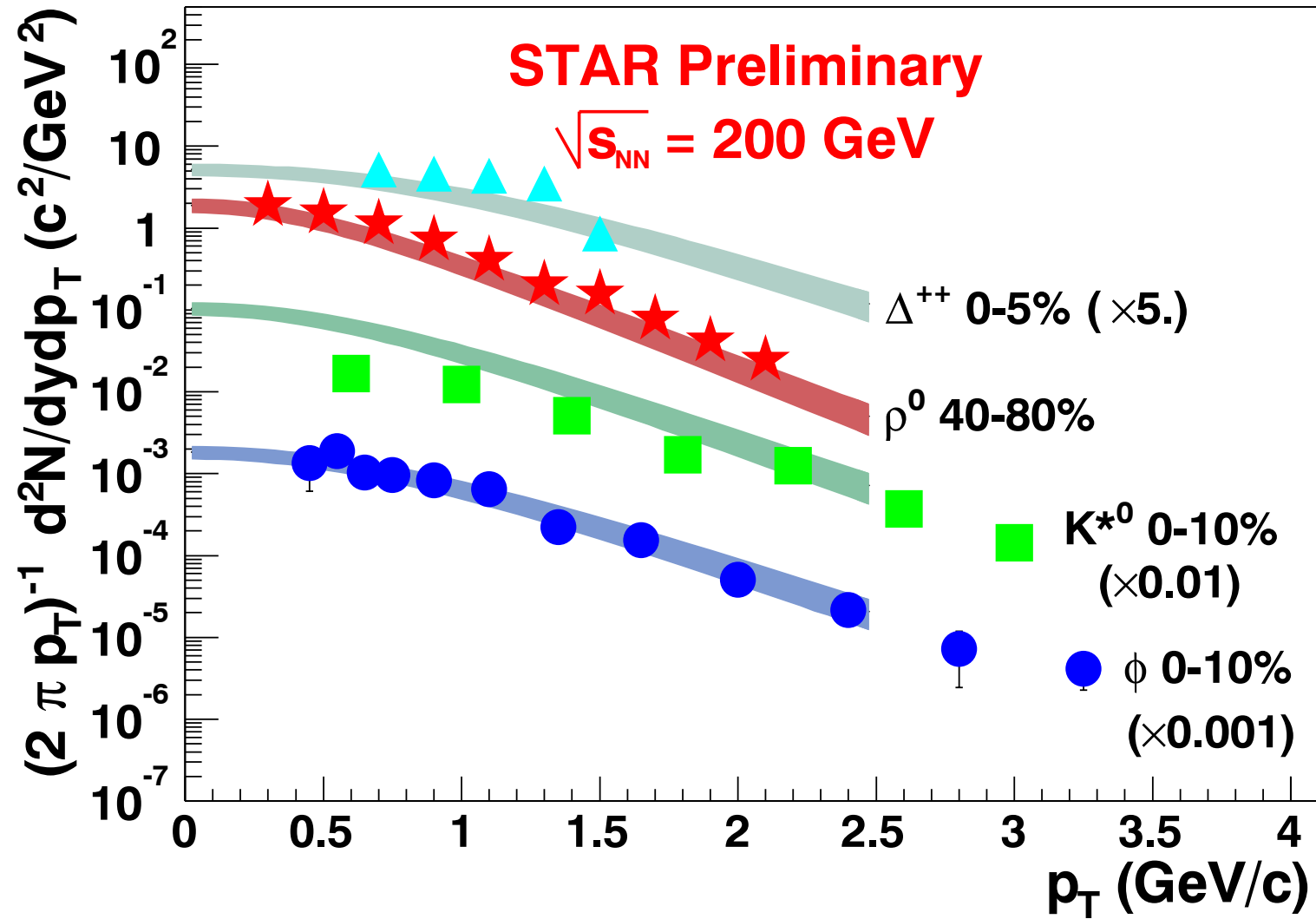
# STAR spectra vs. single freeze-out model

