

# Rechenaufgabe 1

a)  $W = \frac{1}{2} \epsilon_0 \vec{E}^2 = \frac{1}{2} \epsilon_0 \rho(x)\rho(L_1-x)\rho(y)\rho(L_2-y)\rho(z)\rho(L_3-z) (A^2+B^2+C^2)$

(This is because  $\rho^2(x) = \rho(x)$ ).

b) Die gesamte Energie lautet

$$E_{\text{feld}} = \int_{-a}^a dx \int_{-a}^a dy \int_{-a}^a dz W =$$

$$= \int_0^{L_1} dx \int_0^{L_2} dy \int_0^{L_3} dz \frac{1}{2} \epsilon_0 (A^2+B^2+C^2) = \frac{1}{2} \epsilon_0 (A^2+B^2+C^2) L_1 L_2 L_3.$$

c) Da  $\vec{B} = \vec{0}$  hat man  $\vec{S} = \vec{0}$ .

d) Bewegungsgleichung:  $\vec{F} = q\vec{E}$

$$\begin{cases} m\ddot{x} = qA \\ m\ddot{y} = qB \\ m\ddot{z} = qC \end{cases} \quad \mapsto \quad \begin{cases} x = \frac{qA}{m} t^2 + \frac{L_1}{2} \\ y = \frac{qB}{m} t^2 + \frac{L_2}{2} \\ z = \frac{qC}{m} t^2 + \frac{L_3}{2} \end{cases} \quad \text{for } t \in (0, t^*).$$

Um  $t^*$  zu bestimmen, muss man aufpassen. Es muss nämlich gelten, dass:

$$x \leq L_1 \wedge y \leq L_2 \wedge z \leq L_3.$$

$$x \leq L_1 \text{ bedeutet: } \frac{qA}{m} t^2 + \frac{L_1}{2} \leq L_1 \mapsto t^2 \leq \frac{L_1}{2} \frac{m}{qA} \mapsto t \leq \sqrt{\frac{L_1 m}{2 qA}}.$$

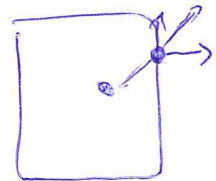
Ähnlich für  $y$  und  $z$ .

Letztendlich bedeckt es:

$$\begin{cases} x = \frac{q_A}{m} t^2 + \frac{L_1}{2} \\ y = \frac{q_B}{m} t^2 + \frac{L_2}{2} \\ z = \frac{q_C}{m} t^2 + \frac{L_3}{2} \end{cases} \quad \text{für } 0 \leq t \leq t_* = \min \left\{ \sqrt{\frac{L_1 m}{2 q_A}}, \sqrt{\frac{L_2 m}{2 q_B}}, \sqrt{\frac{L_3 m}{2 q_C}} \right\}$$

Für  $t > t_*$  geht die Bewegung mit konstanter Geschwindigkeit fort:

$$\begin{cases} x = \frac{2q_A}{m} t_* \cdot t + \left( \frac{q_A}{m} t_*^2 + \frac{L_1}{2} \right) \tilde{x}_1 \\ y = \frac{2q_B}{m} t_* \cdot t + \left( \frac{q_B}{m} t_*^2 + \frac{L_2}{2} \right) \tilde{x}_2 \\ z = \frac{2q_C}{m} t_* \cdot t + \left( \frac{q_C}{m} t_*^2 + \frac{L_3}{2} \right) \tilde{x}_3 \end{cases}$$



Für  $\tilde{x}_1 =$

$$\frac{2q_A}{m} t_*^2 + \tilde{x}_1 = \frac{q_A}{m} t_*^2 + \frac{L_1}{2}$$

$$\tilde{x}_1 = -\frac{q_A}{m} t_*^2 + \frac{L_1}{2}$$

(wenn  $t_* = \sqrt{\frac{L_1 m}{2 q_A}}$ , dann  $\tilde{x}_1 = 0$ )

und so für  $\tilde{x}_2$  und  $\tilde{x}_3$ .