compiled by Patricia Fachini





compiled by Patricia Fachini

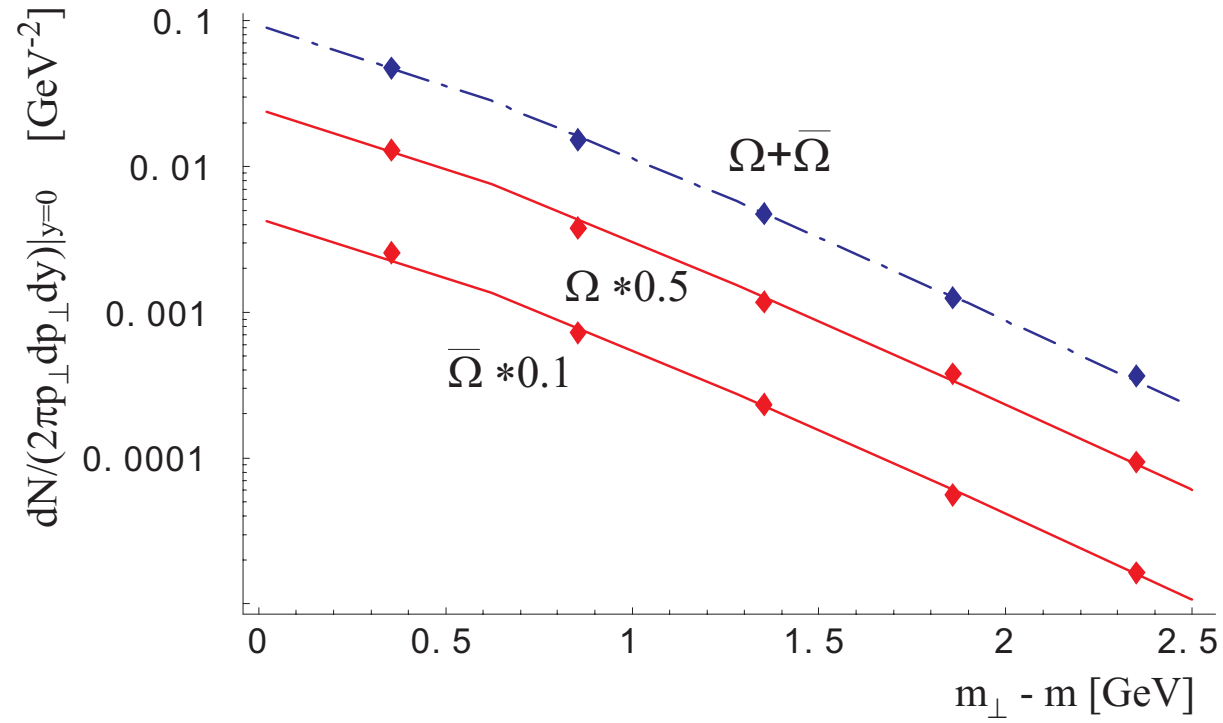


## Ratios including resonances

	$m_\rho^* = 770 \text{ MeV}$	$m_\rho^* = 700 \text{ MeV}$	Experiment
$T \text{ [MeV]}$	$T = 165.6 \pm 4.5$	$T = 167.6 \pm 4.6$	
$\mu_B \text{ [MeV]}$	$\mu_B = 28.5 \pm 3.7$	$\mu_B = 28.9 \pm 3.8$	
$\eta/\pi^-$	$0.120 \pm 0.001$	$0.112 \pm 0.001$	
$\rho^0/\pi^-$	$0.114 \pm 0.002$	$0.135 \pm 0.001$	$0.183 \pm 0.028 \text{ (40-80\%)}$
$\omega/\pi^-$	$0.108 \pm 0.002$	$0.102 \pm 0.002$	
$K^*(892)/\pi^-$	$0.057 \pm 0.002$	$0.054 \pm 0.002$	
$\phi/\pi^-$	$0.025 \pm 0.001$	$0.024 \pm 0.001$	
$\eta'/\pi^-$	$0.0121 \pm 0.0004$	$0.0115 \pm 0.0003$	
$f_0(980)/\pi^-$	$0.0102 \pm 0.0003$	$0.0097 \pm 0.0003$	$0.042 \pm 0.021 \text{ (40-80\%)}$
$K^*(892)/K^-$	$0.33 \pm 0.01$	$0.33 \pm 0.01$	$0.205 \pm 0.033 \text{ (0-10\%)}$ $0.219 \pm 0.040 \text{ (10-30\%)}$ $0.255 \pm 0.046 \text{ (30-50\%)}$ $0.269 \pm 0.047 \text{ (50-80\%)}$
$\Lambda(1520)/\Lambda$	$0.061 \pm 0.002$	$0.062 \pm 0.002$	$0.022 \pm 0.010 \text{ (0-7\%)}$ $0.025 \pm 0.021 \text{ (40-60\%)}$ $0.062 \pm 0.027 \text{ (60-80\%)}$
$\Sigma(1385)/\Sigma$	$0.484 \pm 0.004$	$0.485 \pm 0.004$	

see talk about the resonance production by Christina Markert on Sunday

# $\Omega$ spectra from STAR @ 200 GeV



C. Suire, QM2002

## Elliptic flow

WB + WF + Anna Baran, Proceedings of the Coimbra Workshop on Hadron Physics, nucl-th/0212053, AIP 660

Ph. D. Thesis by Anna Baran, to be published

when the nuclei collide at non-zero impact parameter,  $b \neq 0$ , the momentum distribution of the produced particles carries azimuthal asymmetry

$$\left. \frac{dN}{d^2p_{\perp} dy} \right|_{y=0} = \left. \frac{dN}{2\pi p_{\perp} dp_{\perp} dy} \right|_{y=0} (1 + 2 v_2 \cos 2\phi + 2 v_4 \cos 4\phi + \dots)$$

experimentally determined **centrality** may be used to determine  $b$  from the geometric formula

$$c \simeq \frac{b^2}{(2R)^2}$$

**eccentricity** is obtained from the measured values of  $R_{\text{side}}(\phi)$ , STAR Collaboration,

$$\epsilon = \frac{\langle y \rangle^2 - \langle x \rangle^2}{\langle y \rangle^2 + \langle x \rangle^2}$$

modification of the freeze-out hypersurface (almond shape)

$$r_x = \rho_{\max} \sqrt{1 - \epsilon} \cos \phi$$

$$r_y = \rho_{\max} \sqrt{1 + \epsilon} \sin \phi$$

modification of the flow profile (stronger in-plane)

$$u_x = \frac{r_x}{N} \sqrt{1 + \delta} \cos \phi$$

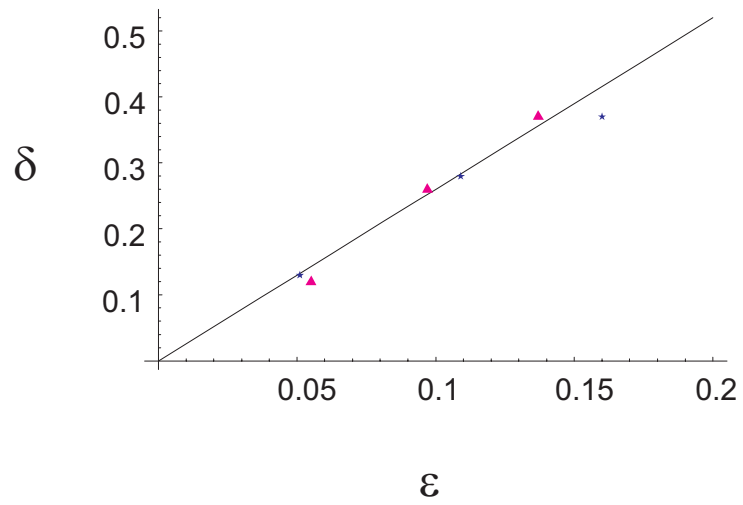
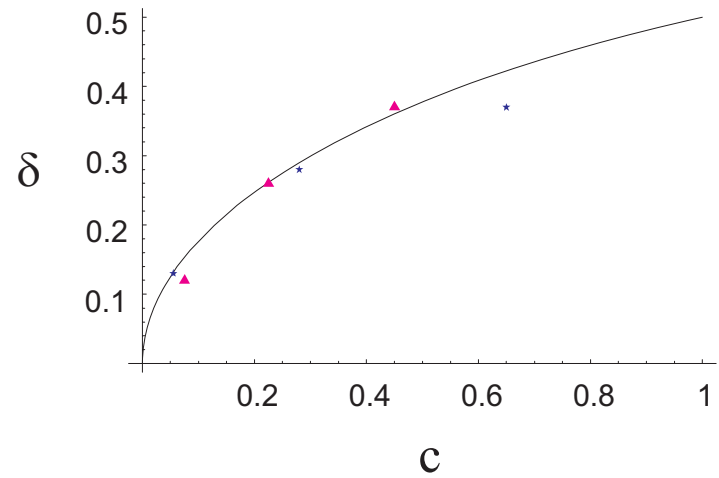
$$u_y = \frac{r_y}{N} \sqrt{1 - \delta} \sin \phi$$