$$A^\mu = \cancel{1} E^\mu \cos\left(\frac{\omega t - kz}{R}\right)$$

a)  $\square A^\mu = 0$

$$\square A^\mu - \partial^\mu (\partial_\nu A^\nu) = 0.$$

$$\partial_\nu A^\nu = 0 \rightarrow \square A^\mu = 0.$$

(siehe b)

$$\square A^\mu = 0 \rightarrow \overset{\text{achtung:}}{\frac{\omega^2}{c^2} = k^2} \rightarrow \boxed{\omega = kc!} \quad (\text{da } \square = \frac{1}{c^2} \partial_\nu^2 - \Delta)$$

b)  $\partial_\nu A^\nu = \partial_y A^3 = \partial_y \left( \cancel{1} \cos\left(\frac{\omega t - kz}{R}\right) \right) = 0.$

c)  $E^i = c F^{i0} = c (\partial^i A^0 - \partial^0 A^i) = c (-\partial_x A^0 - \partial_0 A^i).$

$$A^0 = 0, \text{ ergo}$$

$$E^i = -c \partial_0 A^i.$$

$$A^1 = A^3 = 0. \text{ Nur } A^2 \neq 0. \text{ ergo:}$$

$$E^1 = E^3 = 0.$$

$$E^2 = -c \partial_t \left( \cancel{1} \cos\left(\frac{\omega t - kz}{R}\right) \right) = +c \cancel{1} \cdot \frac{\omega}{R} \sin\left(\frac{\omega t - kz}{R}\right)$$

$$\vec{E} = \left( 0, \frac{c\omega}{R} \sin\left(\frac{\omega t - kz}{R}\right), 0 \right).$$

d)

$$F^{i5} = -\epsilon^{i5K} B^K$$

$$F^{i5} = \partial^i A^5 - \partial^5 A^i$$

$$F^{12} = \partial^1 A^2 - \partial^2 A^1 = \partial^1 A^2 = 0, \text{ da } A^2 \text{ nicht von } x \text{ abhängt}$$

$$F^{13} = \partial^1 A^3 - \partial^3 A^1 = 0$$

$$F^{23} = \partial^2 A^3 - \partial^3 A^2 = -\partial^3 A^2 = \partial_3 A^2 = \partial_z A^2$$

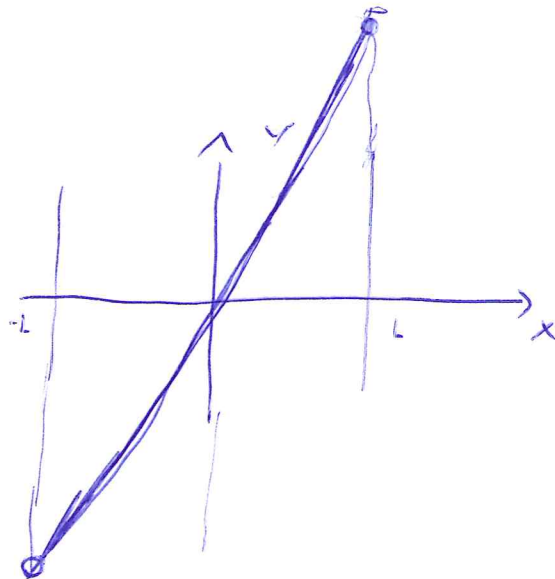
$$= \partial_z \left( b \cos\left(\frac{\omega t - kz}{\hbar}\right) \right) = \frac{bk}{\hbar} \sin\left(\frac{\omega t - kz}{\hbar}\right)$$

$$F^{23} = \epsilon^{231} B^1 = -B^1$$

$$B^1 = -\frac{bk}{\hbar} \sin\left(\frac{\omega t - kz}{\hbar}\right), \quad B^2 = B^3 = 0.$$

$$\rho = \sigma_0 \delta(2x-y) \mathcal{R}(4L^2-z^2) \mathcal{R}(L^2-x^2);$$

1) xy-Ebene  
 $|x| < L$

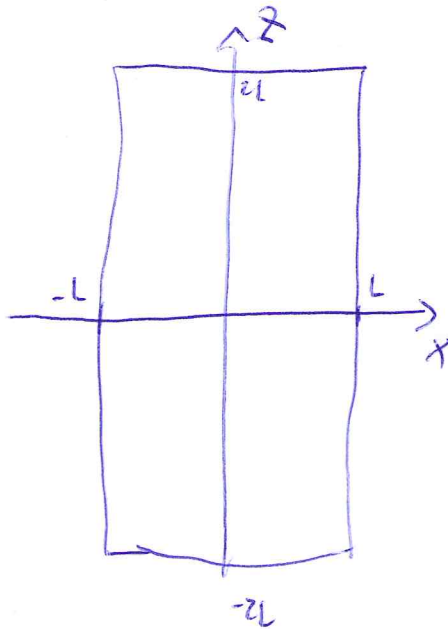


$$2x = y$$

2) xz-Ebene

$$|z| < 2L$$

$$|x| < L$$



c) 
$$Q = \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dz \sigma_0 \delta(2x-y) \mathcal{R}(4L^2-z^2) \mathcal{R}(L^2-x^2)$$