$$u_z = \frac{r_z}{N}$$

$$u_t = \frac{t}{N}$$

$N$  obtained from the normalization condition  $u^\mu u_\mu = 1$

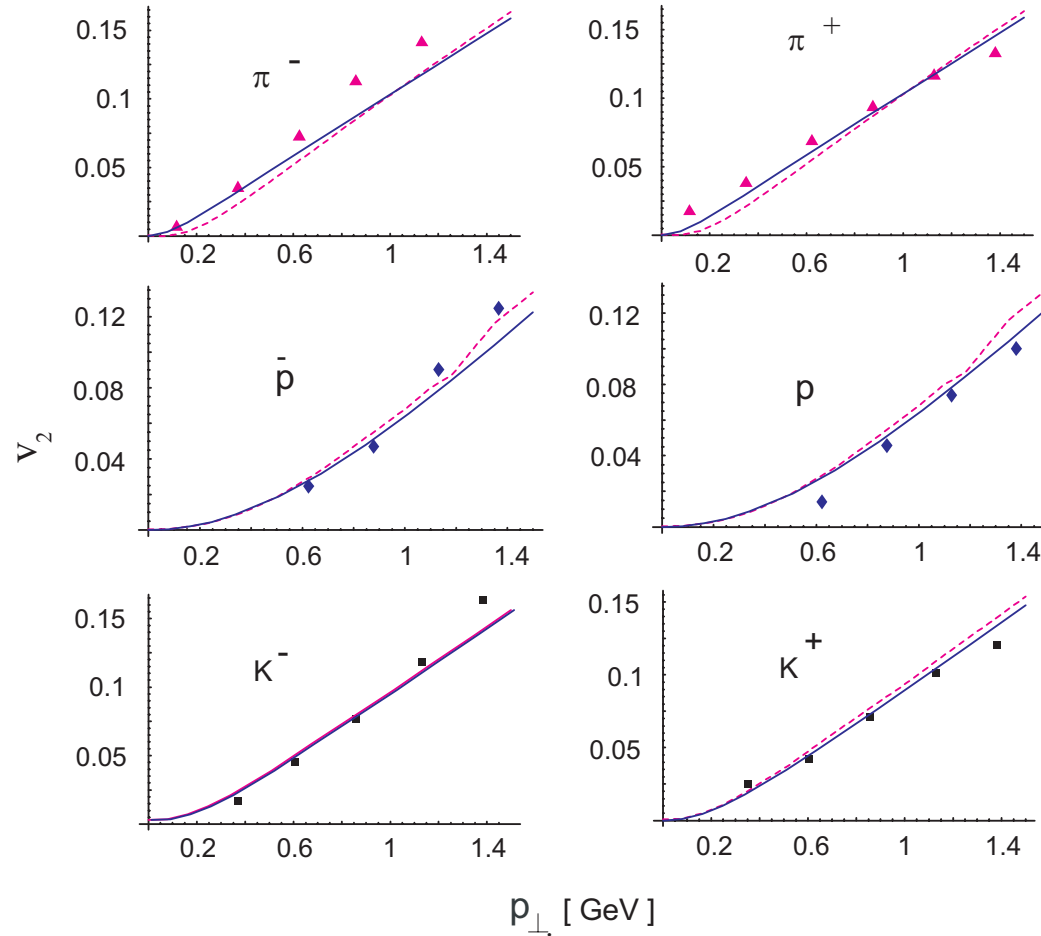
results of the fit procedure of  $v_2$  for pions, kaons, and protons



▲ PHENIX      ★ STAR

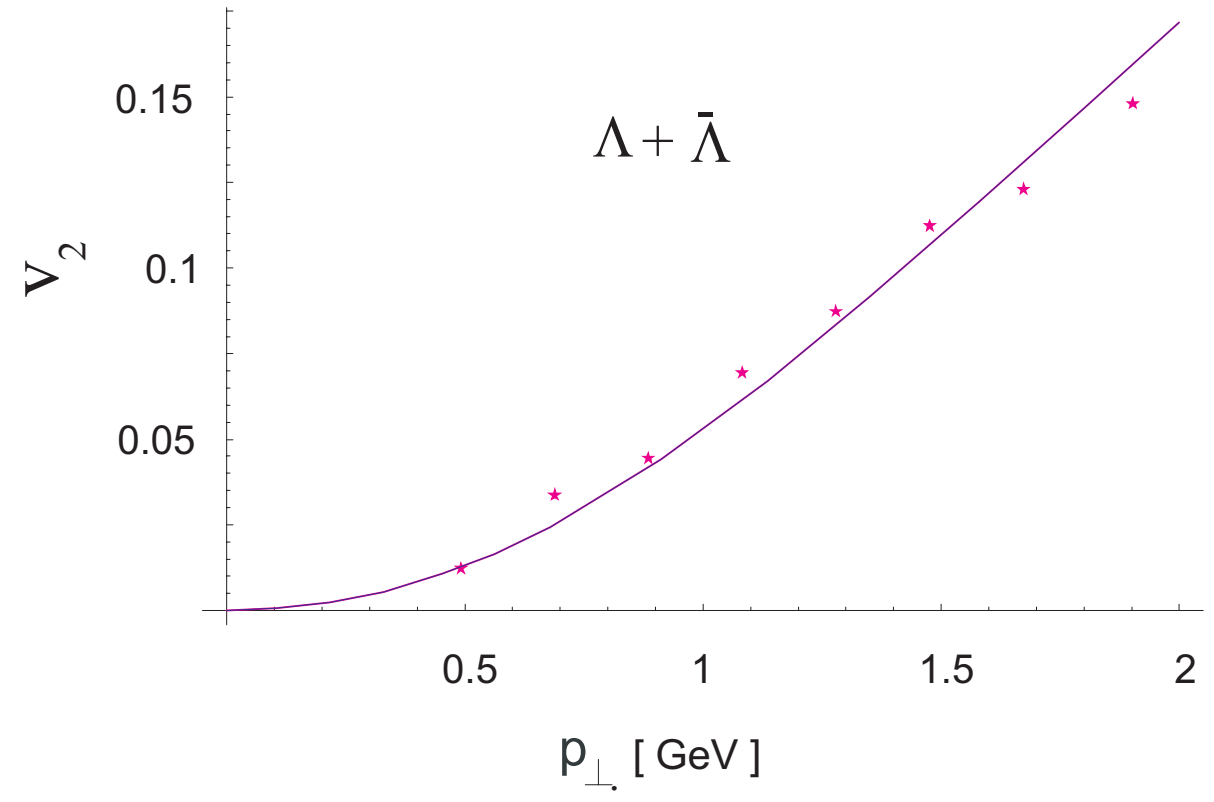
preliminary minimum bias data from PHENIX @ 200 GeV

S. A. Voloshin, Nucl. Phys. **A715** (2003) 379c



single-freeze-out model fit:  $T = 165$  MeV,  $\mu_B = 26$  MeV (from the ratios),  $\tau = 4.04$  fm,  $\rho_{\max} = 3.70$  fm (from the spectra),  $\epsilon = 0.13$ ,  $\delta = 0.25$  (from  $v_2$ )