$$= \sigma_0 \left( \int_{-2L}^{2L} dz \right) \cdot \left[ \underbrace{\int_{-L}^L dx \mathcal{R}(L^2-x^2)}_{2L} \underbrace{\int_{-\infty}^{\infty} \delta(2x-y) dy}_1 \right] = \sigma_0 \cdot 4L \cdot 2L = 8\sigma_0 L^2.$$

check

$$= \sigma_0 \cdot 4L \cdot 2L.$$

$$\int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dx \operatorname{sech}(L^2 - x^2) \frac{1}{2} \delta(x - \frac{y}{2})$$

$$|x| < L$$

$$\frac{1}{2} \text{ for } |y| < 2L$$

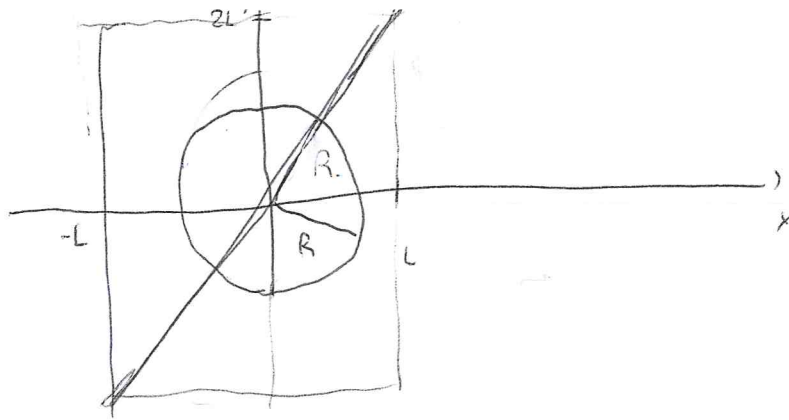
$$\int_{-\infty}^{\infty} dx \operatorname{sech}(L^2 - x^2) \frac{1}{2} \delta(x - \frac{y}{2}) =$$

$$0 \text{ for } |y| > 2L$$

Each wire

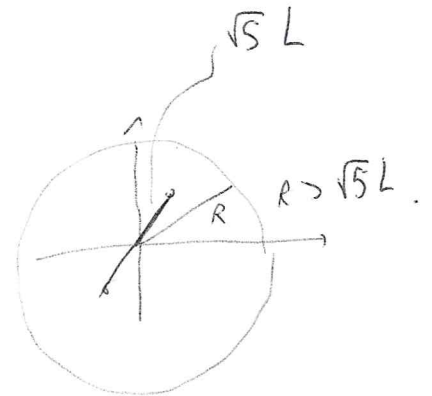
$$\frac{1}{2} \int_{-2L}^{2L} dy = 2L$$

OK.

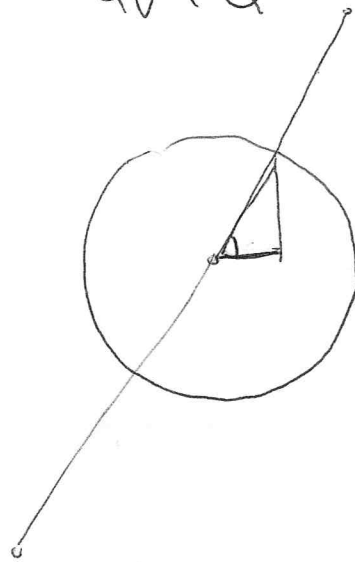


$$Q_v = \text{?}$$

$$\text{if } R > \sqrt{5} L \quad Q_v = Q \rightarrow \varphi = \frac{Q}{\sigma_0}$$



$$\text{if } R < \sqrt{5} L \quad Q_v < Q$$



$$x_{\max} = R \cos \varphi \quad \text{aber } \cos \varphi = \frac{1}{\sqrt{5}} \rightarrow x_{\max} = \frac{R}{\sqrt{5}}$$

$$Q_v = \sigma_0 \int_{-2L}^{2L} \int_{-\infty}^0 dx \cdot \frac{1}{\sqrt{5}} \left( \frac{R^2}{5} - x^2 \right) = \sigma_0 \cdot 4L \cdot \left( \frac{2R}{\sqrt{5}} \right) = 8 \sigma_0 L \frac{R}{\sqrt{5}} \quad \text{for } R < \sqrt{5} L$$

check: For  $R = \sqrt{5} L$   $Q_v = Q$  as expected!