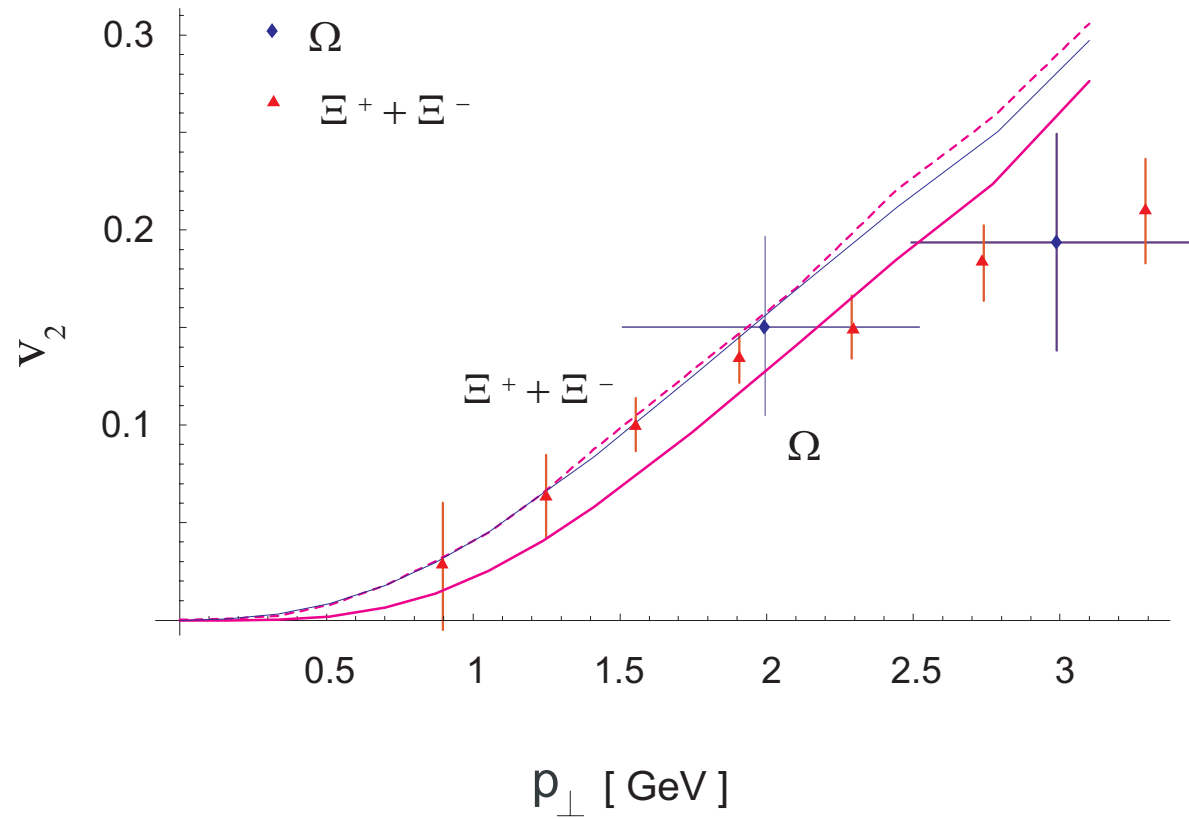
minimum bias (0-80%) data from STAR @ 200 GeV, PRL 92 (2004) 052302



model parameters as above

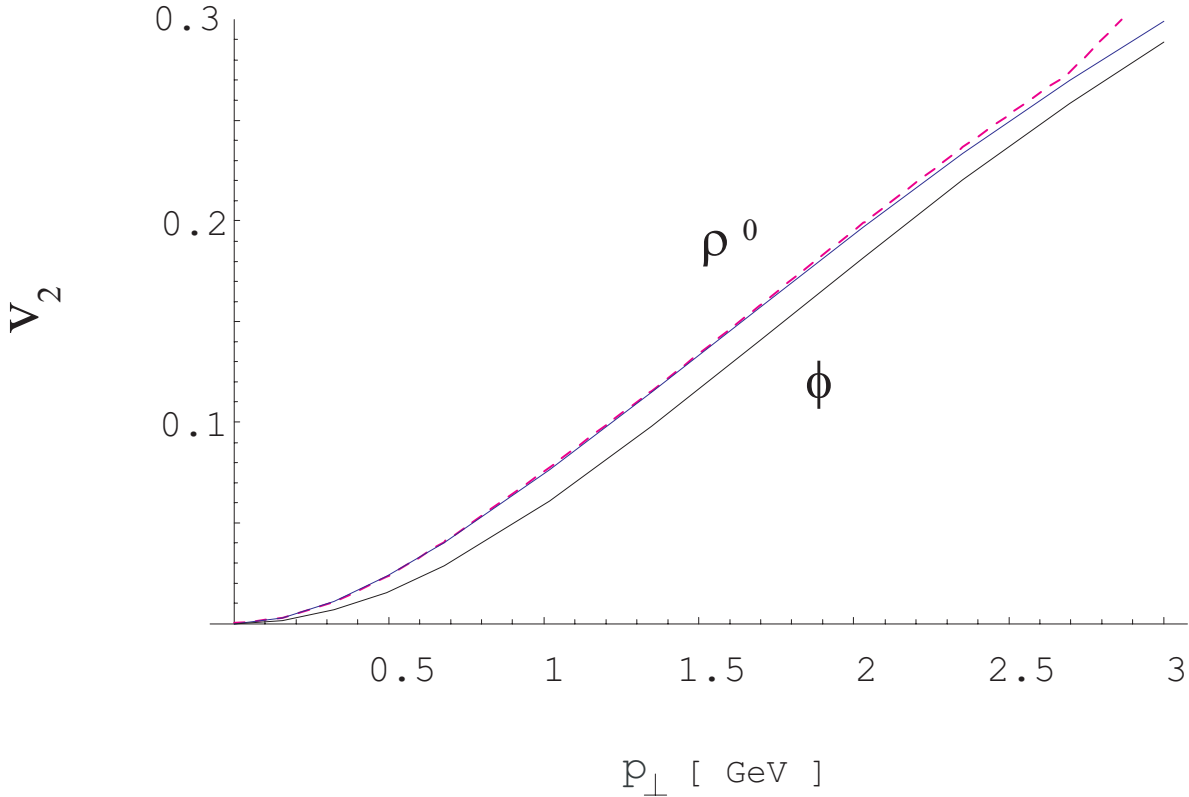


the first measurement of the elliptic flow for multistrange baryons, J. Castillo,  
contribution to QM04, J. Phys. G30 (2004) S1207



model parameters as above

again with the same parameters predictions for  $\rho$  and  $\phi$



# Summary

1. Production of strange particles is well described in a thermal model with single freeze-out
2. Thermodynamic parameters are determined from the ratios of hadron abundances, whereas the expansion parameters are determined from the spectra of pions, kaons, and protons. Then, the predictions about the transverse-momentum spectra and  $v_2$  are made for strange particles
3. Single freeze-out model yields  $R_{\text{out}}/R_{\text{side}} \approx 1$ , and describes well pion invariant masses and balance functions

# Back-up slides

# The phase-shift formula for the density of resonances

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Beth,Uhlenbeck (1937); Dashen, Ma, Bernstein, Rajaraman; **Weinhold, Friman, Nörenberg**; WB+WF+BH, PRC 68 (2003) 034911; Pratt, Bauer, nucl-th/0308087

$$\frac{dn}{dM} = f \int \frac{d^3p}{(2\pi)^3} \frac{d\delta_{12}(M)}{\pi dM} \frac{1}{\exp\left(\frac{\sqrt{M^2+p^2}}{T}\right) \pm 1}$$

In some works the spectral function of the resonance is used instead of the derivative of the phase shift. For narrow resonances this does not make a difference, since then  $d\delta_{12}(M)/dM \simeq \pi\delta(M - m_R)$ , and

$$n^{\text{narrow}} = f \int \frac{d^3p}{(2\pi)^3} \frac{1}{\exp\left(\frac{\sqrt{m_R^2+p^2}}{T}\right) \pm 1}$$

For wide resonances, or for effects of tails, the difference between the correct formula and the one with the spectral function is significant

# Concept of the balance functions

S. Bass, P. Danielewicz, and S. Pratt, PRL 85 (2000) 2689

$$B(\delta, Y) = \frac{1}{2} \left\{ \frac{\langle N_{+-}(\delta) \rangle - \langle N_{++}(\delta) \rangle}{\langle N_+ \rangle} + \frac{\langle N_{-+}(\delta) \rangle - \langle N_{--}(\delta) \rangle}{\langle N_- \rangle} \right\}$$

$N_{+-}$  and  $N_{-+}$  numbers of the unlike-sign pairs

$N_{++}$  and  $N_{--}$  numbers of the like-sign pairs

two members of a pair fall into the rapidity window  $Y$ , their relative rapidity is

$$\delta = \Delta y = |y_2 - y_1|$$

$N_+$  ( $N_-$ ) number of positive (negative) particles in the interval  $Y$

## Two contributions for the $\pi^+\pi^-$ balance function

- 1) **RESONANCE CONTRIBUTION (R)** is determined by the decays of neutral hadronic resonances which have a  $\pi^+\pi^-$  pair in the final state

$$K_S, \eta, \eta', \rho^0, \omega, \sigma, f_0$$

- 2) **NON-RESONANCE CONTRIBUTION (NR)** other possible correlations among the charged pions

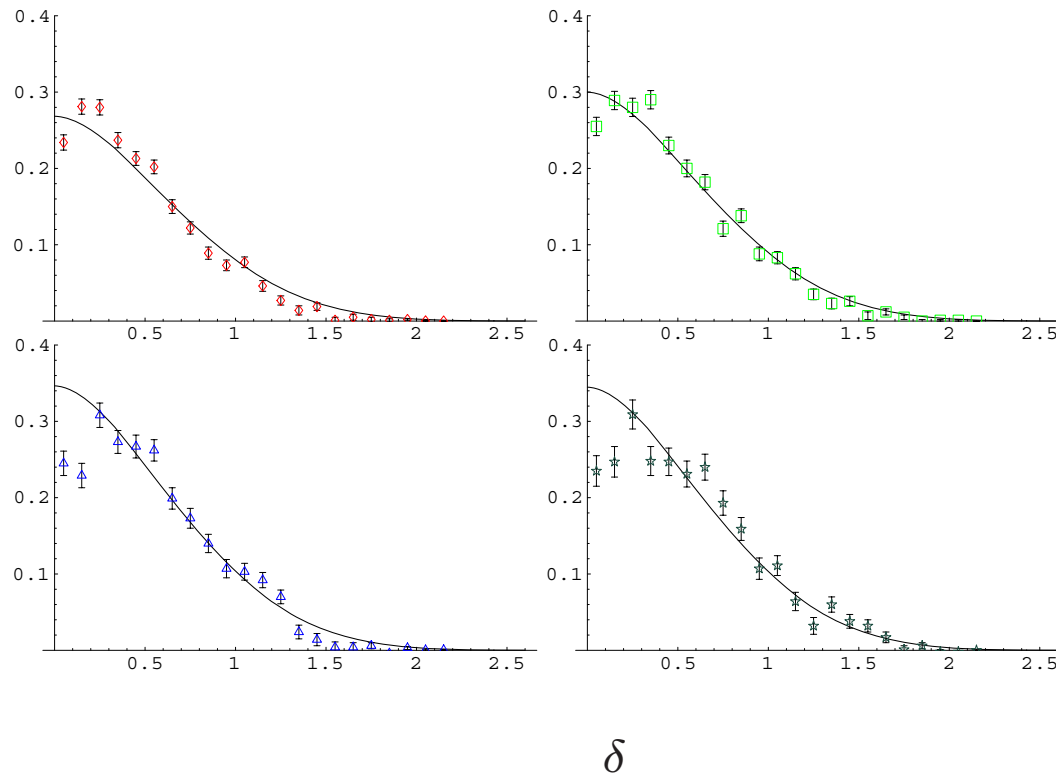
in our approach the non-resonance two-particle distribution is determined by the local relative thermal momenta of particles

The pion balance function is constructed as a sum of the two terms

$$B(\delta, Y) = B_R(\delta, Y) + B_{NR}(\delta, Y)$$

# Fit to the STAR data

$B(\delta)$



four different centralities: 0-10%, 10-40%, 40-70%, 70-96%

rescaling factors: 0.40, 0.44, 0.51, 0.51 ( $\chi^2$  fits)

poor man's way of taking into account the detector efficiency